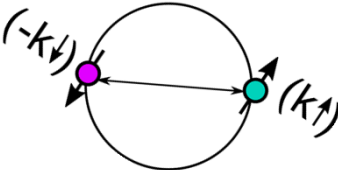

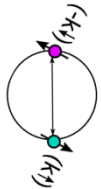
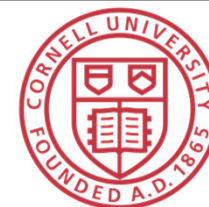


A Three-Fluid Model of Dissipation at Surfaces in Superconducting Radiofrequency Cavities

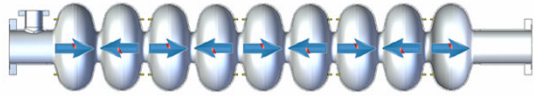


11am Tuesday June 27th, 2023; SRF 2023 Grand Rapids, MI

Michelle M. Kelley, S. Deyo, N. Sitaraman, T. Oseroff, D. B. Liarte,
M. Liepe, J. P. Sethna, T. A. Arias (Cornell University)

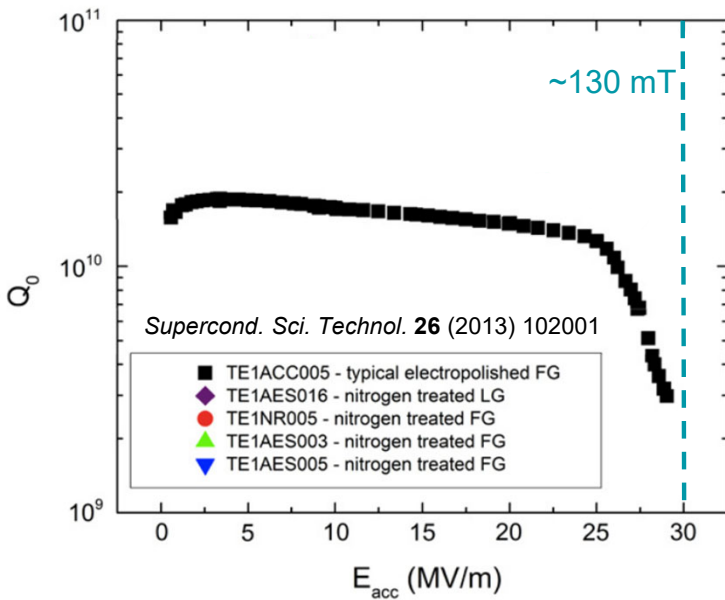


Field dependence of efficiency (quality factor)



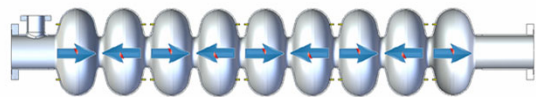
Superconducting radiofrequency (SRF) cavity

$$Q \equiv 2\pi \frac{\text{energy stored}}{\text{energy dissipated per RF cycle}}$$



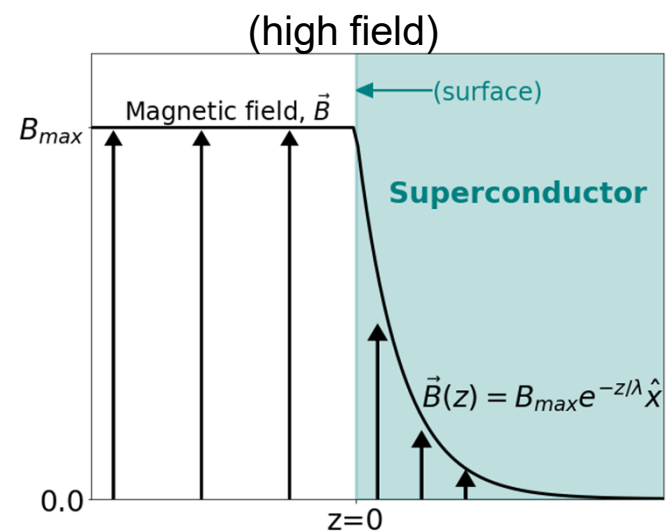
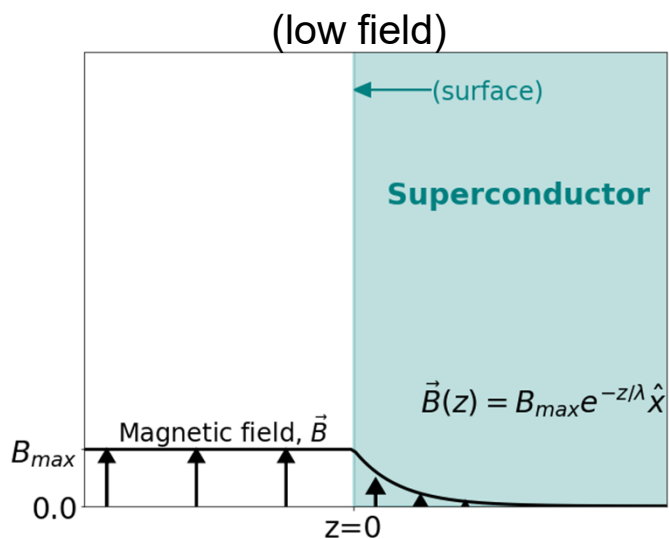
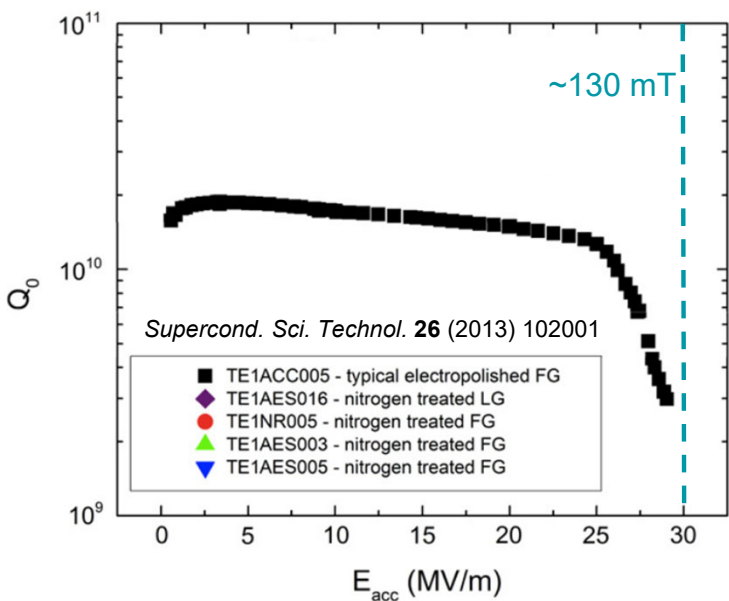


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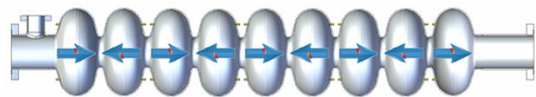
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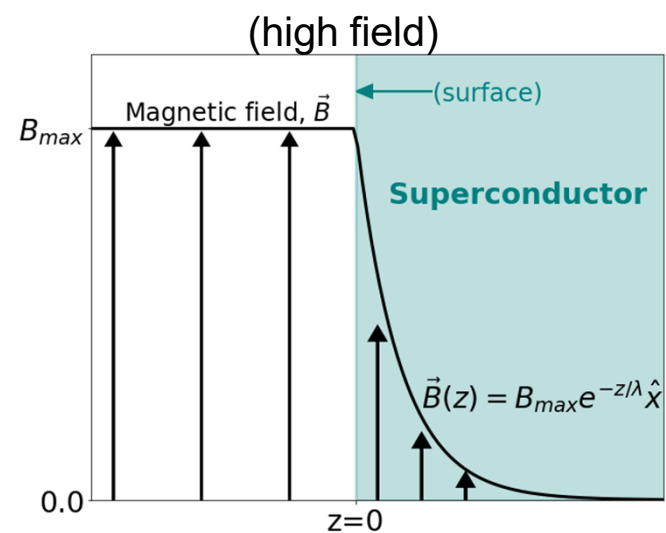
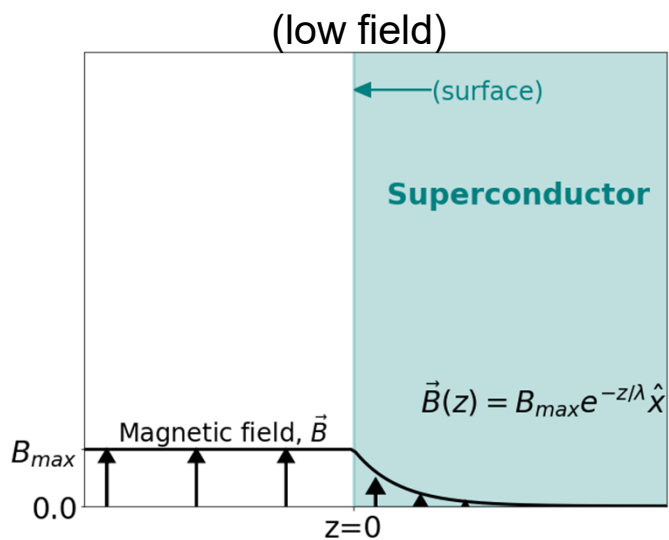
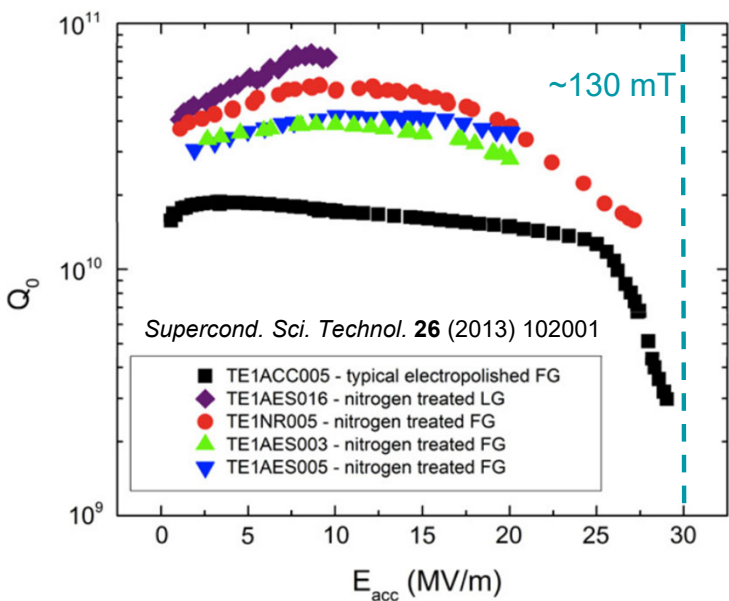


