

Superconducting Non-Elliptical Cavities (TEM Cavity Designs)

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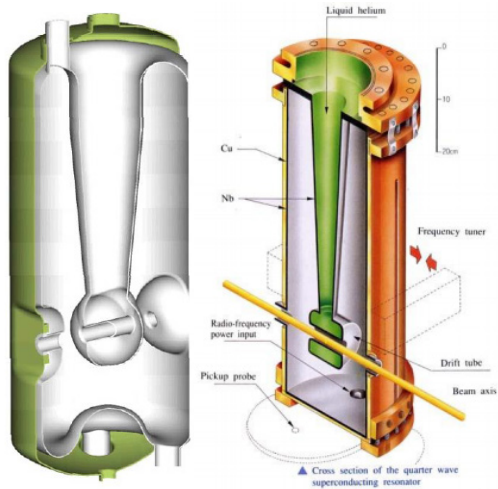
Outline

- Why Non-Elliptical Cavities?
- Electromagnetic modes
- Non-Elliptical accelerating cavities
 - Basic principles
 - Types of cavities
 - Current accelerator applications
- Non-Elliptical deflecting and crabbing cavities
 - Basic principles
 - Types of cavities
 - Current collider applications
- Design issues for non-elliptical cavities

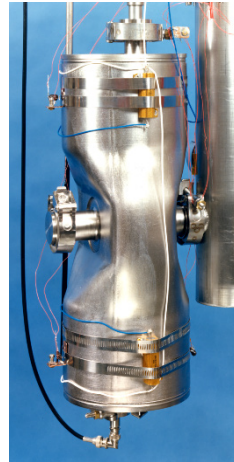
World of Superconducting Non-Elliptical Cavities

RF Cavities of interesting shapes for particle acceleration

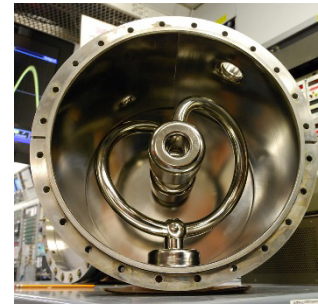
Quarter Wave Cavities



Half Wave Cavities



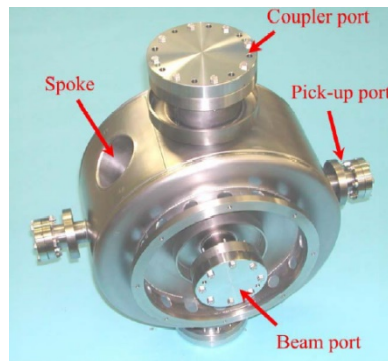
Split Ring Resonator



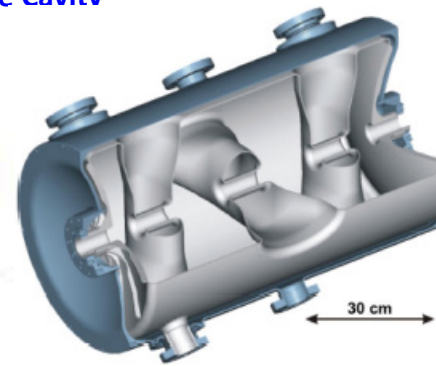
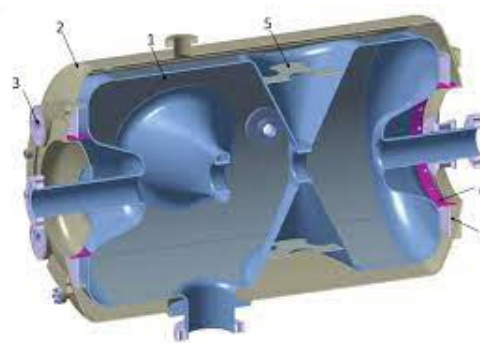
Superconducting RFQ Cavity



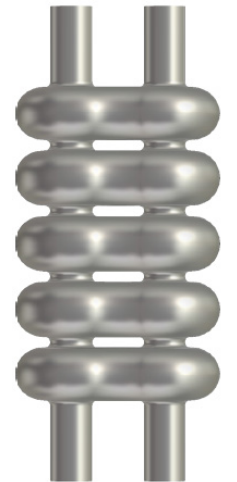
Single Spoke Cavities



Multi Spoke Cavity



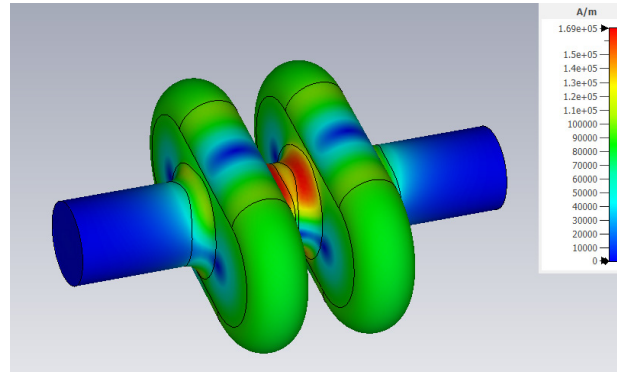
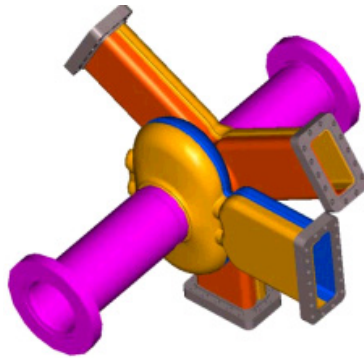
Twin Axis Cavity



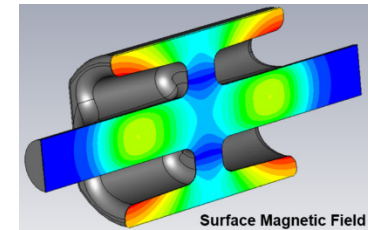
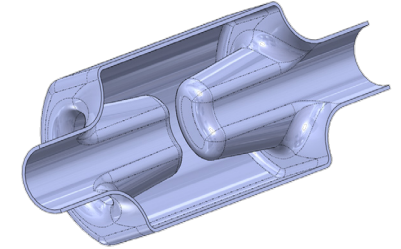
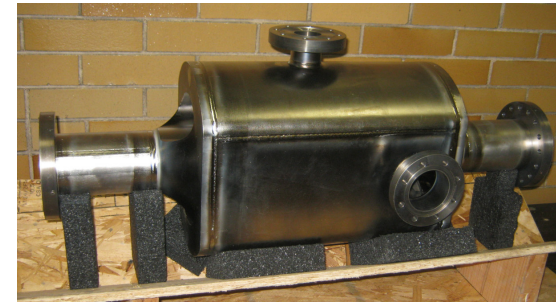
World of Superconducting Non-Elliptical Cavities

RF Cavities of interesting shapes for deflecting and crabbing applications

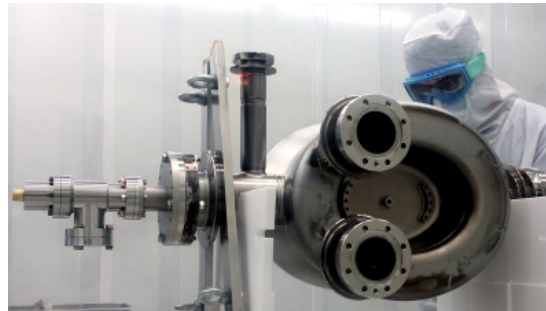
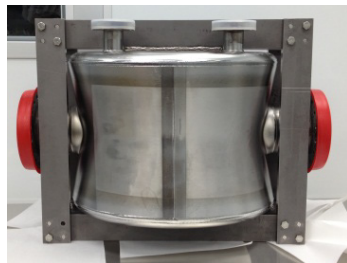
Squashed Elliptical Cavities



4-Rod Cavity



Double Quarter Wave Cavity



RF-Dipole Cavities



Electron vs Proton/Heavy Ion Acceleration

- Electrons and protons are distinctly different due to mass of the particles

- Electrons $\rightarrow 0.511 \text{ MeV}/c^2$

- 550 keV – $\gamma = 2.08$, $\beta = 0.88$
- 550 MeV – $\gamma = 1077$, $\beta = 1$

$$E_T = \gamma m_0 c^2 = \text{KE} + E_0$$

$$\text{KE} = (\gamma - 1) m_0 c^2$$

- Protons $\rightarrow 938 \text{ MeV}/c^2$

- 550 keV – $\gamma = 1.0005$, $\beta = 0.032$
- 550 MeV – $\gamma = 1.59$, $\beta = 0.78$
- 550 GeV – $\gamma = 587.4$, $\beta = 1$

$$\beta = \sqrt{1 - \frac{1}{\gamma^2}}$$

- For heavy ions: Mass number (A) > 1

$$E_T = \text{KE} + AE_0$$

$$\text{KE} = (\gamma - 1) Am_0 c^2$$

$$p = \gamma \beta Am_0 c = \frac{\gamma \beta AE_0}{c}$$

- Protons $\rightarrow A = 1$ and $q = 1$

$$\text{KE} = qV_{\text{eff}} \cos \phi$$

$$B\rho = mv = \frac{p}{q} = \frac{\gamma \beta E_0}{c}$$

- Heavy ions $\rightarrow A > 1$ and q

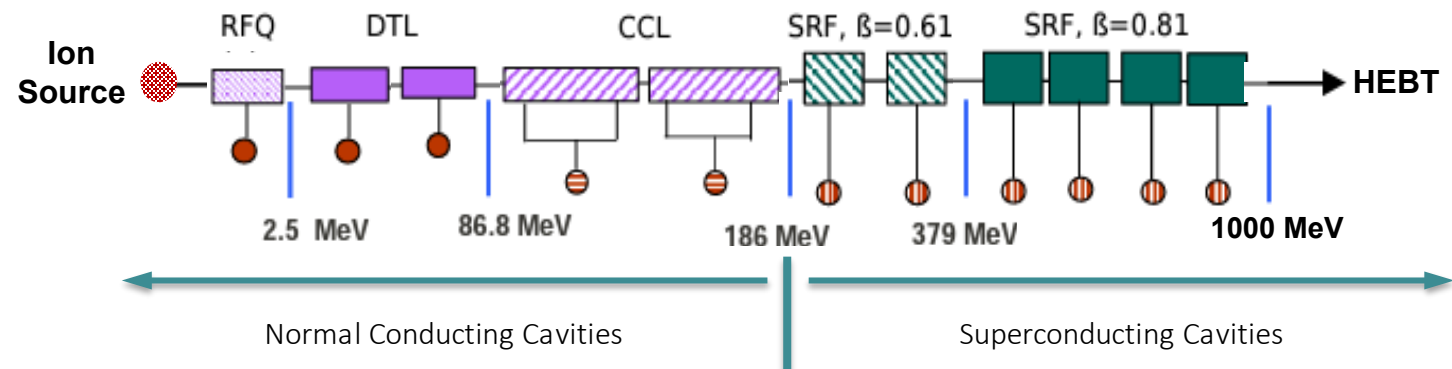
$$\text{KE} = \frac{q}{A} V_{\text{eff}} \cos \phi$$

$$B\rho = p = \frac{A}{q} \frac{\gamma \beta E_0}{c}$$

Why Non-Elliptical Cavities?

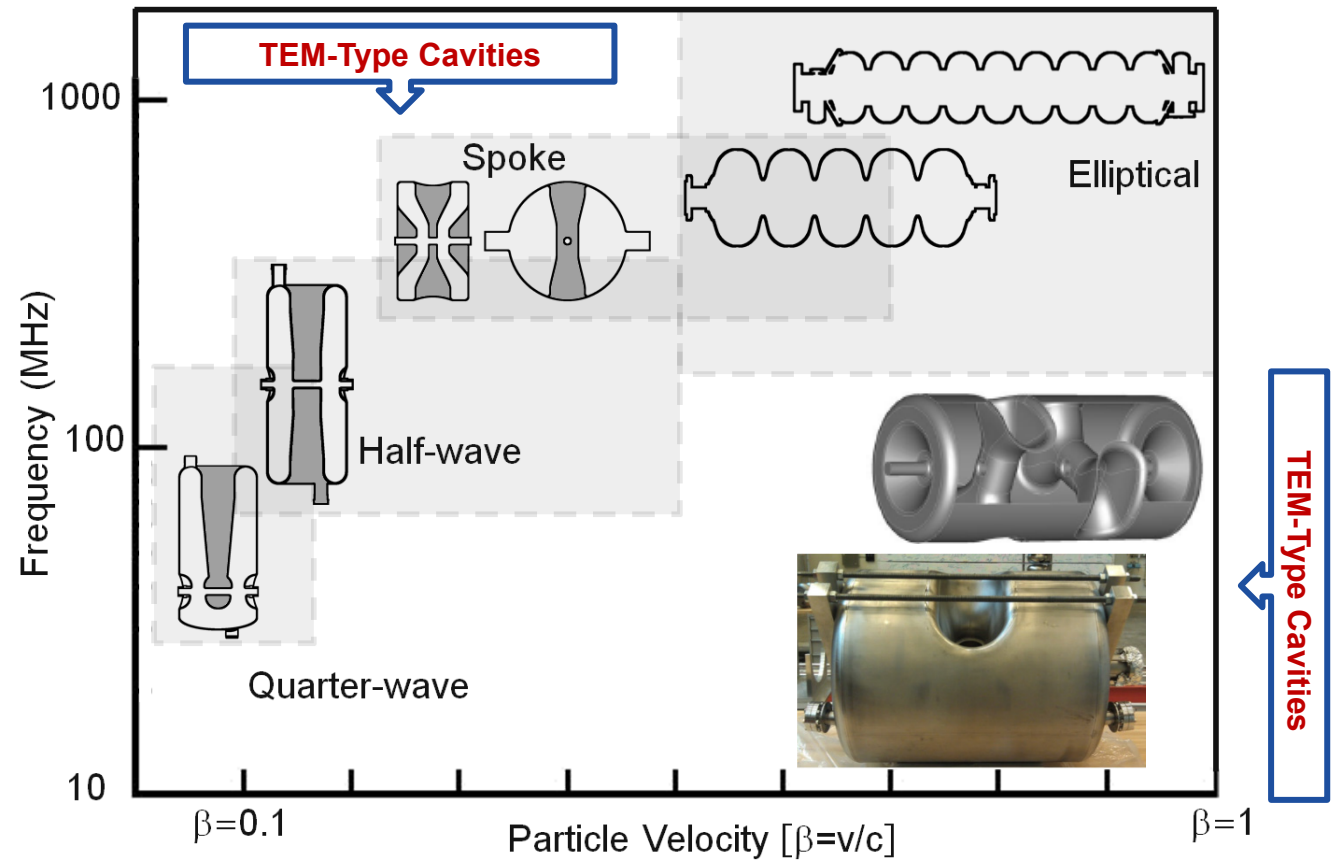
1. To accelerate Protons/Heavy Ions → Need $\beta < 1$ ($\beta = v/c < 1$) accelerating cavities