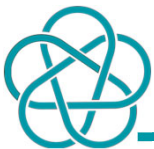
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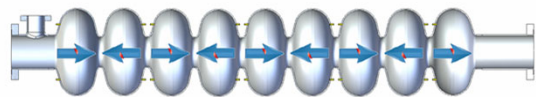
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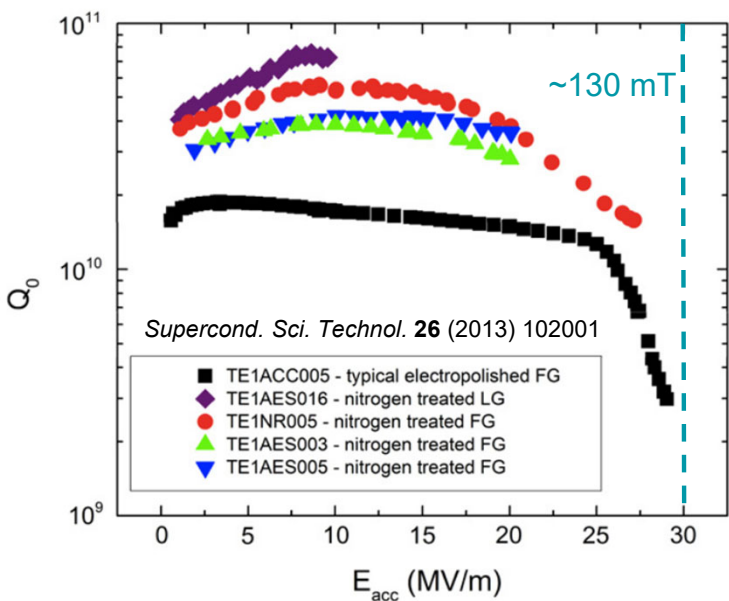


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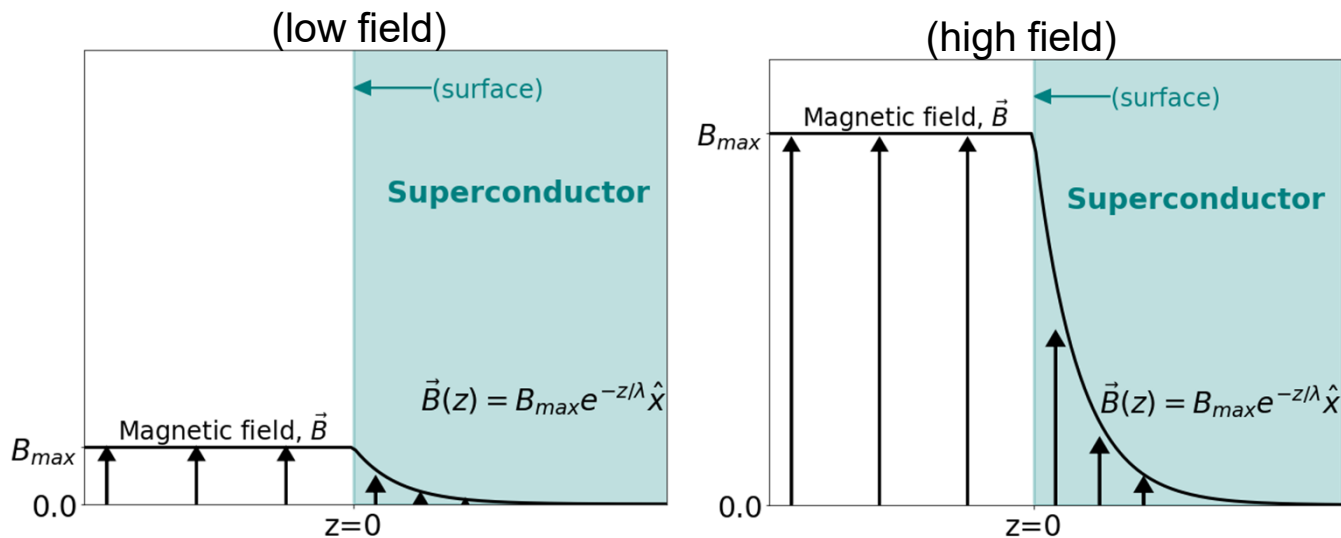


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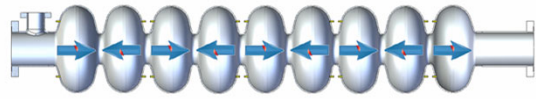


How can stronger fields increase the efficiency of an SRF cavity???



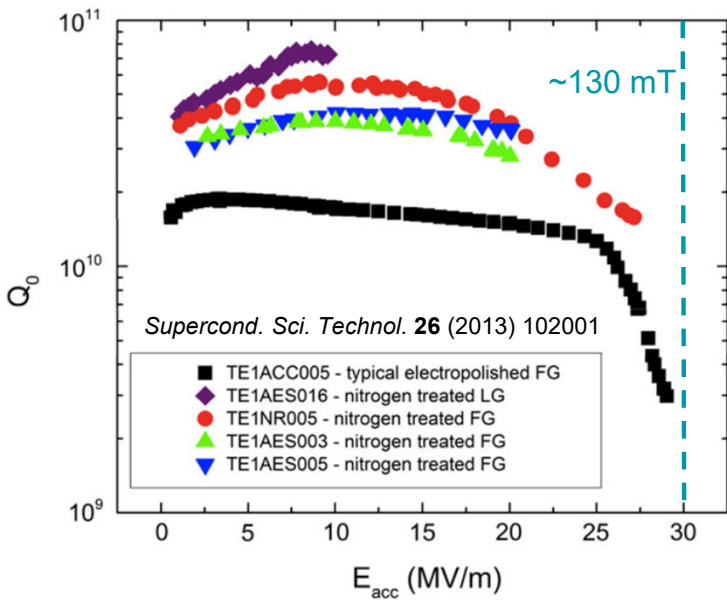


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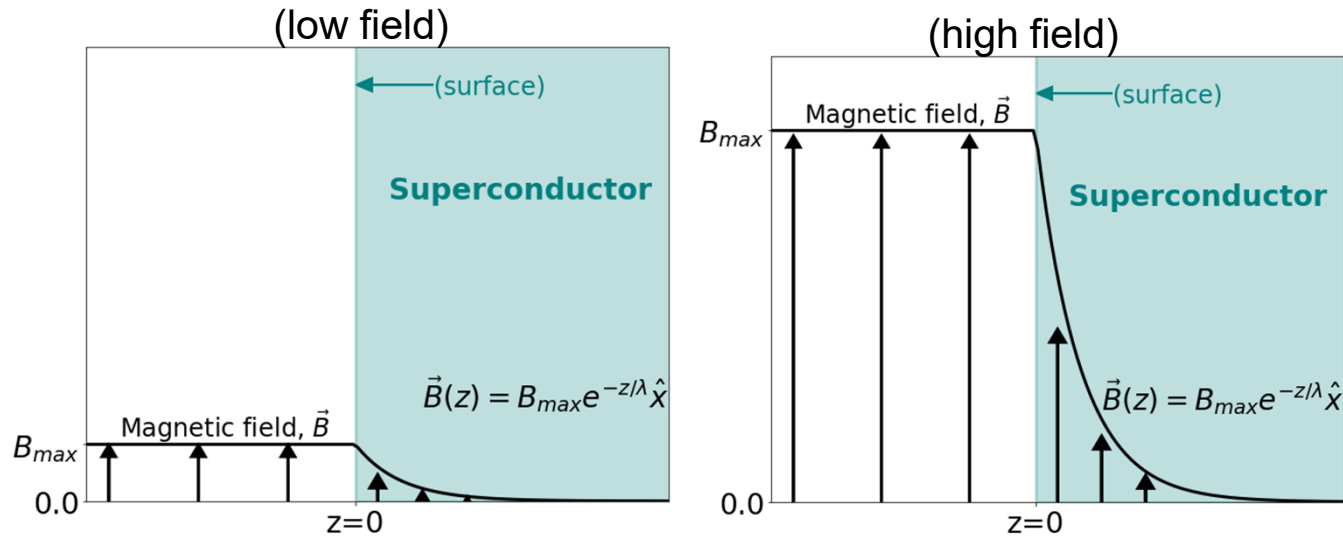


Superconducting radiofrequency (SRF) cavity

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How can stronger fields increase the efficiency of an SRF cavity???