- In accelerating electrons, all cavities are designed at $\beta = 1$
- In Hadron acceleration
 - Velocity of the particle beam increases with energy
 - Depends on the synchronization of the particle bunches and rf voltage in the cavity
 - Require various types of cavities each optimized to accelerate different velocity ranges
- Elliptical cavities has intrinsic problem as β goes down
 - Due to mechanical problems, multipacting, low RF efficiency
- Solution: Use of TEM-type cavities



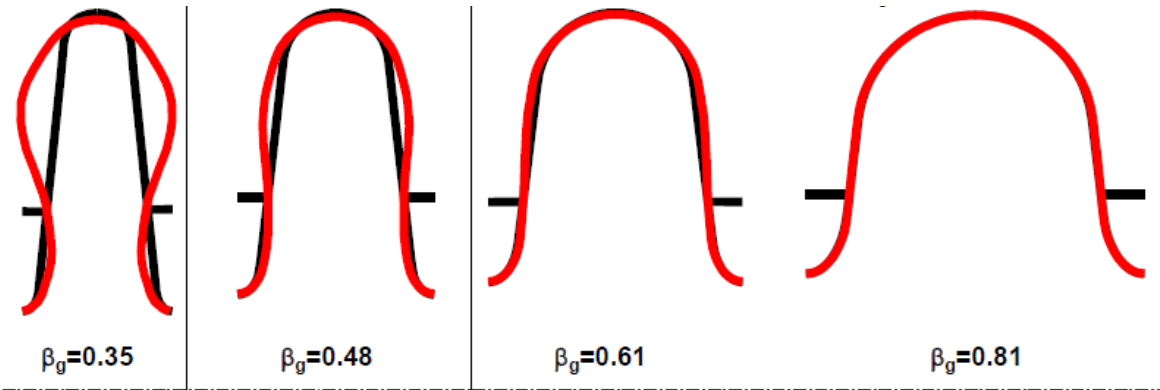
Proton/Heavy Ion Acceleration

- Protons/Heavy Ions require acceleration by many cavities to reach velocities approaching speed of light
- Cavities are designed for specific regions of velocity



Limitation on $\beta=1$ Elliptical Cavities

- Elliptical cavities have been designed for $\beta > 0.5$ for cw applications and $\beta > 0.6$ for pulsed high energy acceleration
- At very low β elliptical cavities start to look like bellows
 - In π mode cell-to-cell distance $\sim \beta\lambda/2$ and cavity diameter is $\sim \lambda$
 - Ratio of cavity length/diameter $\sim \beta/2$



- Low rf efficiency
- Poor mechanical stability
- Possibility of strong multipacting

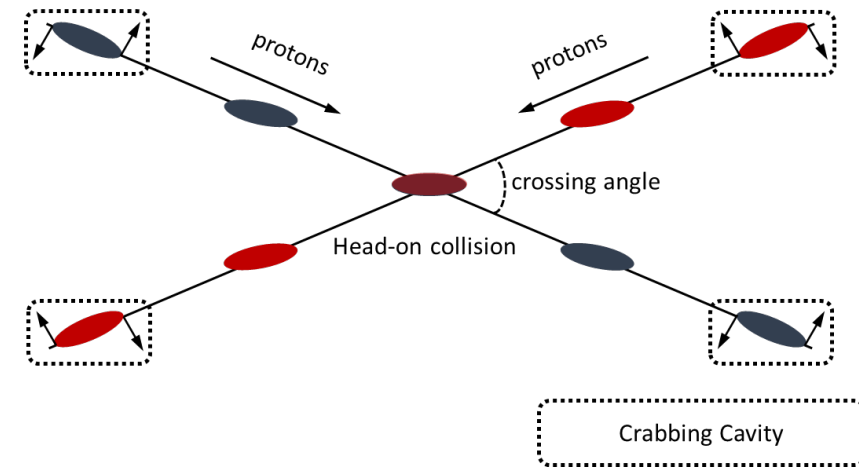
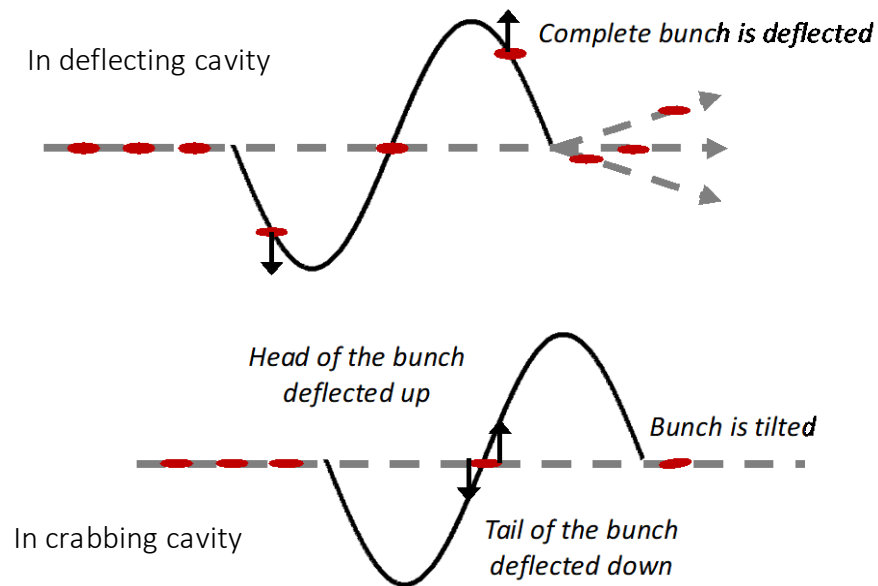
- Will work in CW applications
- Pessimistic in pulsed applications

Suitable for all CW and pulsed applications

Why Non-Elliptical Cavities?

2. To deflect or crab a beam

- Requires to provide a transverse kick to the beam
- Standard elliptical cavities operating in TM_{010} mode only produce a longitudinal gradient
- Deflecting/crabbing cavities operate mostly at $\beta=1$
- Solution: Use of TE_{11} -like mode, TM_{11} -type mode, or TEM-type mode cavities



Applications of Non-Elliptical Cavities

Application	Maximum β	Beam	Maximum Current	Operation
Linacs for nuclear physics research	~ 0.2 (0.5)	Light & Heavy Ions	$\sim 1 \mu\text{A}$	cw
Drivers for radioactive ion beam (RIB) facilities and accelerator driven systems (ADS)	$\sim 0.3 - 0.9$	Light & Heavy Ions	$\sim 0.1 - 30 \text{ mA}$	cw
Linacs for radioisotope production	~ 0.3	p, d	$\sim 1 - 10 \text{ mA}$	cw
Neutron spallation sources	~ 1	p	$\sim 10 - 100 \text{ mA}$	pulsed
Accelerators for material irradiation	~ 0.3	d	$\sim 100 \text{ mA}$	cw
Compact high β linacs (proposed)	1	e	$\sim 1 \text{ mA}$	cw
Deflecting and crabbing applications	1	e, p	$\sim 1 \text{ A}$	cw

Applications of Non-Elliptical Cavities

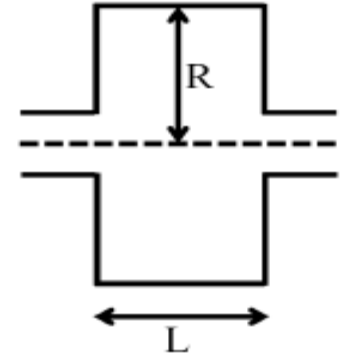
- Superconducting technology allows cw and high duty cycle operation
 - Also allows increase bore (transverse acceptance) as highest shunt impedance is not essential and TEM cavities allow lower frequencies with the associated larger longitudinal acceptance
- Drivers – RIB production (ISOL, fragmentation) (ions), ADS (transmutation, energy) (p, H-, d), Spallation neutron sources (p, H-)
 - Longer machines typically, large velocity swing, several cavity regimes
 - Treat as almost fixed gradient machines
 - Beam loss (halo) an issue, careful beam dynamics required, favor symmetric rather than asymmetric cavities
 - Beam loading is typically an important consideration
- Post-accelerators (Radioactive ion beam and nuclear physics) (ions)
 - Shorter machines typically, broad velocity acceptance
 - Utilize maximum cw gradient to improve performance and/or reduce cost
 - Short independently phased cavities give flexibility to beam delivery
 - Beam loading typically not an issue

Advantages of Non-Elliptical Cavities

- Traditionally low β superconducting resonators were quarter wave (or split rings) used as post-accelerators for heavy ion tandems serving the nuclear physics community (ATLAS, INFN-LNL)
- Increased interest in Radioactive Ion Beam (RIBs) has created a renaissance in low and medium β superconducting cavity development in the last 20 years for both post-accelerators and drivers – ISAC-II, NSCL-ReA, FRIB
- High duty cycle driver linacs are now being built with superconducting sections beginning at lower β values (SPIRAL-II, C-ADS, IFMIF, ESS)
 - Rise in performance of spoke cavities and half-wave resonators (HWR)
 - Shapes are being optimized for performance with more emphasis on forming
 - Clean room assembly, high pressure water rising and separated vacuum cryostats are now standard
- Spoke resonators are now being investigated at velocities at or near $\beta=1$ for compact machines
- Deflecting/crabbing cavities have seen a rise in interest due to high performance, compact designs for applications such as crabbing cavities for LHC high luminosity upgrade

Electromagnetic Fields

- Resonance cavity mode types: TM type, TE type, TEM type
- For a cylindrical geometry (Simplest form of a resonant cavity)



$$\vec{E}(x, y, z, t) = \vec{E}(x, y)e^{j(kz - \omega t)}$$

$$\vec{H}(x, y, z, t) = \vec{H}(x, y)e^{j(kz - \omega t)}$$

TM Modes

- Modes with longitudinal electric fields and no transverse magnetic fields

$$\text{TM Modes: } \begin{cases} E_z = E_0 \cos\left(\frac{p\pi z}{L}\right) J_m\left(\frac{x_{mn}r}{R}\right) \cos(m\phi), \\ E_r = -E_0 \frac{p\pi R}{Lx_{mn}} \sin\left(\frac{p\pi z}{L}\right) J'_m\left(\frac{x_{mn}r}{R}\right) \cos(m\phi), \\ E_\phi = E_0 \frac{mp\pi R^2}{rLx_{mn}^2} \sin\left(\frac{p\pi z}{L}\right) J_m\left(\frac{x_{mn}r}{R}\right) \sin(m\phi), \\ H_z = 0, \\ H_r = jE_0 \frac{m\omega R^2}{c\eta r x_{mn}^2} \cos\left(\frac{p\pi z}{L}\right) J_m\left(\frac{x_{mn}r}{R}\right) \sin(m\phi), \\ H_\phi = jE_0 \frac{\omega R}{c\eta x_{mn}} \cos\left(\frac{p\pi z}{L}\right) J'_m\left(\frac{x_{mn}r}{R}\right) \cos(m\phi), \end{cases}$$

$$\omega_{TM_{mnp}} = c\sqrt{\left(\frac{x_{mn}c}{R}\right)^2 + \left(\frac{p\pi}{L}\right)^2}$$

x_{mn} is the n^{th} root of J_m

TE Modes

- Modes with longitudinal magnetic fields and no transverse electric fields

$$\text{TE Modes: } \begin{cases} H_z = H_0 \sin\left(\frac{p\pi z}{L}\right) J_m\left(\frac{x'_{mn}r}{R}\right) \cos(m\phi), \\ H_r = H_0 \frac{p\pi R}{Lx'_{mn}} \cos\left(\frac{p\pi z}{L}\right) J'_m\left(\frac{x'_{mn}r}{R}\right) \cos(m\phi), \\ H_\phi = -H_0 \frac{mp\pi R^2}{rL(x'_{mn})^2} \cos\left(\frac{p\pi z}{L}\right) J_m\left(\frac{x'_{mn}r}{R}\right) \sin(m\phi), \\ E_z = 0, \\ E_r = jH_0 \frac{m\eta\omega R^2}{cr(x'_{mn})^2} \sin\left(\frac{p\pi z}{L}\right) J_m\left(\frac{x'_{mn}r}{R}\right) \sin(m\phi), \\ E_\phi = jH_0 \frac{\eta\omega R}{cx'_{mn}} \sin\left(\frac{p\pi z}{L}\right) J'_m\left(\frac{x'_{mn}r}{R}\right) \cos(m\phi), \end{cases}$$

$$\omega_{TE_{mnp}} = c\sqrt{\left(\frac{x'_{mn}c}{R}\right)^2 + \left(\frac{p\pi}{L}\right)^2}$$

x'_{mn} is the n^{th} root of J'_m

Modes in a Pill Box Cavity

SRF 2023 Tutorial –
SRF Fundamental/Nb Material,
Akira Miyazak

- TM_{010}
 - Electric field is purely longitudinal
 - Electric and magnetic fields have no angular dependence
 - Mode of interest for acceleration in elliptical cavities
 - Frequency depends only on radius, independent of length
- TM_{0np}
 - Monopole modes that can couple to the beam and exchange energy
- TM_{1np}
 - Dipole modes that can deflect the beam
- TE modes
 - No longitudinal E field
 - Cannot couple to the beam
- TEM modes → For coaxial geometries
 - Transverse Electro Magnetic (TEM) mode