→ need to study dissipation at superconducting surfaces



Superconducting RF dissipation & Bogoliubov quasiparticles



Bogoliubov transformation

$$\Psi(r \uparrow) = \sum_n \left(\gamma_{n\uparrow} u_n(r) + \gamma_{n\downarrow}^\dagger v_n^*(r) \right)$$

$|u|^2 + |v|^2 = 1$

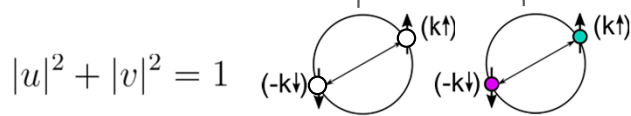


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Bogoliubov-de Gennes
self-consistent field method

$$\begin{aligned} (H_e + U)u + \Delta v &= Eu \\ -(H_e^* + U)v + \Delta^* u &= Ev \end{aligned}$$

$$\begin{aligned} U &= -V \sum_n |u_n|^2 f_n + |v_n|^2 (1 - f_n) \\ \Delta &= V \sum_n u_n v_n^* (1 - 2f_n), \end{aligned}$$

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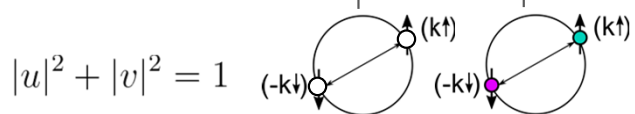


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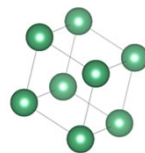
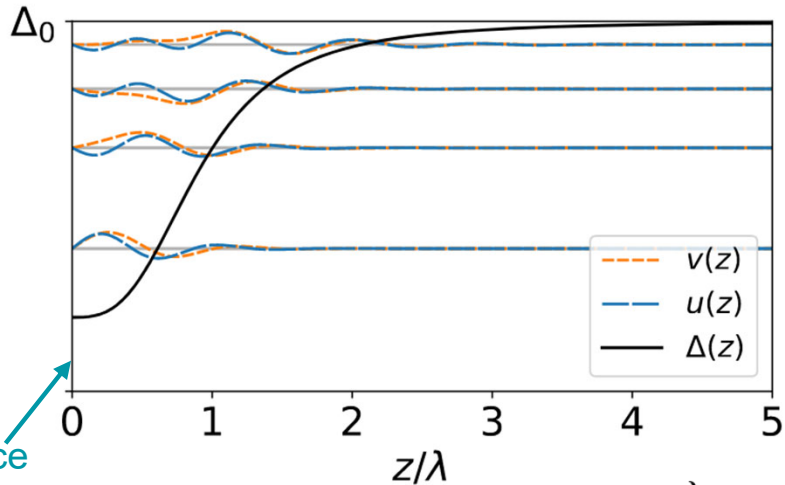
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Solve for spatially inhomogeneous energy gap, $\Delta(z)$



$$\lambda_{\text{Nb}} \sim 40 \text{ nm} \gg \sim 0.33 \text{ nm}$$



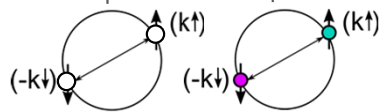
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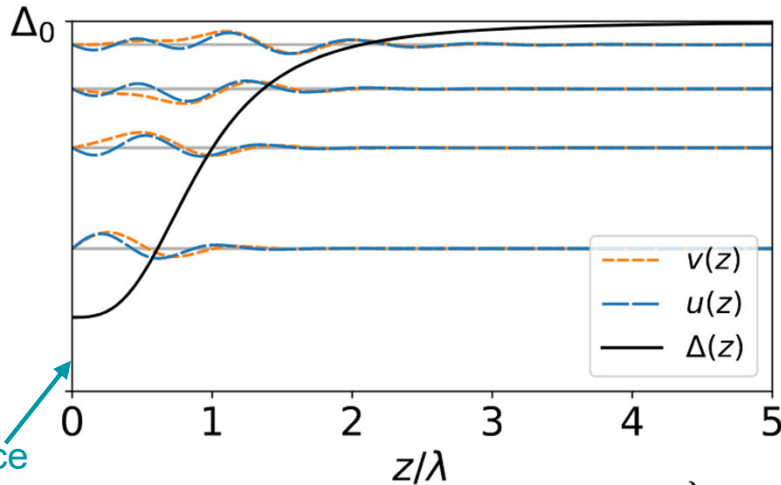
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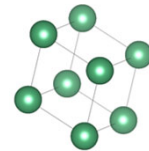
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Ohmic dissipation

$$\begin{aligned} P_{\text{diss}} &= \int dV \Re \left[\frac{1}{\sigma} \right] j^2(r) \\ \mathbf{j} &= \frac{1}{2m} [(\Psi^* \hat{\mathbf{p}} \Psi - \Psi \hat{\mathbf{p}} \Psi^*) - 2q\mathbf{A}|\Psi|^2] \end{aligned}$$



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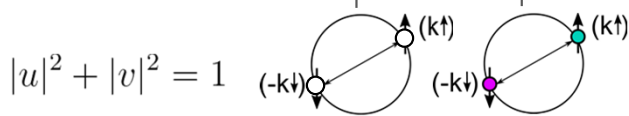


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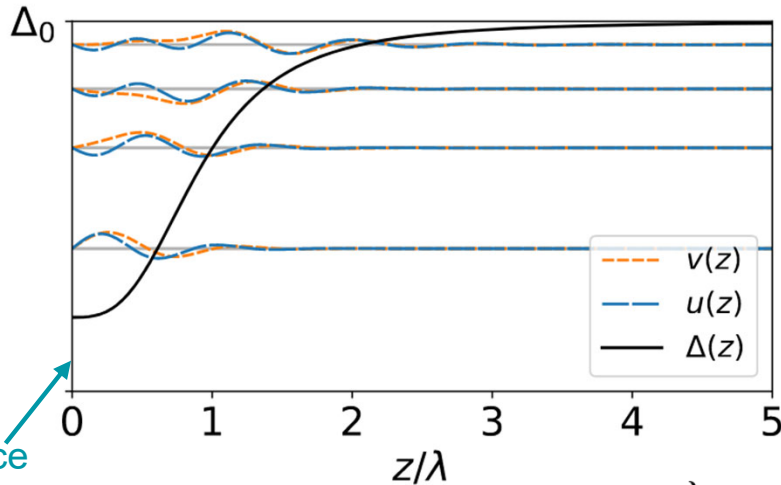
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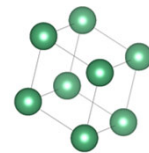


Surface

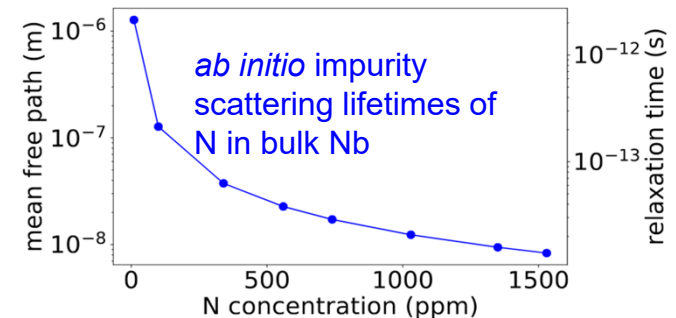
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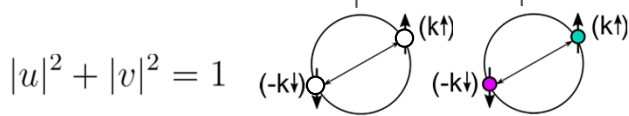


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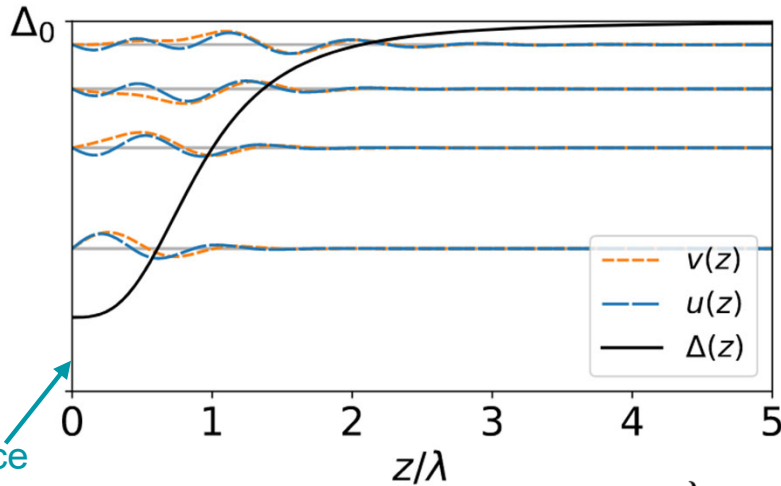
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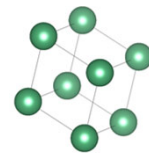
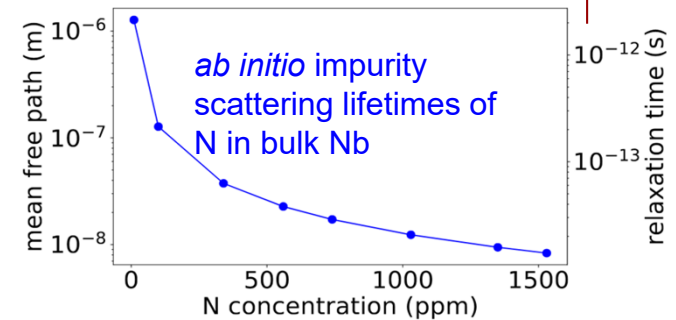
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much faster than RF period (~1 ns)



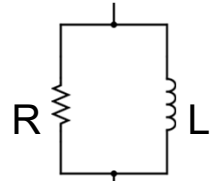
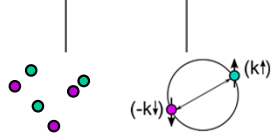
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A “three”-fluid model with Bogoliubov quasiparticles



Two-fluid
model

$$n_{\text{tot}} = n_n + n_s$$



$$\sigma = \frac{n_n e^2 \tau}{m} + i \frac{n_s e^2}{m \omega}$$

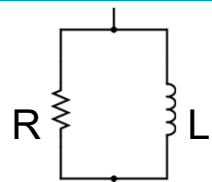
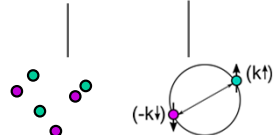


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BCS prediction

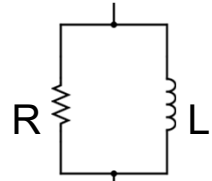
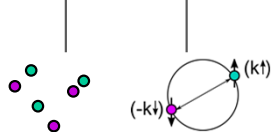
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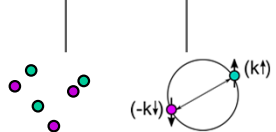
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A "three"-fluid model with Bogoliubov quasiparticles



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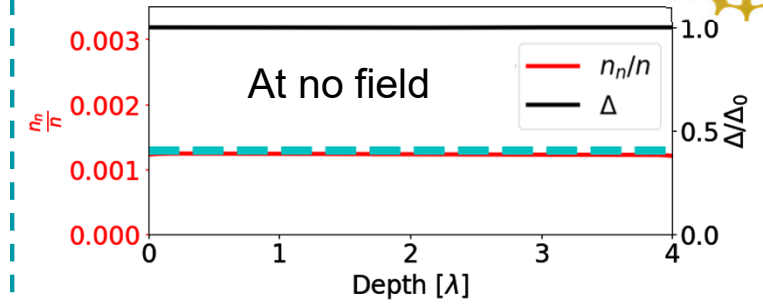


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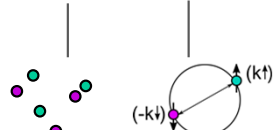


A "three"-fluid model with Bogoliubov quasiparticles

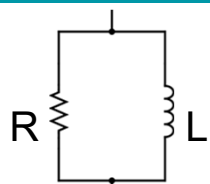


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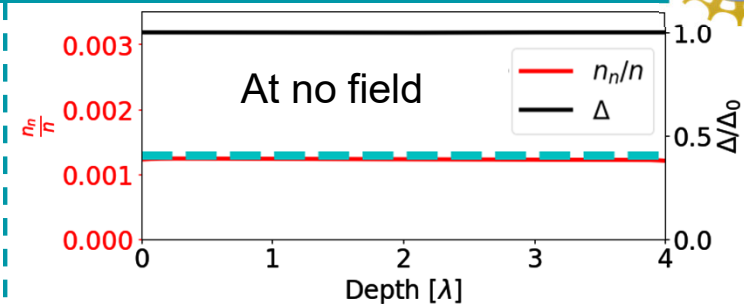


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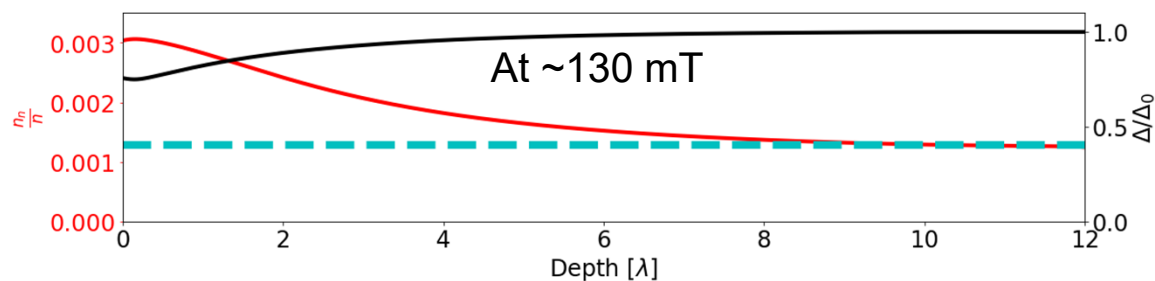


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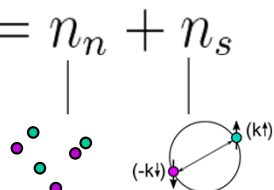


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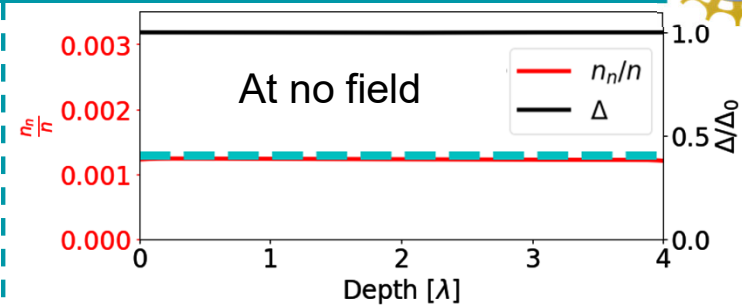


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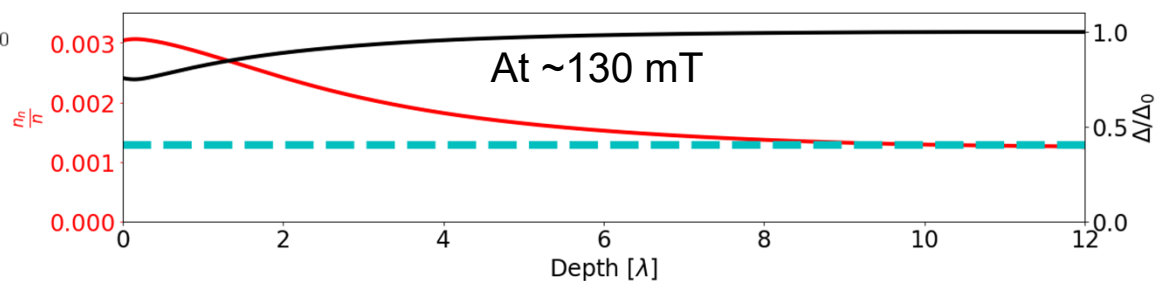
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bound $\sum_{i: |E_i| < \Delta_0}$ or continuum $\sum_{i: |E_i| \geq \Delta_0}$

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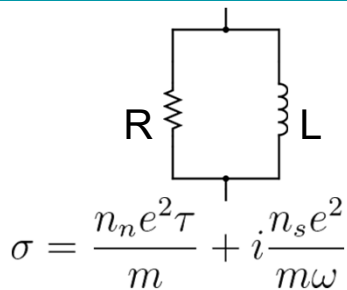
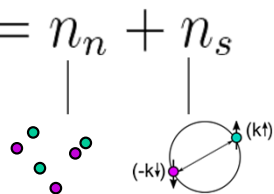


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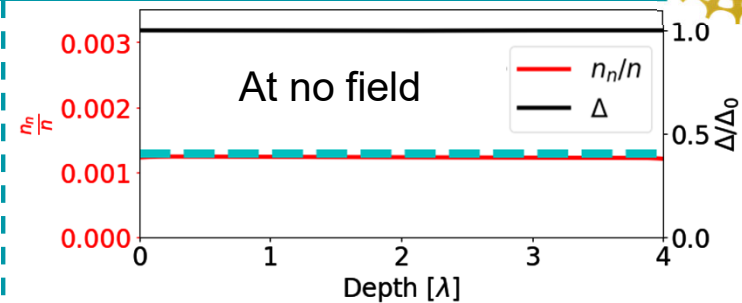
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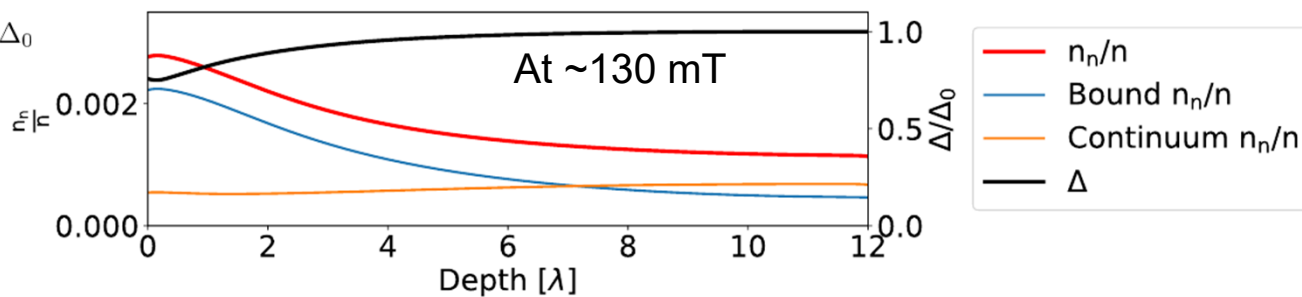
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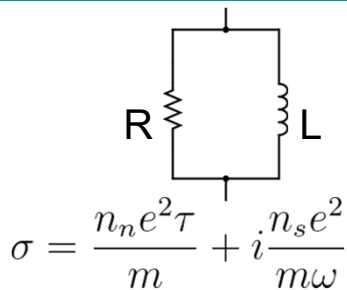
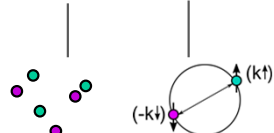