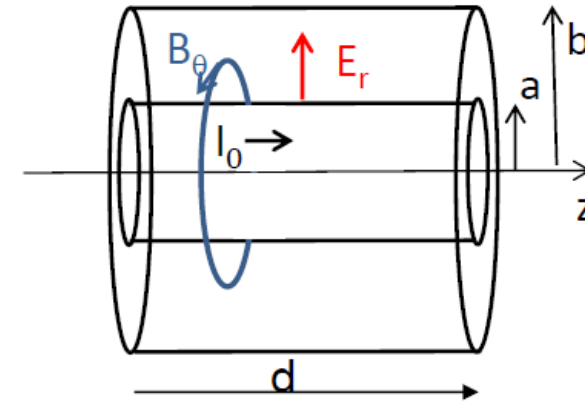
Coaxial Resonator (TEM Mode)

- Consider a coaxial geometry with grounded plates at the ends with inner radius a , outer radius b and length d
- A standing wave occurs with E_r vanishing at the end walls at $z=0$ and $z=d$
- Fields components

$$B_\theta = \frac{\mu_0 I_0}{\pi r} \cos \frac{p\pi z}{d} e^{j\omega t}$$

$$E_r = -2j \sqrt{\frac{\mu_0}{\epsilon_0}} \frac{I_0}{2\pi r} \sin \frac{p\pi z}{d} e^{j\omega t}$$

$$\text{where } \omega = k_z c = \frac{p\pi c}{d}, \quad p = 1, 2, 3, \dots$$



- Peak voltage on the inner conductor is found by integrating the radial electric field between the grounded outer conductor and the inner conductor

$$\hat{V}(z) = \int_a^b E_r(z) dr$$

$$\hat{V}(z) = \sqrt{\frac{\mu_0}{\epsilon_0}} \frac{I_0}{\pi} \ln\left(\frac{b}{a}\right) \sin \frac{p\pi z}{d}$$

Types of Superconducting Non-Elliptical Cavities

TM Type

- Accelerating cavities

Twin axis cavity
TM₁₁₀-like mode



- Deflecting and crabbing cavities

Squashed elliptical cavity
TM₁₁₀-like mode



TE Type

- Accelerating cavities

RF quadrupole cavity
TE₂₁-like mode

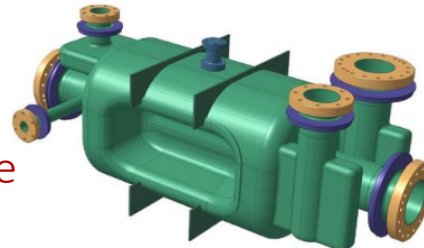


- Deflecting and crabbing cavities

Double quarter wave cavity
TE₁₁-like mode



RF-dipole cavity
TE₁₁-like mode



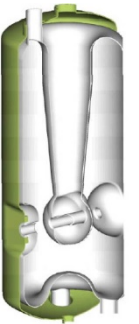
TEM Type

- Accelerating cavities

Quarter Wave Cavity



Half Wave Cavity

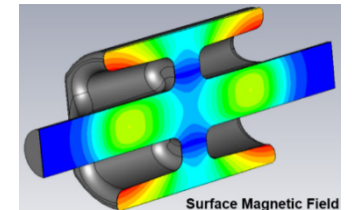


Spoke Cavity



- Deflecting and crabbing cavities

4-Rod Cavity



Designing Non-Elliptical Cavities

- No universal design for any of the cavity geometries
- Has many degrees of freedom with many parameters for optimization
- *May lead to complicated designs*

BUT IT IS MORE FUN !!

- Non-elliptical cavities are 3D geometries
- Requires 3D simulations to optimize cavity designs
- Available simulation packages
 - ✓ CST Studio, HFSS, ACE3P Code Suite, COMSOL

NON-ELLIPTICAL ACCELERATING CAVITIES

(1) BASIC PRINCIPLES

(2) TYPES OF NON-ELLIPTICAL ACCELERATING CAVITIES

- i. TEM-Type Cavities**
- ii. TM-Type Cavities**

(3) DESIGN CONSIDERATIONS

BASIC PRINCIPLES OF NON-ELLIPTICAL ACCELERATING CAVITIES

RF Acceleration

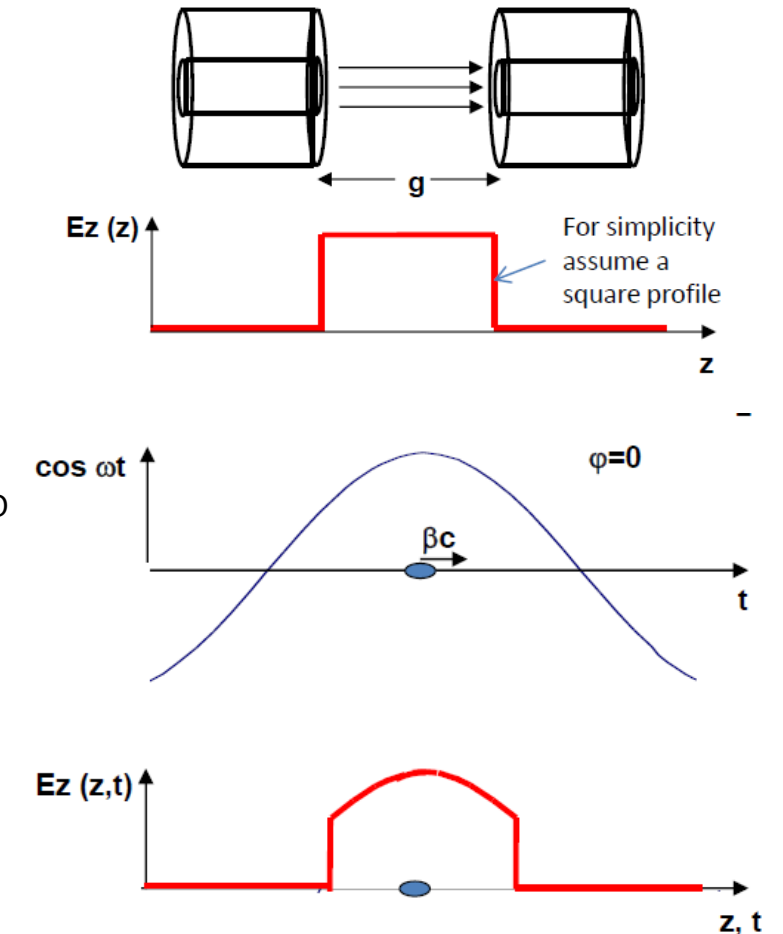
- A standing wave is established in the resonator with a time varying E_z field on axis
- For a particle travelling on axis with velocity βc and sees a field that is the product of spatial variation and time modulation

$$E_z(z, t) = E_z(\rho = 0, z) \cos(\omega t + \varphi)$$

- φ is a constant (rf phase) that defines the time of arrival of the particle with respect to the rf time modulation
- A phase corresponds to the maximum acceleration that can be given to the particle (on crest acceleration)
- One can calculate the accelerating voltage (V_{eff}) imparted to the particle by

$$V_{\text{eff}}(\varphi) = \int_{-\infty}^{\infty} E_z(\rho = 0, z) \cos(\omega t(z) + \varphi) dz$$

Example: Gap between drift tubes



Accelerating Voltage and Gradient

$$V_{eff}(\varphi) = \int_{-\infty}^{\infty} E_z(\rho=0, z) \cos(\omega t(z) + \varphi) dz$$

- Time and position are linked through the velocity (assuming β doesn't change)

$$t = \frac{z}{v} = \frac{z}{\beta c}, \quad c = f\lambda \text{ and } \omega = 2\pi f$$

$$V_{eff}(\varphi) = \int_{-\infty}^{\infty} E_z(\rho=0, z) \cos\left(\frac{2\pi z}{\beta\lambda} + \varphi\right) dz$$

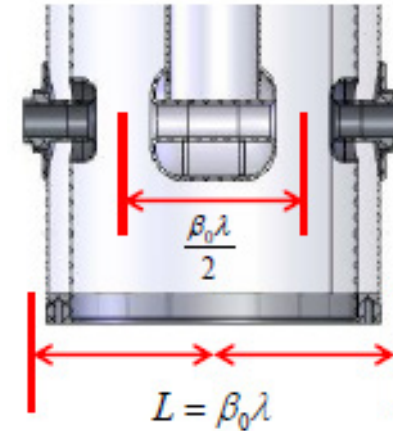
$$V_{eff}(\varphi) = \int_{-\infty}^{\infty} E_z(\rho=0, z) \left[\cos\left(\frac{2\pi z}{\beta\lambda}\right) \cos\varphi - \sin\left(\frac{2\pi z}{\beta\lambda}\right) \sin\varphi \right] dz$$

- Since E_z is typically an even function this simplifies to

$$V_{eff}(\varphi) = \int_{-\infty}^{\infty} E_z(\rho=0, z) \cos\left(\frac{2\pi z}{\beta\lambda}\right) dz \cos\varphi = V_c \cos\varphi$$

$$\text{where } V_c = \int_{-\infty}^{\infty} E_z(\rho=0, z) \cos\left(\frac{2\pi z}{\beta\lambda}\right) dz$$

is called the accelerating voltage (also $V_c = V_{acc} = V_{eff}(0)$)



- Accelerating gradient (E_a) – Gives the effective voltage gain [MV/m]

$$E_a = \frac{V_{eff}}{L} \text{ where } L = n \frac{\beta_0 \lambda}{2}, \text{ n is the number of cells}$$

$$\text{where } V_{eff} = \int_{-\infty}^{\infty} E_z(z, t) dz \text{ with } \beta = \beta_0 \text{ and } \varphi = 0$$

- Very important to specify length in defining E_a

Energy Gain, Transit Time Factor, Velocity Acceptance

- Energy gain:
$$\Delta W = q \int_{-\infty}^{+\infty} E(z) \cos(\omega t + \phi) dz$$

- Velocity acceptance:
$$\Theta = \frac{\int_{-\infty}^{\infty} E(z) \cos\left(\frac{\omega z}{\beta c}\right) dz}{\int_{-\infty}^{\infty} E(z) \cos\left(\frac{\omega z}{\beta_0 c}\right) dz}$$

- Velocity acceptance is the reduction of the energy gain associated with the time dependence of the field
- Transit time factor: Time variation of the field during particle transit through the gap

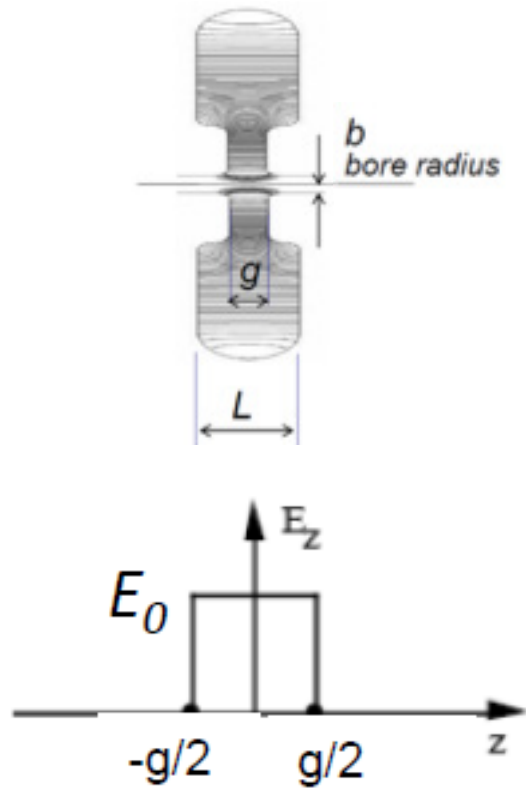
$$T(\beta) = \frac{\int_{-\infty}^{\infty} E(z) \cos\left(\frac{\omega z}{\beta_0 c}\right) dz}{\int_{-\infty}^{\infty} E(z) dz}$$

} Achievable voltage at $\beta = \beta_0$
} Maximum voltage at $\varphi = 0$

- Transit time factor is the additional reduction of energy gain for particles whose velocity is different from optimal velocity (β_0)

Single Gap Structure

- For an accelerating gap with an accelerating field approximated with a square profile
- Here note that as $g \rightarrow 0$ the $T \rightarrow 1$, but small gaps cannot support high fields, so we optimize the gap geometry using a number of considerations