A "three"-fluid model with Bogoliubov quasiparticles



Two-fluid model

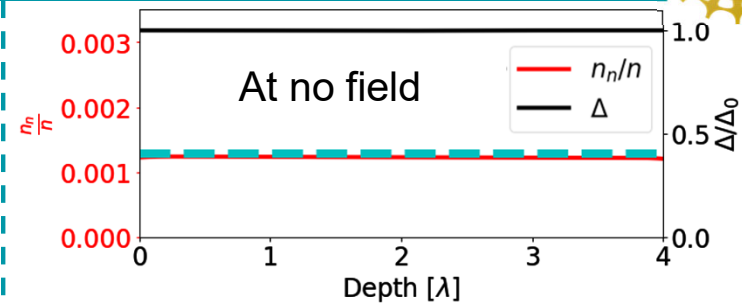
$$n_{\text{tot}} = n_n + n_s$$



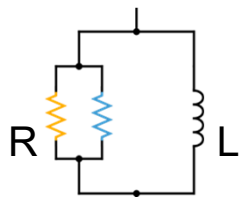
$$\sigma = \frac{n_n e^2 \tau}{m} + i \frac{n_s e^2}{m \omega}$$

$$\frac{n_n}{n_{\text{tot}}} = \frac{1}{N(0)} \sum_i \left(-\frac{\partial f(E_i)}{\partial E_i} \right)$$

BCS prediction
 $E_i = \sqrt{\xi_i^2 + \Delta^2}$



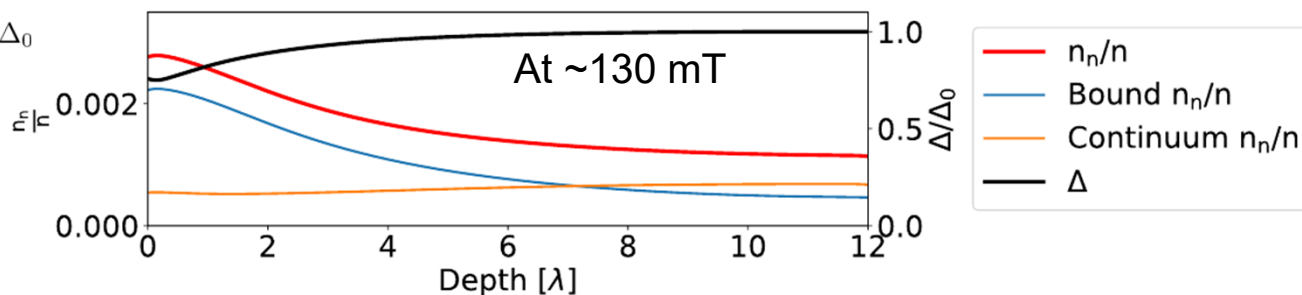
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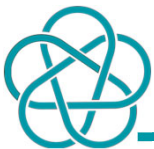


$$\sigma = \frac{e^2(\tau_b n_b + \tau_{ub} n_{ub})}{m} + i \frac{e^2 n_s}{m \omega}$$

bound $\sum_{i: |E_i| < \Delta_0}$ or continuum $\sum_{i: |E_i| \geq \Delta_0}$

$$\frac{n_n}{n_{\text{tot}}}(z) = \frac{1}{N(0)} \sum_i \left(-\frac{\partial f(E_i)}{\partial E_i} \right) [|u_i(z)|^2 + |v_i(z)|^2]$$



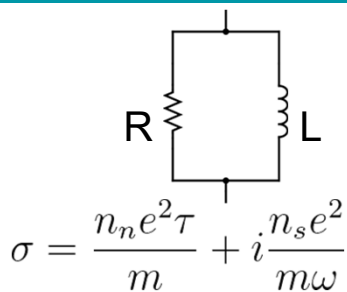
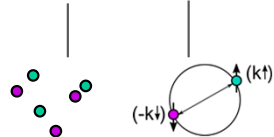


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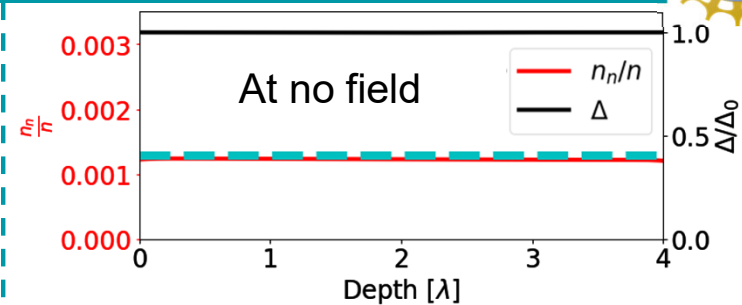
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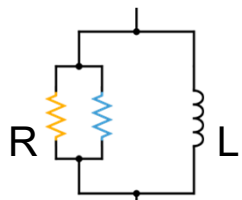


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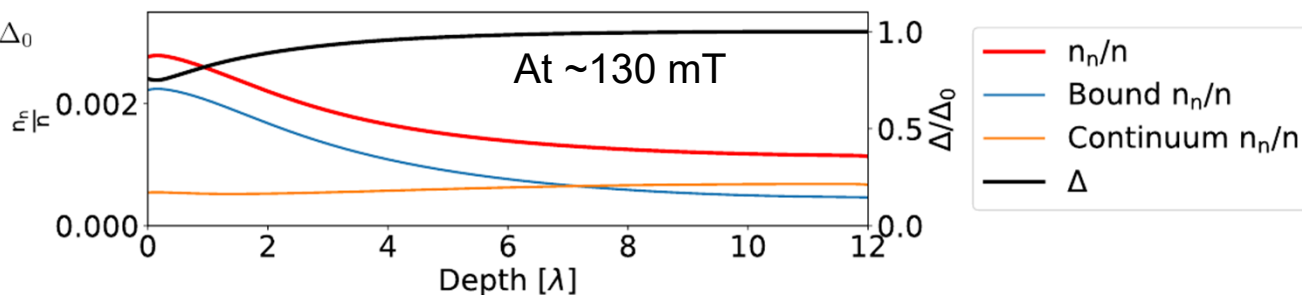


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BdG quasiparticle-impurity scattering

$$\tau_{n\mathbf{k}_{\parallel}}^{-1} = \frac{1}{N_{\mathbf{k}_{\parallel}}} \sum_{\mathbf{k}'_{\parallel}} \frac{2\pi}{\hbar} \alpha^2 n_{\text{imp}}(z_0) (1 - f_{n'\mathbf{k}'_{\parallel}}) \delta(E_{n\mathbf{k}_{\parallel}} - E_{n'\mathbf{k}'_{\parallel}}) |u_{n\mathbf{k}_{\parallel}}(z_0) u_{n'\mathbf{k}'_{\parallel}}(z_0) + v_{n\mathbf{k}_{\parallel}}(z_0) v_{n'\mathbf{k}'_{\parallel}}(z_0)|^2$$

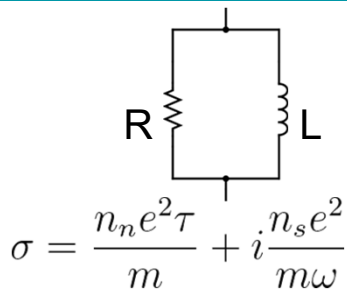
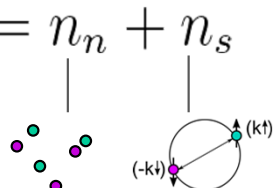


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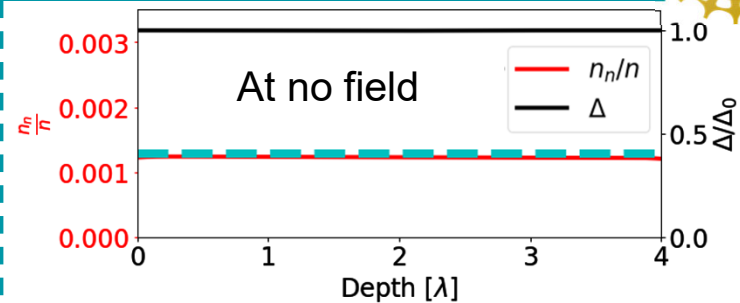
Two-fluid model

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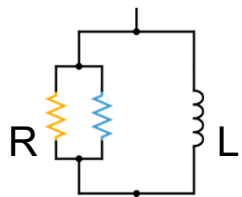


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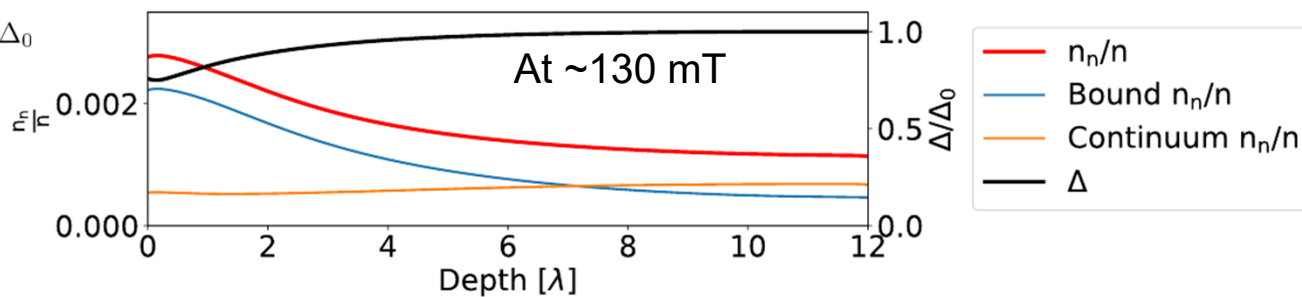