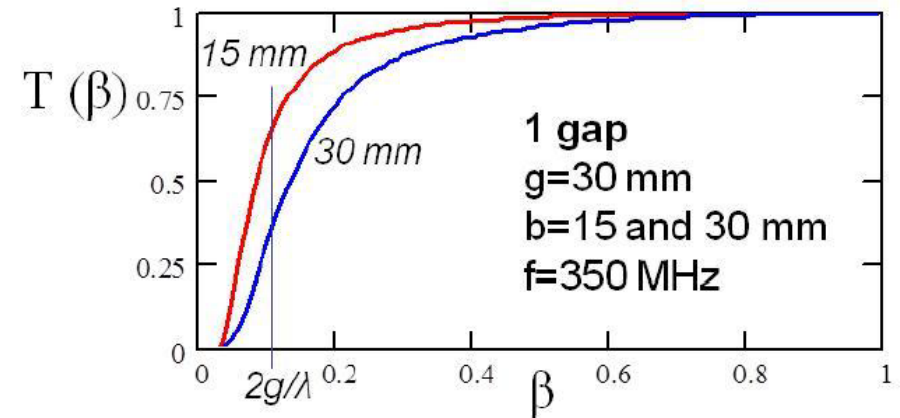
$$E_z(\rho = 0, z) = E_0 \Big|_{-g/2}^{g/2} \text{ else } E_z(0, z) = 0 \text{ so } V_0 = E_0 L = E_0 g$$

$$V_c = \int_{-g/2}^{g/2} E_0 \cos\left(\frac{2\pi z}{\beta\lambda}\right) dz = E_0 \frac{\beta\lambda}{\pi} \sin\left(\frac{\pi g}{\beta\lambda}\right) \quad T = \frac{V_{eff}}{V_0} = \frac{\beta\lambda}{\pi g} \sin\left(\frac{\pi g}{\beta\lambda}\right) \text{ and } V_{eff} = E_0 T L$$

$$T(\beta) \cong \frac{\sin\left(\frac{\pi g}{\beta\lambda}\right)}{\frac{\pi g}{\beta\lambda}}$$

Aperture b contributes to the effective gap length:

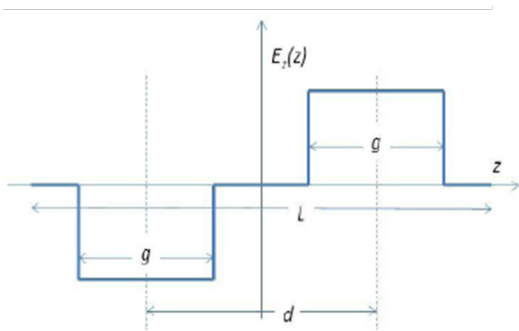
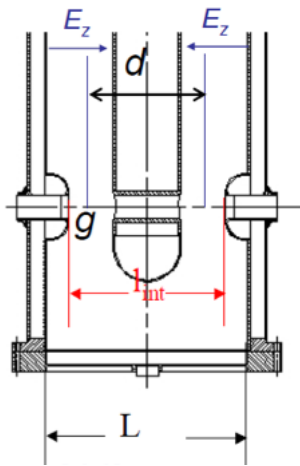
$$g_{eff} \approx \sqrt{g^2 + (2b)^2}$$



Rule of thumb: $g_{eff} < \beta\lambda/2 \rightarrow T(\beta) > 0.63$

Two Gap Structure

- For a two gap with an accelerating field approximated with a square profile in the π mode
- Slower or faster particles will get less acceleration due to poor synchronization with the rf phase

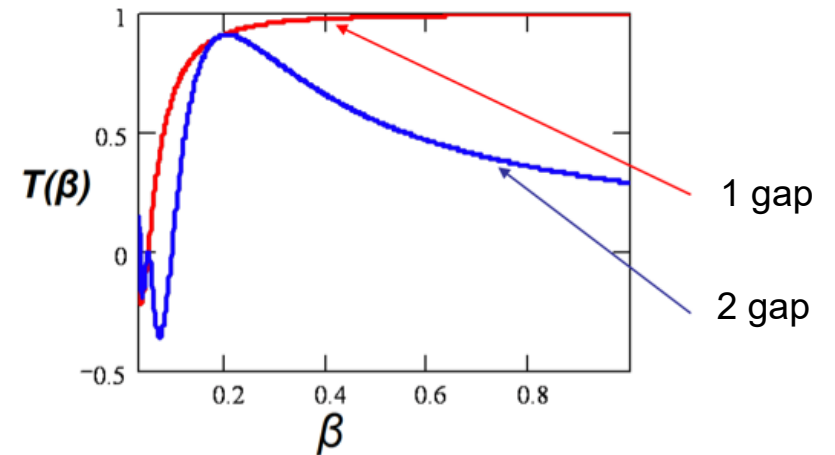


$$T_{2gap}(\beta, g) = \frac{\int_{cav} E(0, z) \sin\left(\frac{2\pi z}{\beta\lambda}\right) dz}{\int_{cav} |E(0, z) dz|} = \frac{\sin\left(\frac{\pi g}{\beta\lambda}\right)}{\frac{\pi g}{\beta\lambda}} \sin\left(\frac{\pi\beta_0}{2\beta}\right) \text{ where } \beta_0 = \frac{2L}{\lambda}$$

$$\Delta W = qV_0 T \cos\varphi = qV_{eff} \frac{T(\beta)}{T(\beta_0)} \cos\varphi$$

$$T(\beta) \cong \underbrace{\frac{\sin\left(\frac{\pi g}{\beta\lambda}\right)}{\frac{\pi g}{\beta\lambda}}}_{\text{1 gap term}} \underbrace{\sin\left(\frac{\pi d}{\beta\lambda}\right)}_{\text{2 gap term}}$$

high if $g < \beta\lambda/2$ high if $d \sim \beta\lambda/2$



Only 2nd term changes for more than 2 equal gaps in π mode

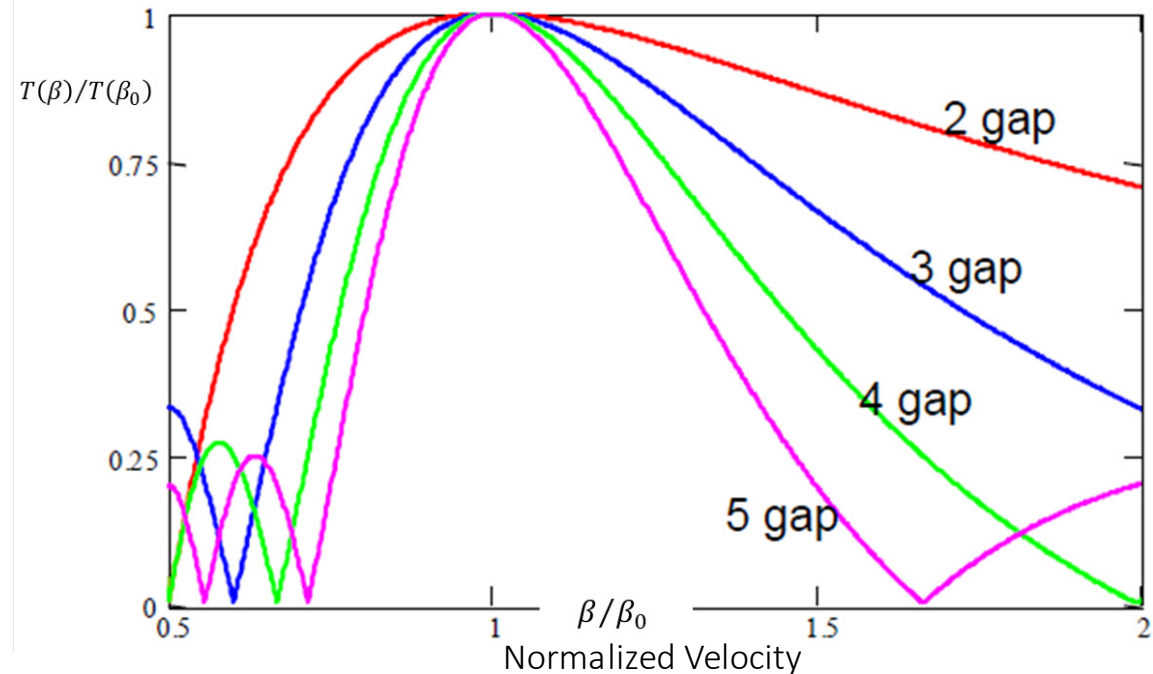
Multi Gap Structures

- Larger the number of gaps (n) larger the energy gain at a given gap voltage V_g
- But larger the number of gaps (n), narrower the velocity acceptance
 - Constant β calls for large n
 - Fast varying β calls for small n
- Higher number of gaps will provide more energy gain over a smaller velocity range

$$V_{\text{eff}} = E_a * L$$

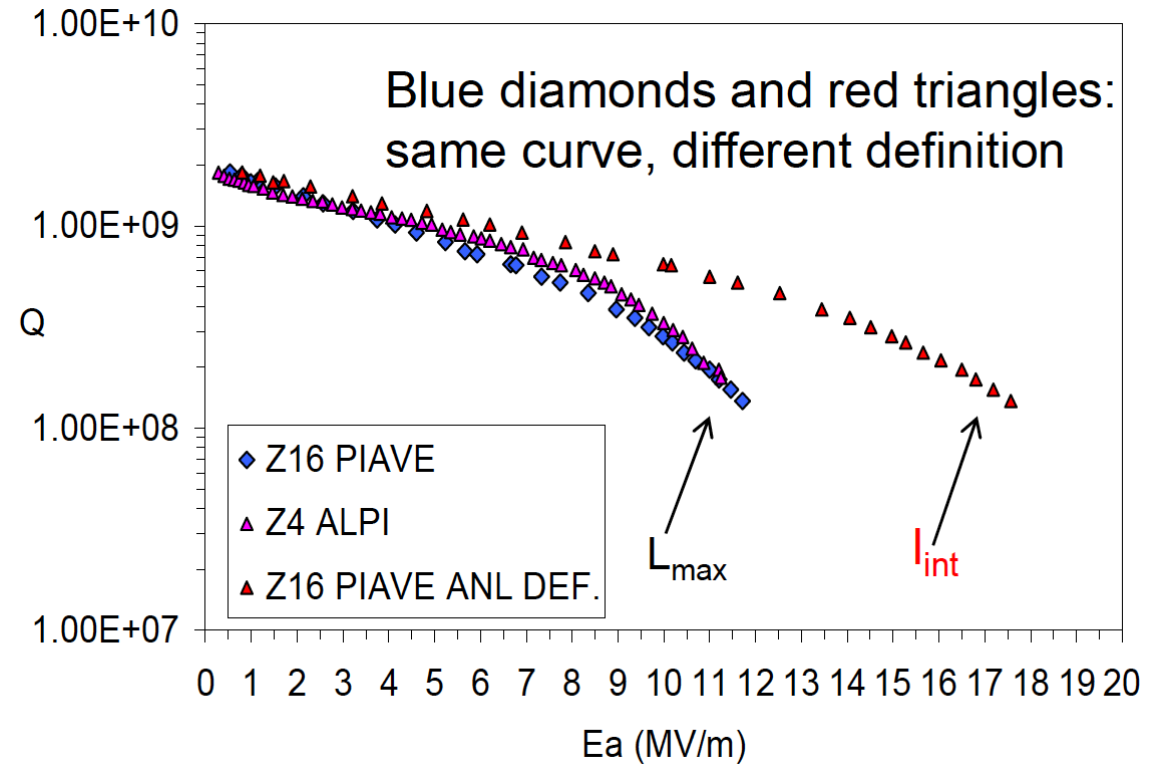
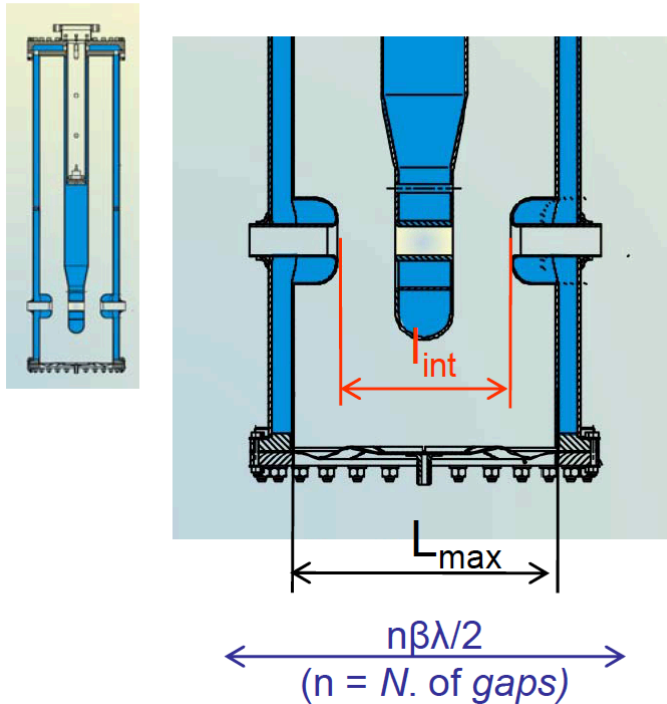
Pay attention to definition of L

Normalized
Transit Time
Factor



Different Definitions of Accelerating Gradient

- Sometimes it is difficult to decide on the definition of L : l_{int} , L_{max} , $n\beta\lambda/2$
- Shorter L is defined, larger E_a appears in Q vs E_a curves
- However, energy gain is always the same and all definitions are constant



Important Parameters

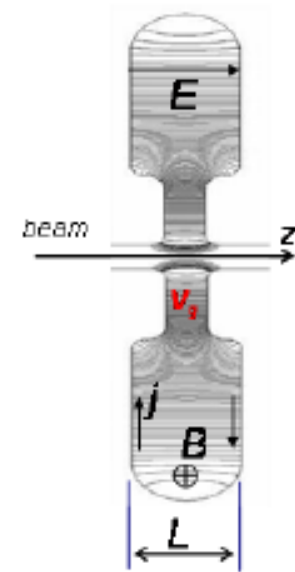
- Parameters are the same for elliptical cavities
- Important to specify the cavity reference length in defining E_a

Avg. accelerating field	$E_a = V_g T(\beta_0) / L$	MV/m	} constants
Stored energy	U / E_a^2	J/(MV/m) ²	
Shunt impedance per meter	$R_{sh} = E_a^2 L / P$	MΩ/m	
Quality Factor	$Q = \omega U / P$		
Geometrical factor	$\Gamma = Q R_s$	Ω	
Peak electric field	E_p / E_a		
Peak magnetic field	B_p / E_a	mT/(MV/m)	
Optimum β	β_0		
Cavity length	L	m	

where:

R_s = surface resistance of the cavity walls

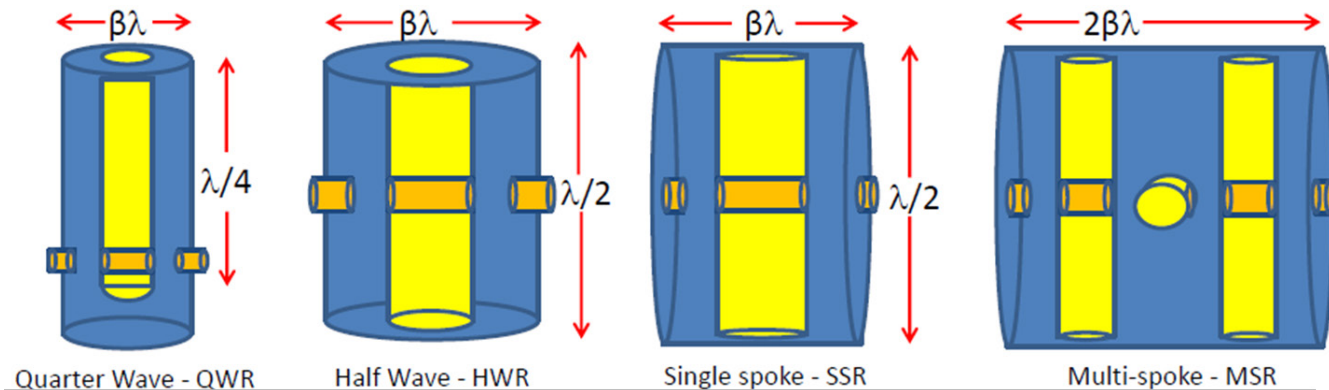
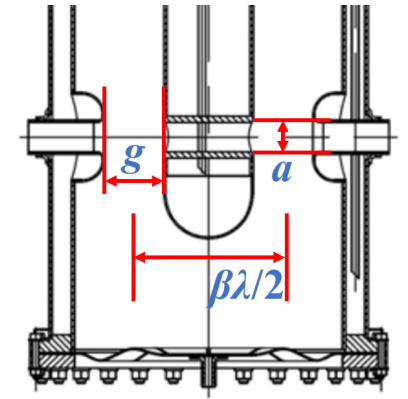
P = rf power losses in the cavity, proportional to R_s



TYPES OF NON-ELLIPTICAL ACCELERATING CAVITIES

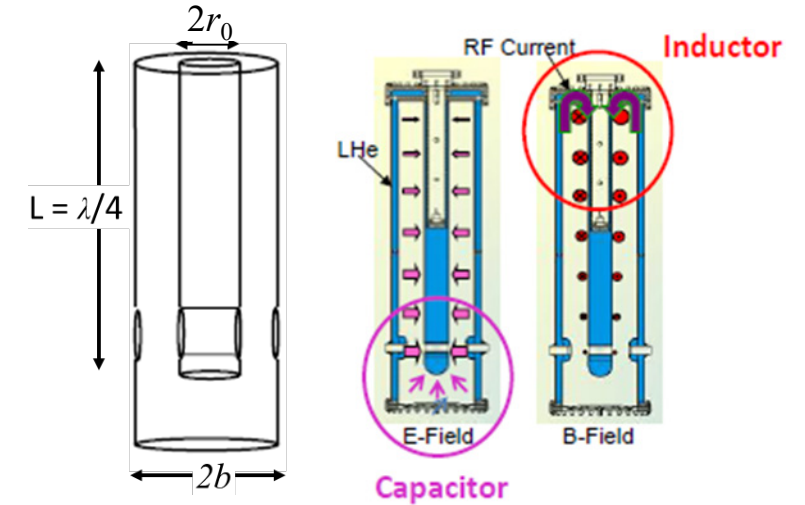
TEM-Type Cavities

- Transverse Electro Magnetic (TEM) mode cavities
 - Mode is related to the cavity symmetry axis
 - Produce accelerating voltages across the coaxial gap with variable gap distance and with transverse dimensions ~ 2 -4 times smaller than an elliptical cavities for the same frequency
- Acceleration typically uses π mode with rf phase advance of 180 deg
 - Requires distance of $\beta\lambda/2$ between gaps for synchronism
- Good transverse acceptance requires a large aperture - efficient rf acceleration requires a gap (g) to aperture (a) ratio $g/a > 1$ and a gap size $\sim 50\%$ of the cell length
 - Low velocities require low frequencies with large wavelengths
 - Low frequency cavities have large accelerating gaps



Quarter Wave Resonator (QWR)

- QWR → A capacitively loaded $\lambda/4$ transmission line
- Maximum voltage builds up on the open tip and maximum current at the plate connecting the center conductor
- Beam tube is placed near the end of the tip to produce a high voltage double gap acceleration geometry



Capacitance per unit length

Inductance per unit length

$$C = \frac{2\pi\epsilon_0}{\ln\left(\frac{b}{r_0}\right)} = \frac{2\pi\epsilon_0}{\ln\left(\frac{1}{\rho_0}\right)}$$

$$L = \frac{\mu_0}{2\pi} \ln\left(\frac{b}{r_0}\right) = \frac{\mu_0}{2\pi} \ln\left(\frac{1}{\rho_0}\right) \quad \rho_0 = \frac{r_0}{b}$$

Center conductor voltage

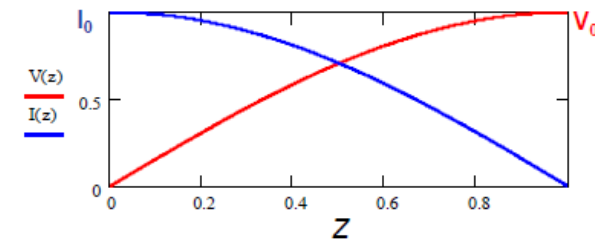
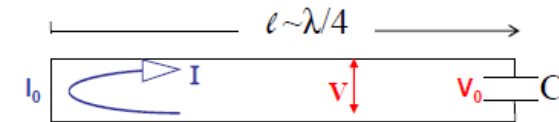
Center conductor current

$$V(z) = V_0 \sin\left(\frac{2\pi}{\lambda} z\right)$$

$$I(z) = I_0 \cos\left(\frac{2\pi}{\lambda} z\right)$$

Line impedance

$$Z_0 = \frac{V_0}{I_0} = \frac{\eta}{2\pi} \ln\left(\frac{1}{\rho_0}\right), \quad \eta = \sqrt{\frac{\mu_0}{\epsilon_0}} \approx 377\Omega$$



Quarter Wave Resonator (QWR)

Optimizing the expected performance of a resonator for given frequency and β

V_p – voltage on the center conductor with outer conductor at ground

Peak Magnetic Field

$$\frac{V_p}{b} = \begin{cases} \eta & H \\ c & B \\ 300 & B \end{cases} \rho_0 \ln\left(\frac{1}{\rho_0}\right) \sin\left(\frac{\pi}{2}\zeta\right) \quad \begin{cases} \text{m, A/m} \\ \text{m, T} \\ \text{cm, G} \end{cases}$$

V_p : Voltage across loading capacitance

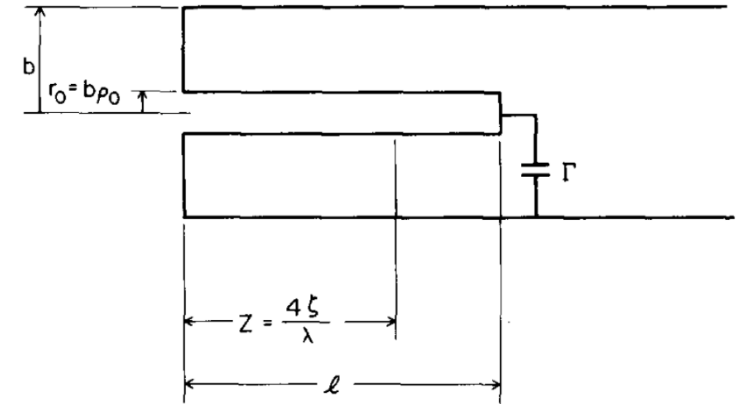
$B \approx 9$ mT at 1 MV/m

Geometrical Factor

$$G = QR_s = 2\pi\eta \frac{b \ln(1/\rho_0)}{\lambda (1+1/\rho_0)}$$

$$G \propto \eta \beta$$

R_s – surface resistance



Energy Content

$$U = V_p^2 \frac{\pi\epsilon_0}{8} \lambda \frac{1}{\ln(1/\rho_0)} \frac{\zeta + \frac{1}{\pi} \sin \pi\zeta}{\sin^2 \frac{\pi}{2}\zeta}$$

$$U \propto \epsilon_0 E^2 \beta^2 \lambda^3$$

Power Dissipation (Ignore losses in the shorting end plate)

$$P = V_p^2 \frac{8}{\pi} \frac{R_s}{\eta^2} \frac{\lambda}{b} \frac{1+1/\rho_0}{\ln^2 \rho_0} \frac{\zeta + \frac{1}{\pi} \sin \pi\zeta}{\sin^2 \frac{\pi}{2}\zeta}$$

$$P \propto \frac{R_s}{\eta^2} E^2 \beta \lambda^2$$

Shunt Impedance $(4V_p^2/P)$

$$R_{sh} = \frac{\eta^2}{R_s} \frac{32}{\pi} \frac{b}{\lambda} \frac{\ln^2 \rho_0}{1+1/\rho_0} \frac{\sin^2 \frac{\pi}{2}\zeta}{\zeta + \frac{1}{\pi} \sin \pi\zeta}$$

$$R_{sh} R_s \propto \eta^2 \beta$$

Other notations R_{sh}, R_a

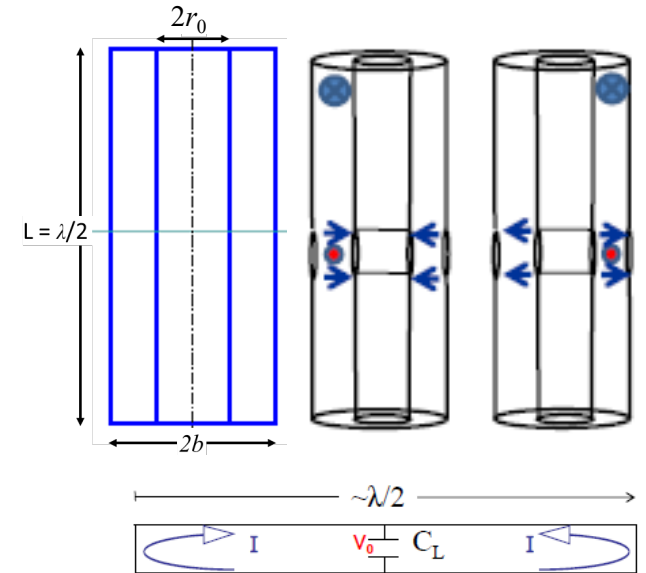
R/Q

$$\frac{R_{sh}}{Q} = \frac{16}{\pi^2} \eta \ln(1/\rho_0) \frac{\sin^2 \frac{\pi}{2}\zeta}{\zeta + \frac{1}{\pi} \sin \pi\zeta}$$

$$\frac{R_{sh}}{Q} \propto \eta$$

Half Wave Resonator (HWR)

- HWR → A $\lambda/2$ transmission line
- Equivalent to 2 QWR facing each other and connected
- Magnetic field loop around the inner conductor with peak fields at the shorted ends
- Beam tube is placed at the center of the inner conductor to coincide with the maximum voltage
- Same accelerating voltage is obtained at about 2 times larger power in QWR ($P_{\text{HWR}} \sim 2P_{\text{QWR}}$)



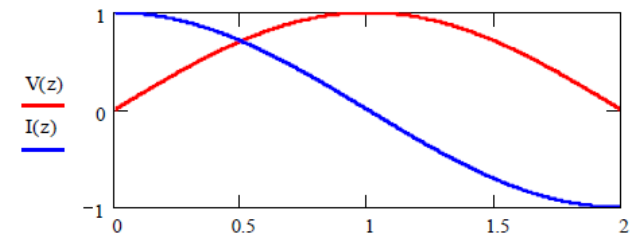
Capacitance per unit length

$$C = \frac{2\pi\epsilon_0}{\ln\left(\frac{b}{a}\right)} = \frac{2\pi\epsilon_0}{\ln\left(\frac{1}{\rho_0}\right)}$$