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potential strength (*ab initio*)
& impurity concentration

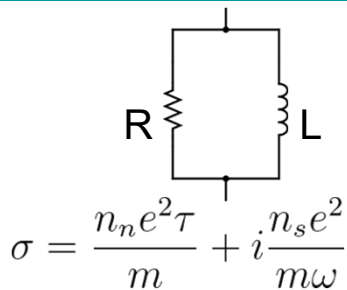
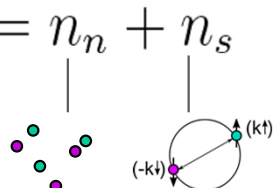


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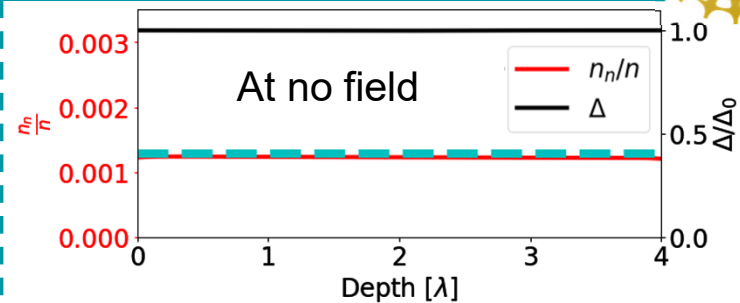
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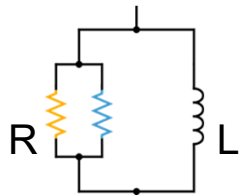


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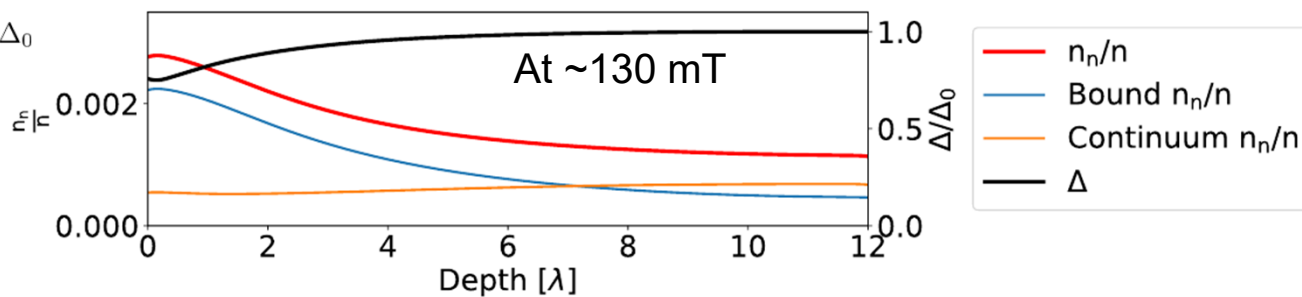


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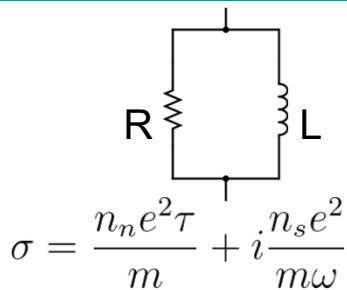
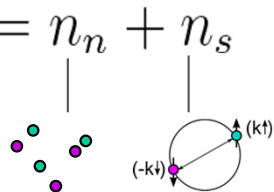


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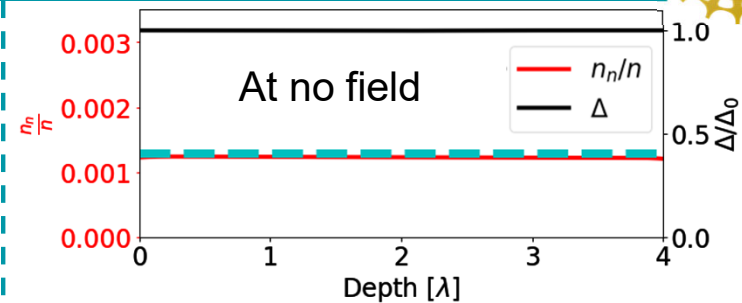
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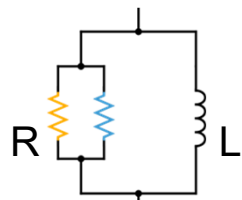


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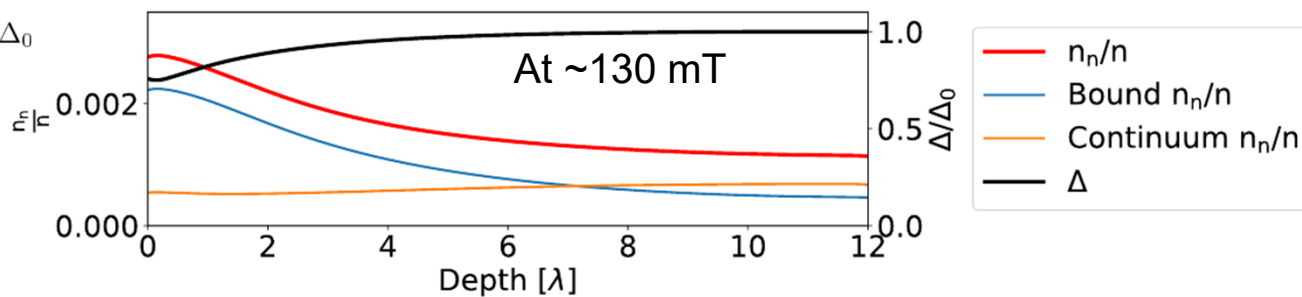


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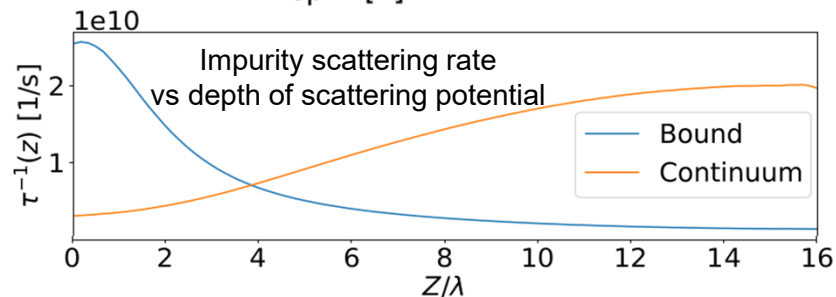


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potential strength (*ab initio*) & impurity concentration

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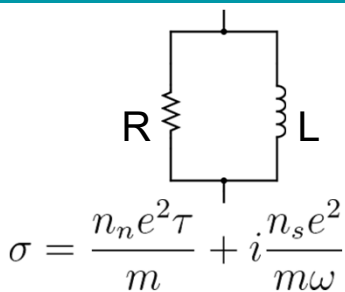
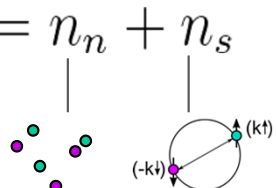


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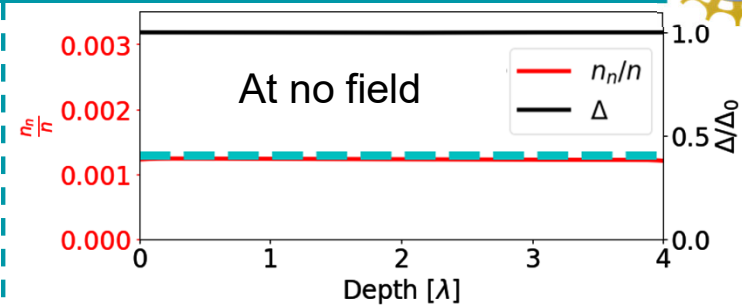
Two-fluid model

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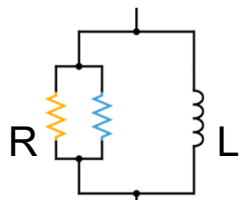


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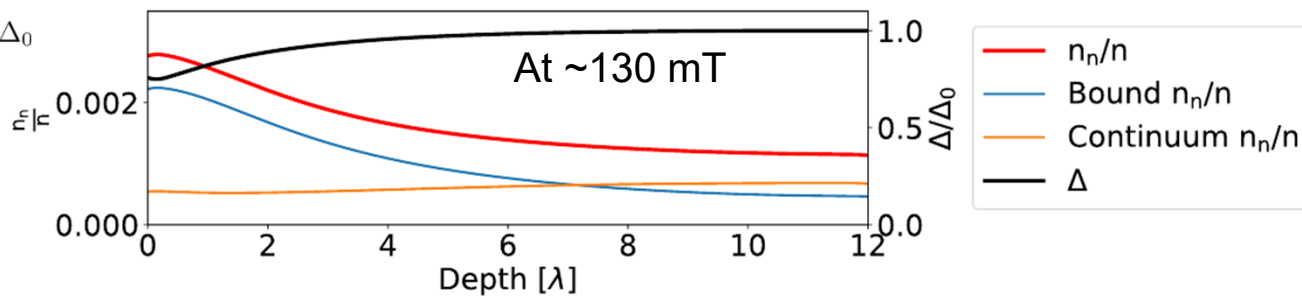


“Three”-fluid model



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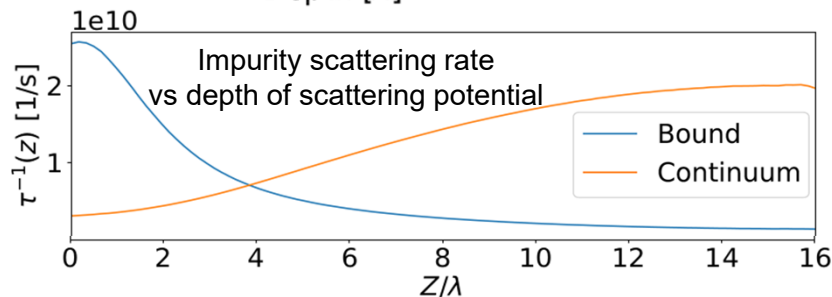


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potential strength (*ab initio*) & impurity concentration

coherence factor



A “dirty” surface can enhance scattering rates of the surface-bound states



“Three”-fluid model results



PHYSICAL REVIEW B **106**, 104502 (2022)

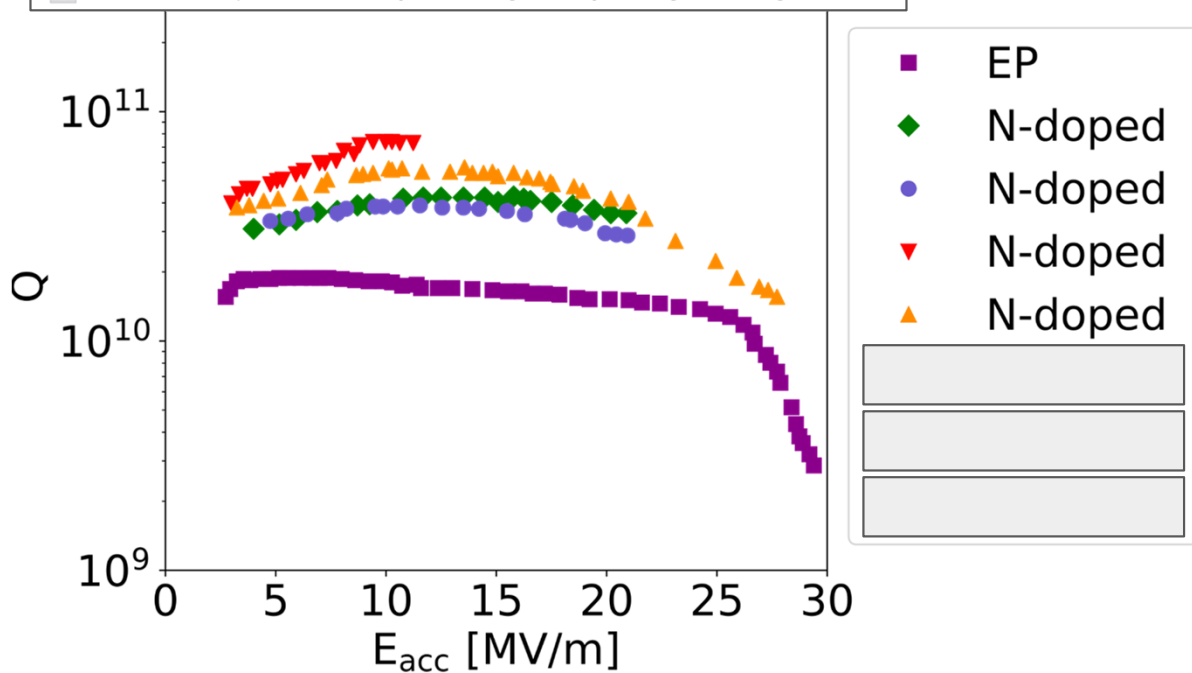
Dissipation by surface states in superconducting radio-frequency cavities