Inductance per unit length

$$L = \frac{\mu_0}{2\pi} \ln\left(\frac{b}{r_0}\right) = \frac{\mu_0}{2\pi} \ln\left(\frac{1}{\rho_0}\right)$$

$$\rho_0 = \frac{r_0}{b}$$



Center conductor voltage

$$V(z) = V_0 \sin\left(\frac{2\pi}{\lambda} z\right)$$

Center conductor current

$$I(z) = I_0 \cos\left(\frac{2\pi}{\lambda} z\right)$$

Line impedance

$$Z_0 = \frac{V_0}{I_0} = \frac{\eta}{2\pi} \ln\left(\frac{1}{\rho_0}\right), \quad \eta = \sqrt{\frac{\mu_0}{\epsilon_0}} \approx 377\Omega$$

Half Wave Resonator (HWR)

Optimizing the expected performance of a resonator for given frequency and β

V_p – voltage on the center conductor with outer conductor at ground

Peak Magnetic Field

$$\frac{V_p}{b} = \begin{Bmatrix} \eta & H \\ c & B \\ 300 & B \end{Bmatrix} \rho_0 \ln\left(\frac{1}{\rho_0}\right) \quad \begin{Bmatrix} \text{m, A/m} \\ \text{m, T} \\ \text{cm, G} \end{Bmatrix}$$

V_p : Voltage across loading capacitance

$B = 9 \text{ mT}$ at 1 MV/m

Power Dissipation (Ignore losses in the shorting end plates)

$$P = V_p^2 \frac{16}{\pi} \frac{R_s}{\eta^2} \frac{\lambda}{b} \frac{1+1/\rho_0}{\ln^2 \rho_0}$$

$$P \propto \frac{R_s}{\eta^2} E^2 \beta \lambda^2$$

Geometrical Factor

$$G = QR_s = 2\pi \eta \frac{b}{\lambda} \frac{\ln(1/\rho_0)}{1+1/\rho_0}$$

$$G \propto \eta \beta$$

R_s – surface resistance

Shunt Impedance

$$R_{sh} = \frac{\eta^2}{R_s} \frac{16}{\pi} \frac{b}{\lambda} \frac{\ln^2 \rho_0}{1+1/\rho_0}$$

$$R_{sh} R_s \propto \eta^2 \beta$$

Energy Content

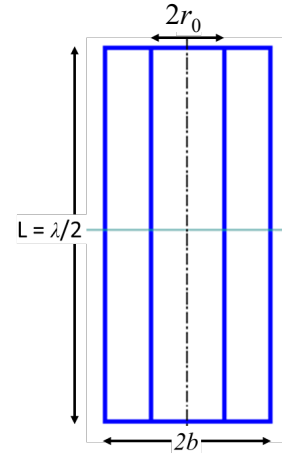
$$U = V_p^2 \frac{\pi \epsilon_0}{4} \lambda \frac{1}{\ln(1/\rho_0)}$$

$$U \propto \epsilon_0 E^2 \beta^2 \lambda^3$$

R/Q

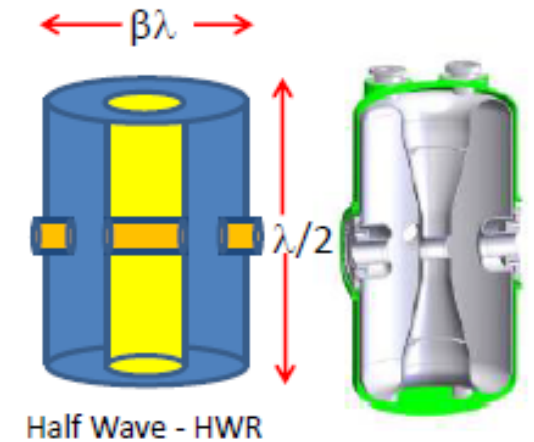
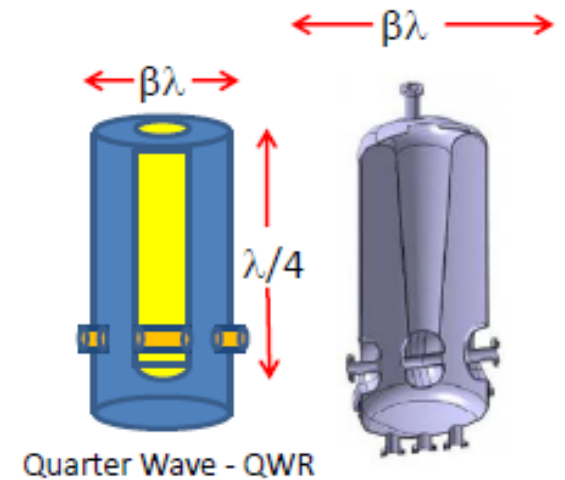
$$\frac{R_{sh}}{Q} = \frac{8}{\pi^2} \eta \ln(1/\rho_0)$$

$$\frac{R_{sh}}{Q} \propto \eta$$



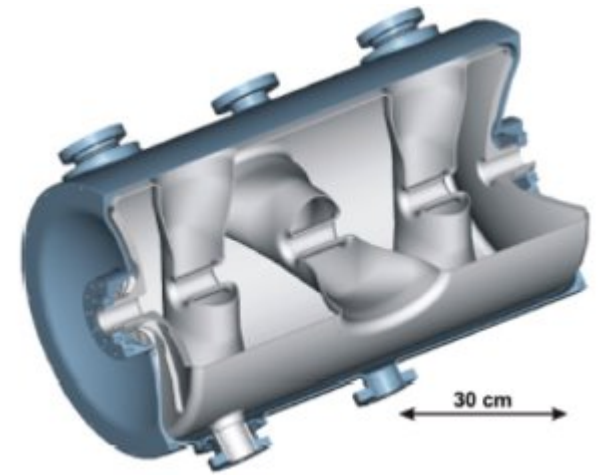
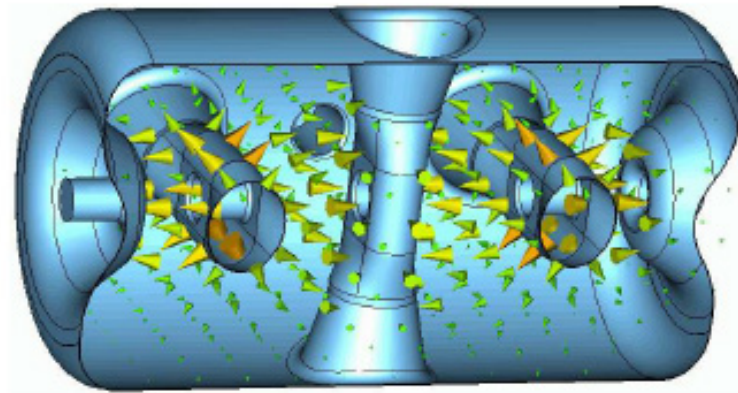
QWR vs HWR

- QWR is the choice for low applications where a low frequency is needed
 - Requires ~50% less structure compared to HWR for same frequency
 - RF power loss is ~50% of HWR for same frequency and β_0
 - Allows low frequency cavities with larger voltage acceptance ($R_{sh}/Q_{QWR} = 2(R_{sh}/Q_{HWR})$)
 - Asymmetric field pattern introduces vertical steering especially for light ions that increase with velocity (Avoid using for $\beta_0 > 0.2$)
 - Mechanically less stable than HWR due to unsupported end
- HWR is chosen for mid velocity range ($\beta_0 > 0.2$) or where steering must be eliminated (ie. High intensity light ion applications)
 - Produces 2× more rf losses for the same frequency and β_0
 - 2× longer for the same frequency
 - Symmetric field pattern and increased mechanical rigidity



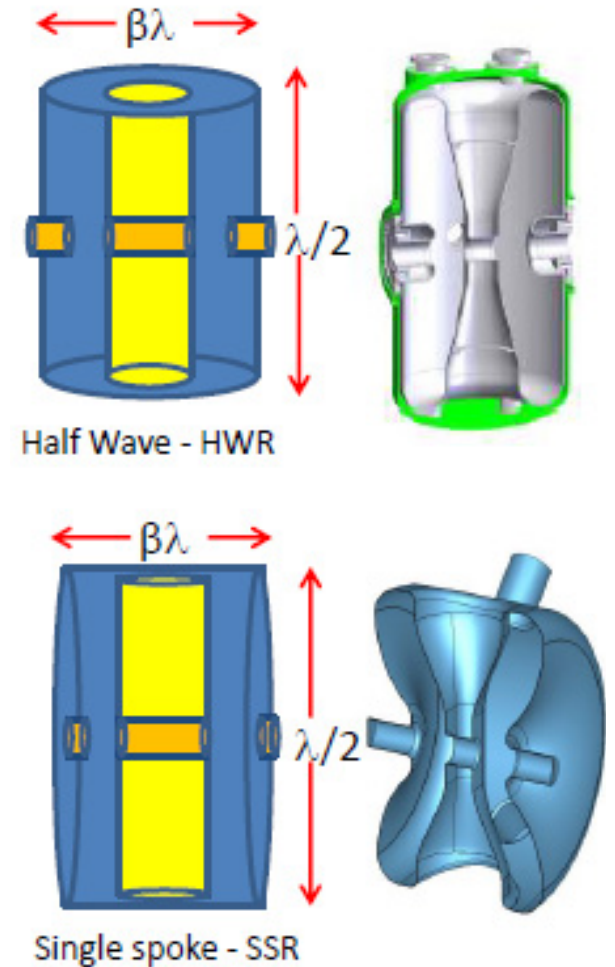
Single and Multi Spoke Cavities

- Supposed to cover ranges $\beta=0.1-0.6$ with $f=300-900$ MHz
- Spoke cavities are also designed at high β and $\beta =1$
- Single spoke cavities are same as TEM-like HW cavities with respect to the spoke axis
- Single spoke geometries allow extension along the beam path to provide multipole spoke
 - In multi spoke cavity spokes are rotated 90 deg from cell to cell
 - Higher effective voltage with low velocity acceptance
 - Strong cell-to-cell coupling with cells linked by the magnetic field



HWR vs Single Spoke Cavities

- Single spoke resonator (SSR) is another variant of the half wave TEM mode cavities
- In HWR the outer conductor (with diameter $\beta_0 \lambda$) is coaxial with the inner conductor
- In SSR the outer cylinder (with diameter $\lambda/2$) is coaxial with beam pipes
 - For $\beta_0 < 0.5$ the SSR has larger overall physical envelop than the HWR for the same frequency
- Cavity choices:
 - For low β applications ($\beta_0 = 0.1-0.25$)
 - HWR is chosen for ~ 160 MHz
 - SSR is chosen for ~ 320 MHz
 - For higher β applications ($\beta_0 = 0.25-0.5$)
 - HWR and SSR are chosen for ~ 320 MHz



High β Spoke Cavities

- High velocity spoke cavities with $\beta_0 > 0.8$ are being designed as an alternative to high β elliptical cavities
- Cavity features
 - Cavities are relatively compact
 - Between 20% - 50% smaller (radially) than a TM cavity of the same frequency and β_0
 - For high β_0 cavities diameter is close to TM counterparts
 - Allows low frequency at reasonable size with high longitudinal acceptance
 - Mechanically stable
 - Allows possible 4 K operation
 - Can achieve high shunt impedance
- Possible applications
 - For pulsed spallation neutron sources
 - Compact light sources



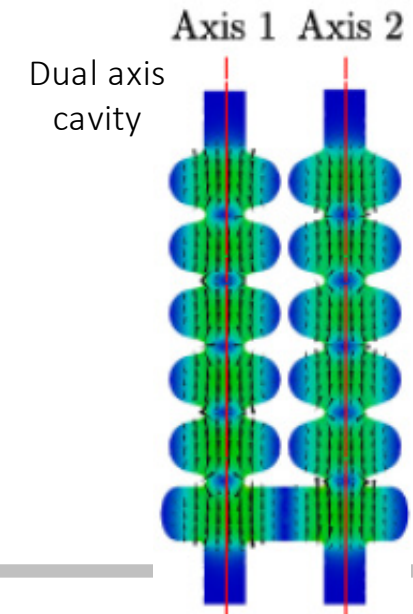
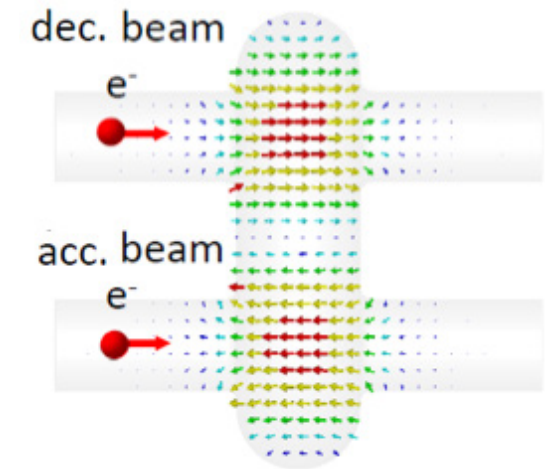
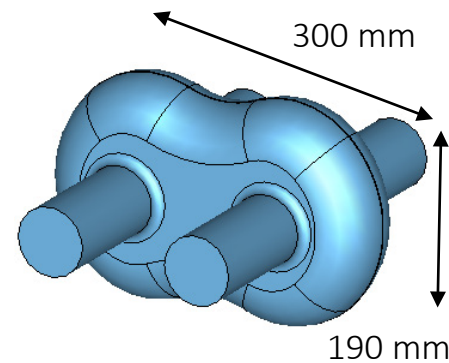
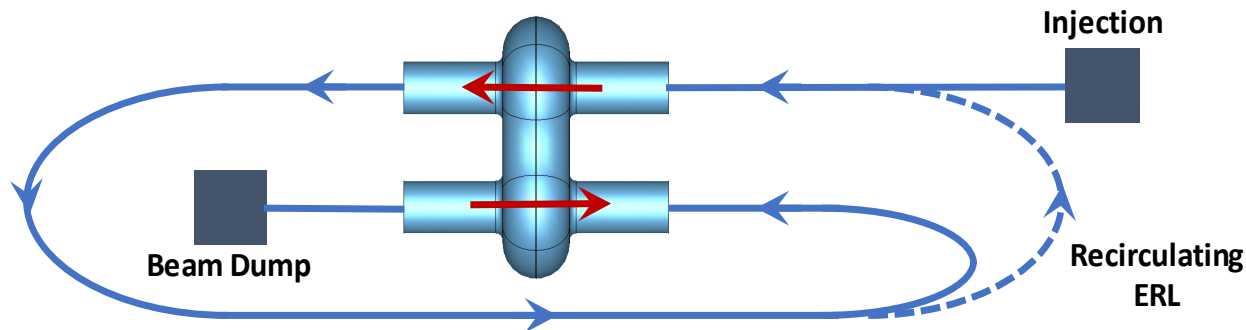
325 MHz $\beta=0.82$ Single Spoke Cavity