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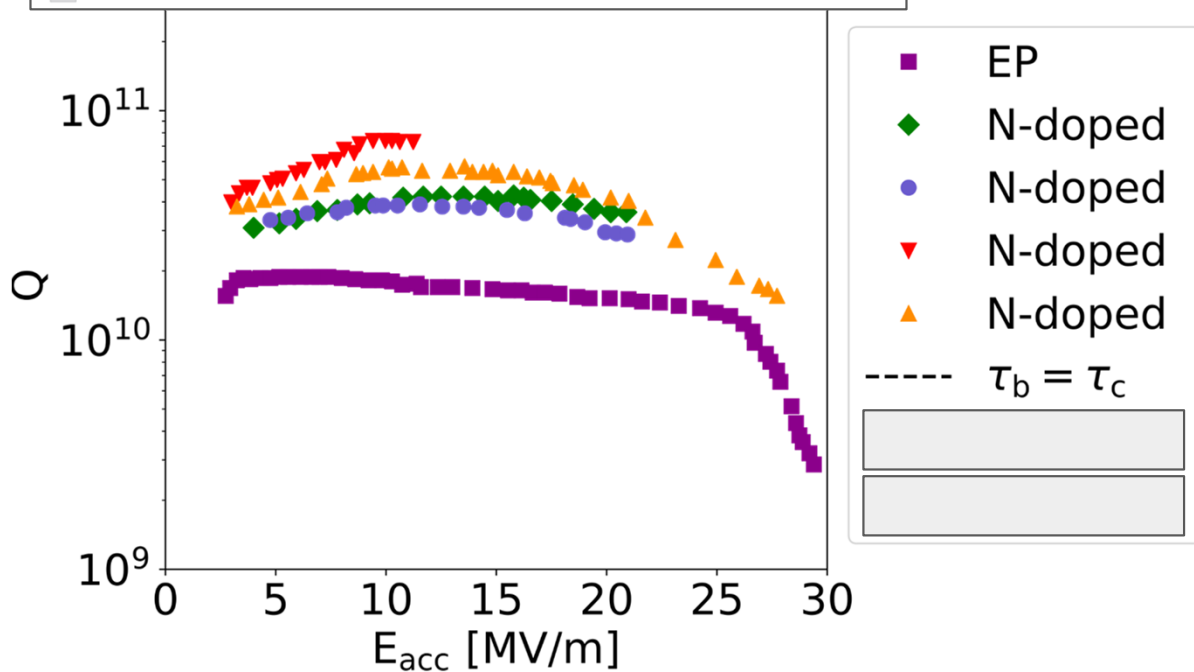
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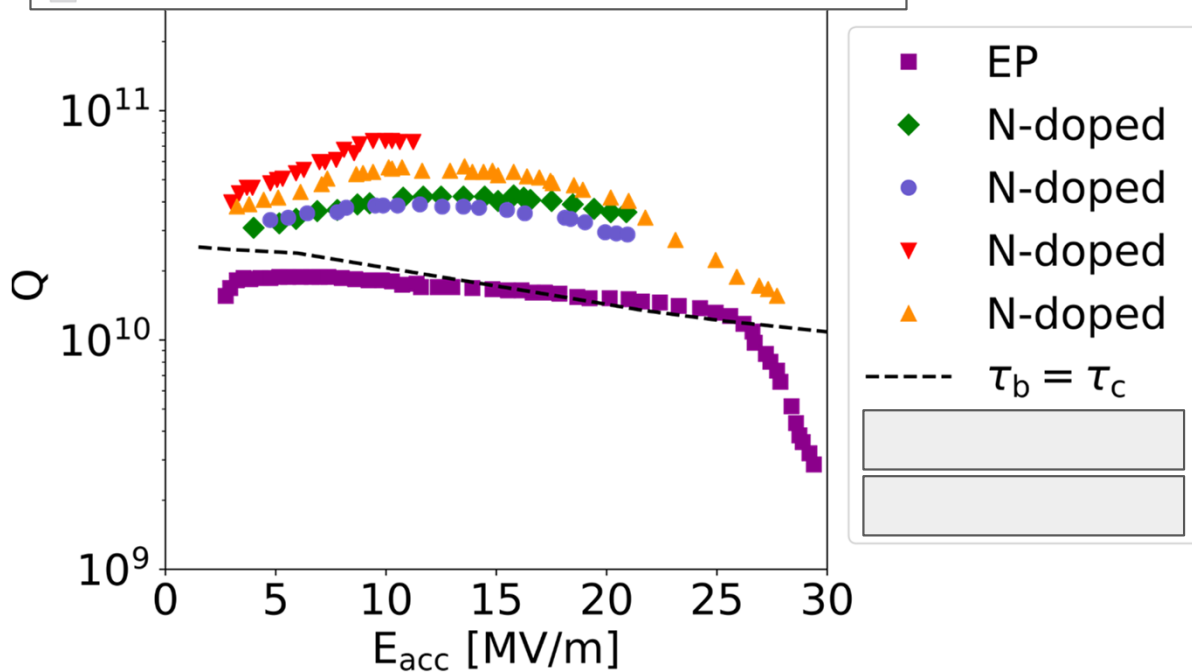
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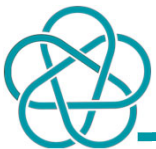
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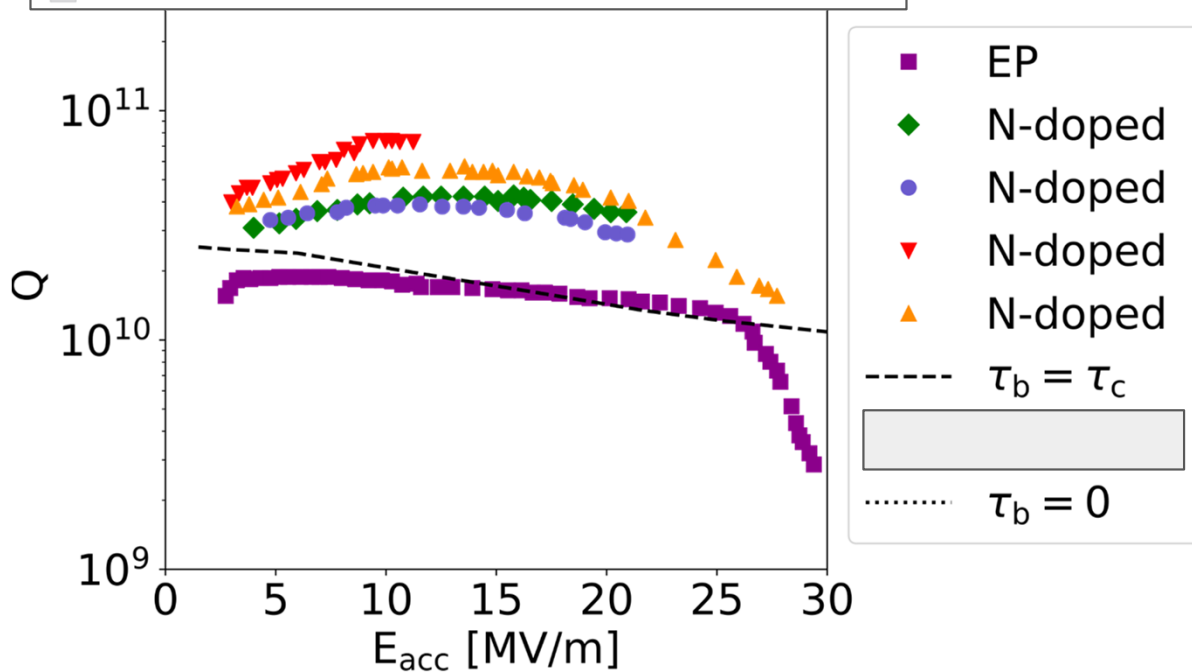
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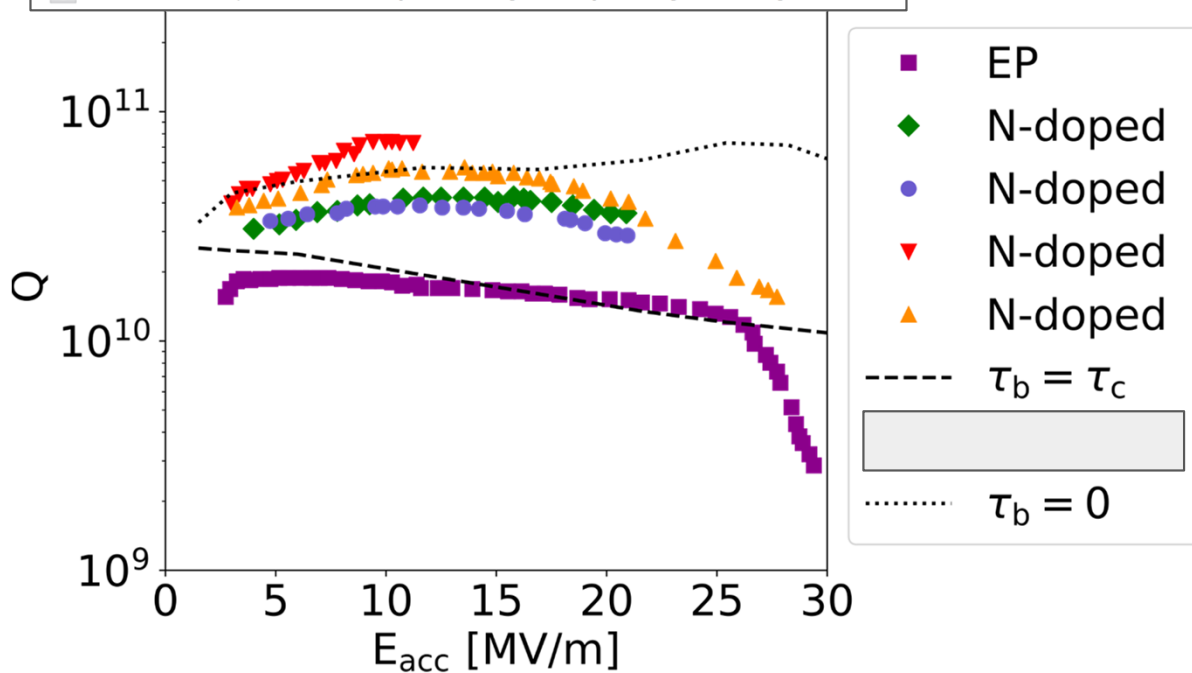
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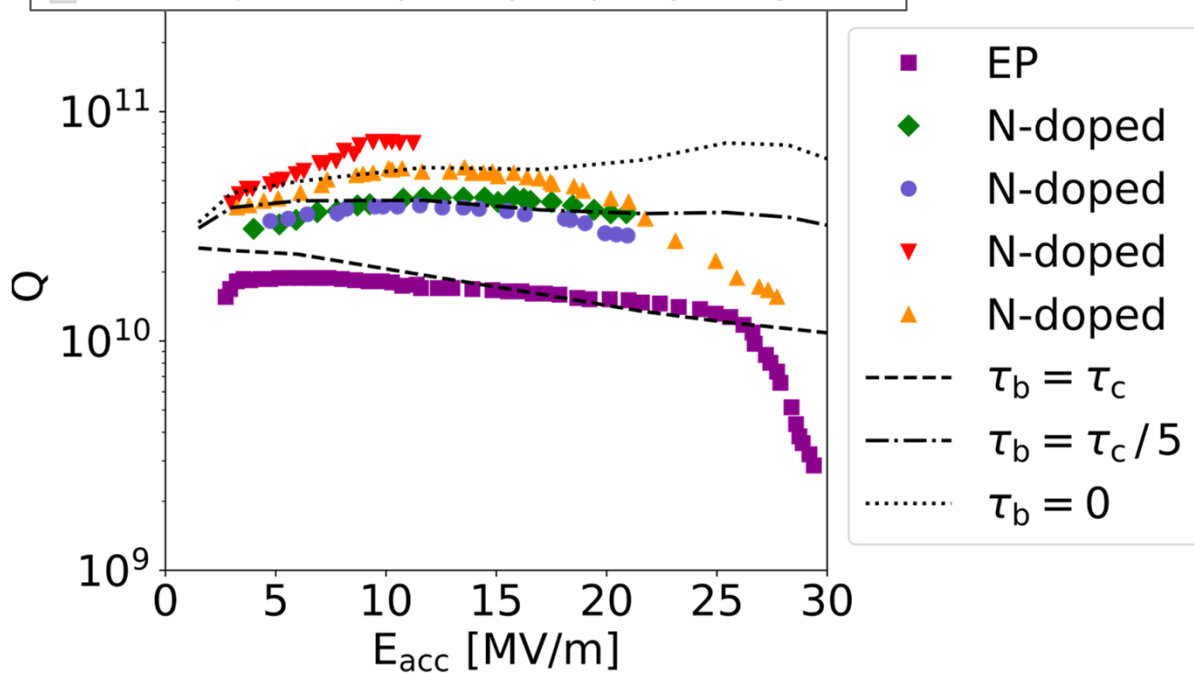
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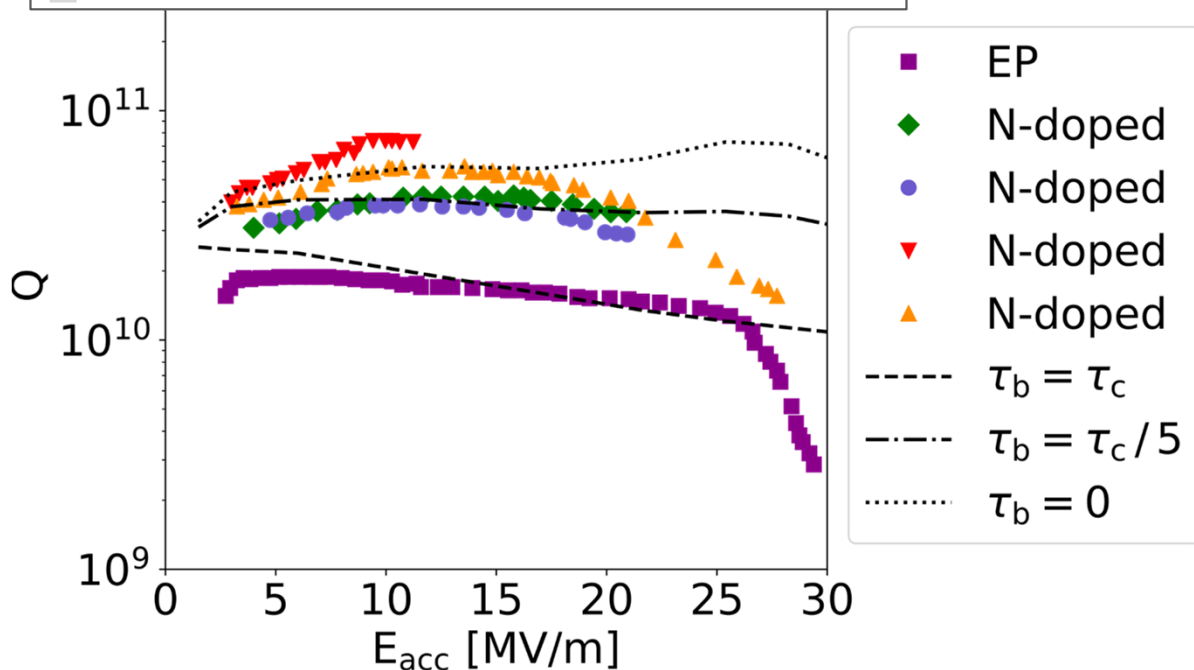
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Questions?



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