500 MHz $\beta=1.0$ Double Spoke Cavity

TM-Type Non-Elliptical Cavities

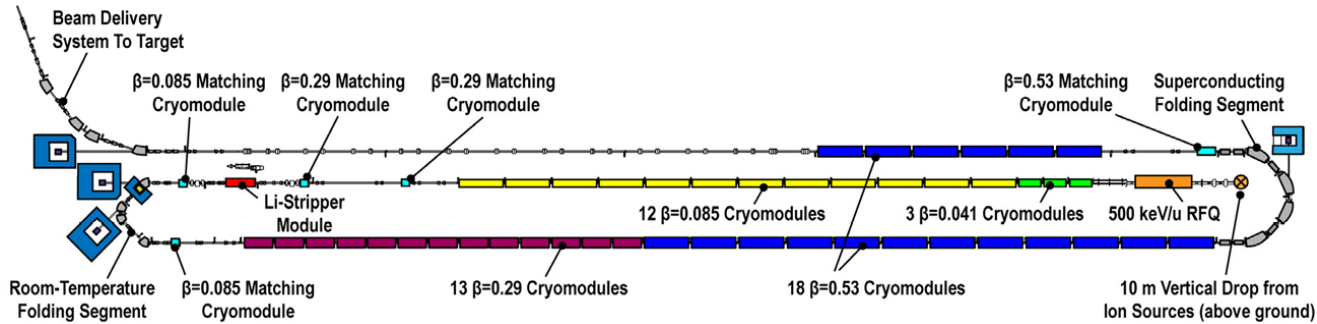
- Cavity with two beam pipes \rightarrow Twin axis cavity / dual axis cavity
- A superconducting cavity designed to accelerate and decelerate two electron beams in the same cavity
- Used in energy recovery with two separated beams traversing the cavity at the same time



REAL CAVITIES AND ACCELERATOR APPLICATIONS

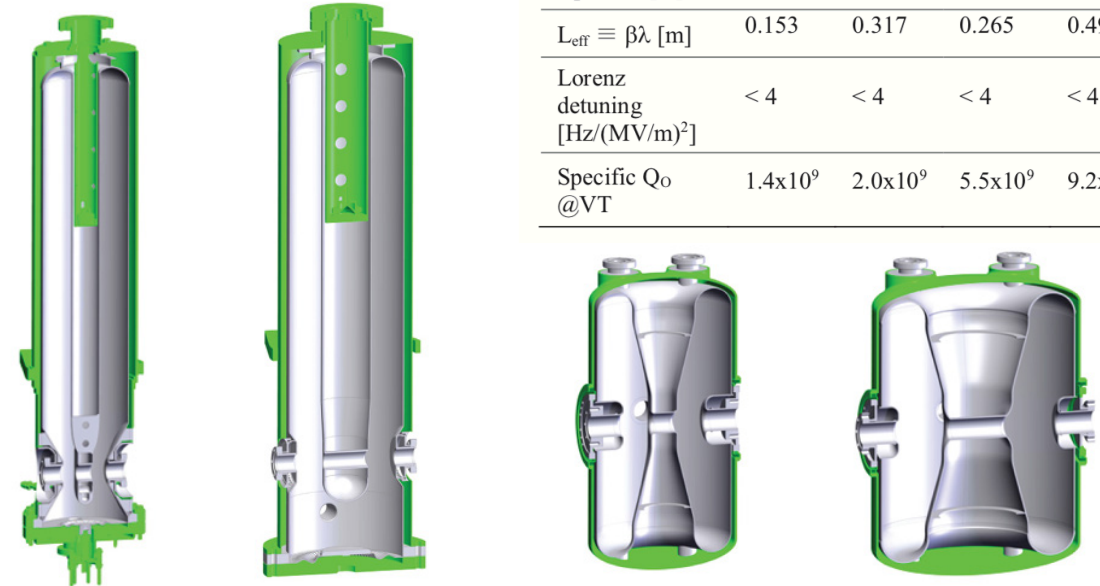
Facility of Rare Isotope Beams (FRIB)

- A high intensity, heavy ion linac
- Accelerate ion species up to ^{238}U ($A/q=7$) with energies of no less than 200 MeV/u



- No. of cavities:
 - QWR ($\beta_0=0.041$) \rightarrow 12
 - QWR ($\beta_0=0.085$) \rightarrow 88
 - HWR ($\beta_0=0.285$) \rightarrow 72
 - HWR ($\beta_0=0.53$) \rightarrow 144

Cavity Type	QWR	QWR	HWR	HWR
β_0	0.041	0.085	0.285	0.53
f [MHz]	80.5	80.5	322	322
V_a [MV]	0.810	1.80	2.09	3.70
E_{acc} [MV/m]	5.29	5.68	7.89	7.51
E_p/E_{acc}	5.82	5.89	4.22	3.53
B_p/E_{acc} [mT/(MV/m)]	10.3	12.1	7.55	8.41
R/Q [Ω]	402	455	224	230
G [Ω]	15.3	22.3	77.9	107
Aperture [m]	0.036	0.036	0.040	0.040
$L_{eff} \equiv \beta\lambda$ [m]	0.153	0.317	0.265	0.493
Lorenz detuning [Hz/(MV/m) 2]	< 4	< 4	< 4	< 4
Specific Q_0 @VT	1.4×10^9	2.0×10^9	5.5×10^9	9.2×10^9



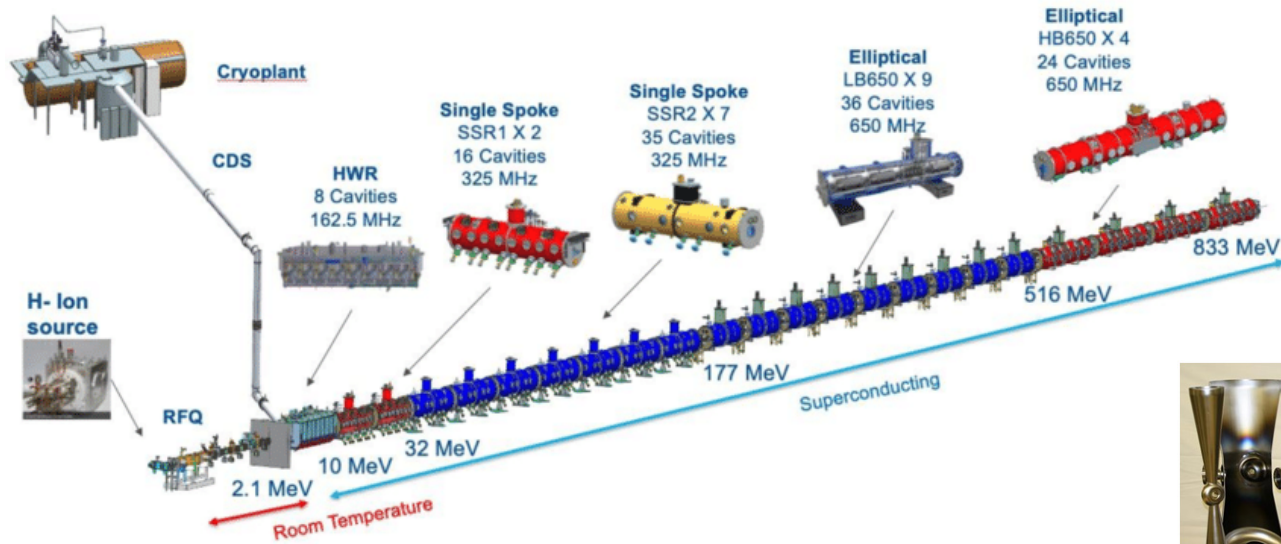
$\beta = 0.041$ cavity $\beta = 0.085$ cavity $\beta = 0.29$ cavity $\beta = 0.53$ cavity

FRIB $\beta=0.041, 0.085$ 80.5 MHz

FRIB $\beta=0.29, 0.53$ 322 MHz

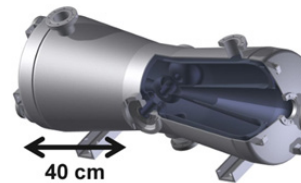
Proton Improvement Plan II (PIP II) at Fermilab

- 800 MeV linac to accelerate protons to generate intense neutrino beam for the Long Baseline Neutrino Facility (LBNF) and Deep-Underground Neutrino Experiment (DUNE)

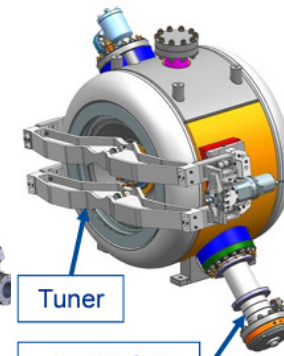


Section	Freq	Energy (MeV)	Cav/mag/CM	Type
RFQ	162.5	0.03-2.1		
HWR ($\beta_{opt}=0.11$)	162.5	2.1-10.3	8/8/1	HWR, solenoid
SSR1 ($\beta_{opt}=0.22$)	325	10.3-35	16/8/2	SSR, solenoid
SSR2 ($\beta_{opt}=0.47$)	325	35-185	35/21/7	SSR, solenoid
LB 650 ($\beta_g=0.61$)	650	185-500	33/22/11	5-cell elliptical, doublet*
HB 650 ($\beta_g=0.92$)	650	500-800	24/8/4	5-cell elliptical, doublet*

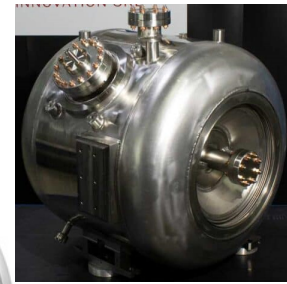
Cavity Type	HWR	SSR1	SSR2	LB650	HB650
Frequency [MHz]	162.5	325	325	650	650
Optimal β	0.112	0.22	0.475	0.61	0.97
Effective Length [cm]	20.7	20.3	43.8	70.3	106.1
Aperture [mm]	33	30	40	41.5	59
E_{peak}/E_{acc}	4.7	3.84	3.38	2.43	2.1
B_{peak}/E_{acc} [mT/(MV/m)]	5.0	5.81	5.93	4.6	3.94
G [Ω]	48	84	115	187	260
R_{sh}/Q [[Ω]]	272	242	297	327.4	576



162.5 MHz HWR



325 MHz SSR1



325 MHz SSR2

HB650 5-cell Elliptical Cavity

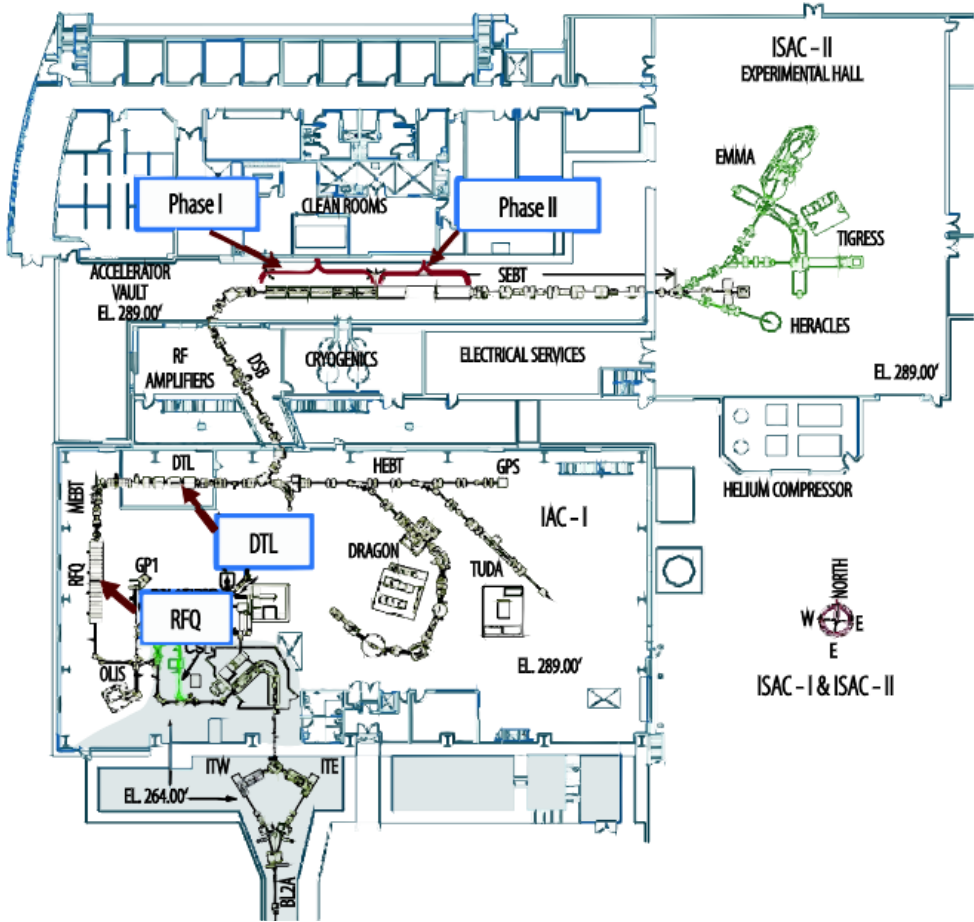


LB650 5-cell Elliptical Cavity

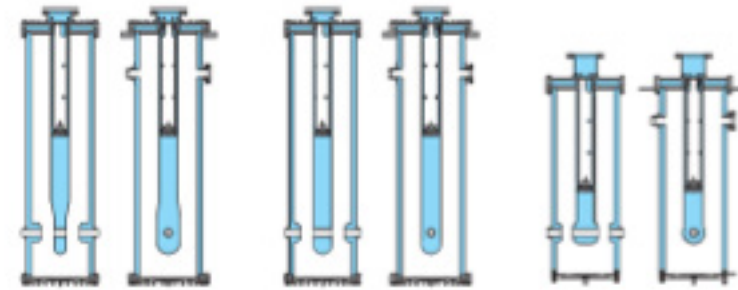
Isotope Separator and Accelerator (ISAC II) and TRIUMF

- To produce rare isotope beams

- Consists of 40 QWRs with three different cavity types of $\beta_0=5.7\%$, 7.1% and 11% with $N = 8, 12$ and 20
- All cavities have same outer and inner conductor diameter to reduce fabrication cost



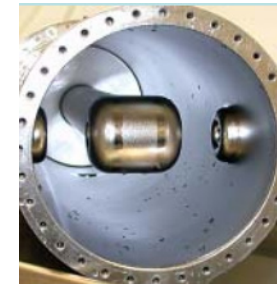
TRIUMF ISAC-II Resonators



SCB low $\beta=0.057$ 106.08 MHz SCB medium $\beta=0.071$ 106.08 MHz SCB high $\beta=0.11$ 141.44 MHz



106.08 MHz $\beta=0.057$

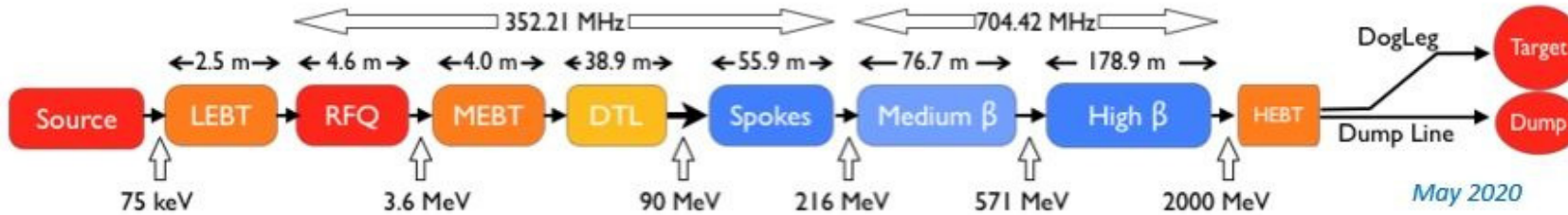


141.44 MHz $\beta=0.11$



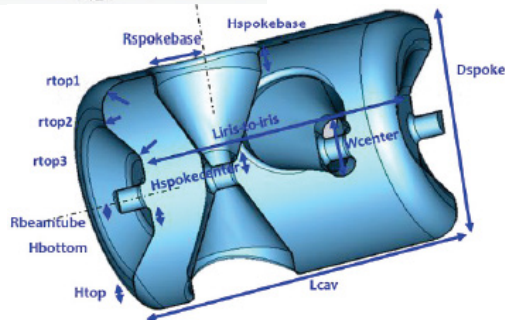
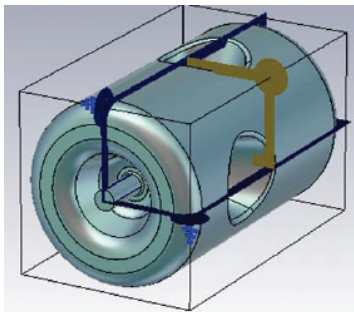
European Spallation Source (ESS)

- 5 MW long pulse proton accelerator

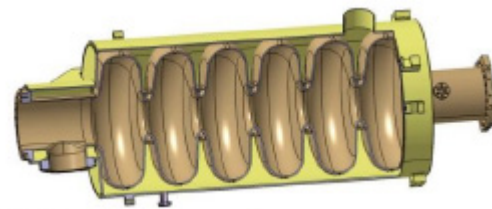


Style	Spoke	Medium- β	High- β
Freq. (MHz)	352.21	704.42	704.42
Cavity #	26	36	84
Velocity range	0.42 to 0.58	0.58 to 0.78	0.78 to 0.95
Nom. Acc. Voltage (MV)	5.74	14.3	18.2
Loaded quality factor	2.85×10^5	8×10^5	7.6×10^5
Dynamic heat load (W)	0.8	4.9	5.5
Max. forward power (kW)	335	1100	1100
RF (baseline)	Tetrode	Klystron	IOT

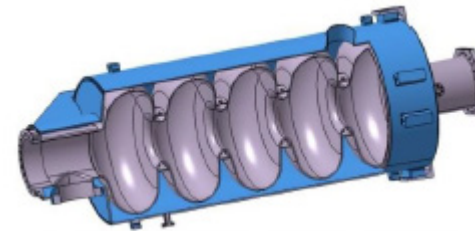
- Cavity types:
 - Spoke cavities – $\beta_0=0.5$
 - Medium β elliptical cavities – $\beta_0=0.67$
 - High β elliptical cavities – $\beta_0=0.86$



$\beta=0.5$ 352 MHz Double Spoke Cavity



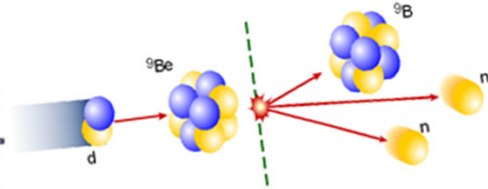
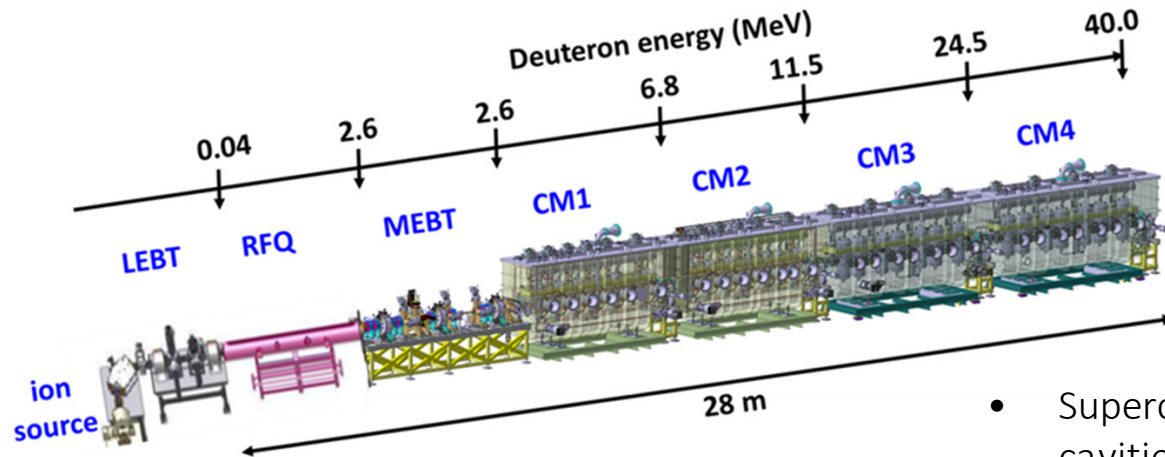
$\beta=0.67$ 704 MHz Medium- β Elliptical Cavity



$\beta=0.86$ 704 MHz High- β Elliptical Cavity

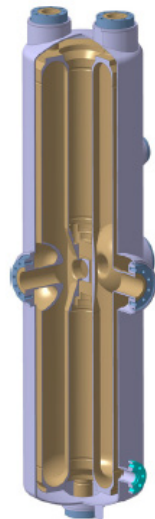
Soreq Applied Research Accelerator Facility (SARAF)

- To accelerate 5 mA beam of protons from 1.3 to 35 MeV or deuterons from 2.6 to 40 MeV

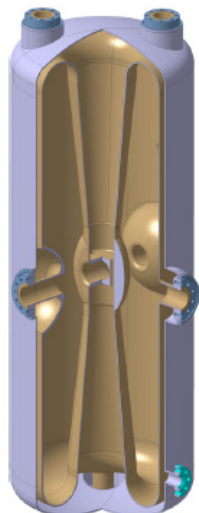


- Superconducting linac with half wave cavities
- Cryomodule layout:
 - $\beta_0=0.09$ – CM1 (6 cavities) and CM2 (7 cavities)
 - $\beta_0=0.18$ – CM3 (7 cavities) and CM4 (7 cavities)

	Low β cavity	High β cavity
β_{opt}	0.091	0.181
E_{acc}	7 MV/m	8.1 MV/m
U_{acc}	1.07 MV	2.49 MV
E_{pk}	34.5 MV/m	35.8 MV/m
B_{pk}	65.6 mT	65.3 mT
$Q_0@45 \text{ n}\Omega$	$8 \cdot 10^8$	$1.2 \cdot 10^9$
Stored energy	5.7 J	16.8 J
$P_{cryo} @45 \text{ n}\Omega$	7.9 W	15.5 W



176 MHz $\beta=0.09$



176 MHz $\beta=0.18$

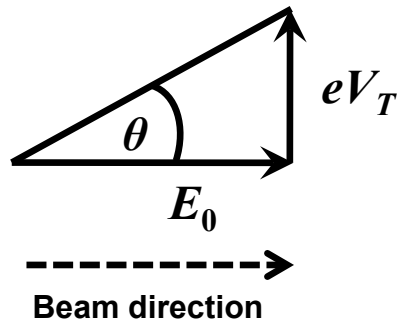
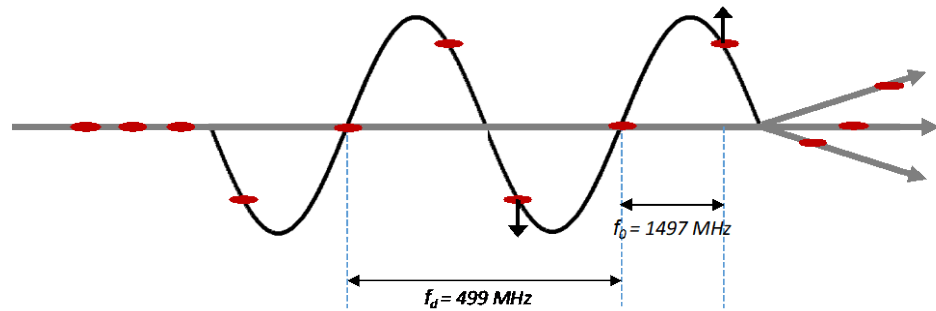
DEFLECTING/CRABBING CAVITIES

Deflecting/Crabbing Concept

- Deflecting/crabbing resonant cavities are required to generate a transverse momentum

Deflecting Cavities

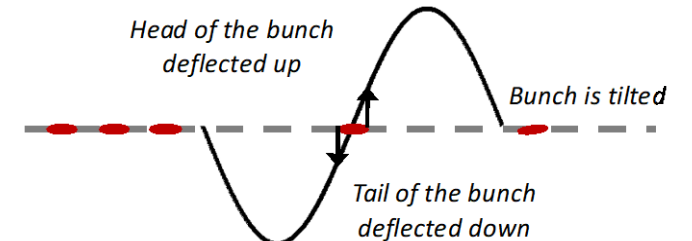
- To separate a single beam to multiple beams



$$\theta = \arctan \left[\frac{eV_t}{E_0} \right] \sim \frac{eV_t}{E_0} \quad V_t = E_0 [eV] \theta [rad]$$

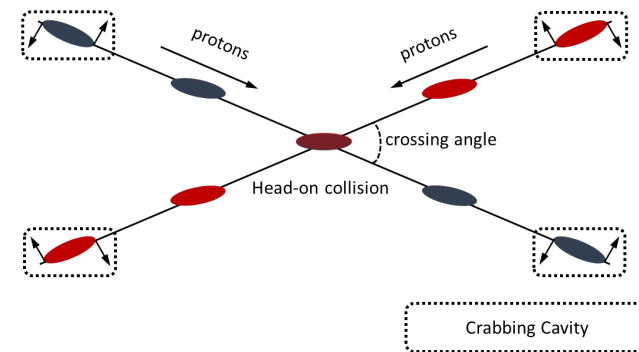
Crabbing Cavities

- To increase luminosity in colliding bunches by allowing head-on-collision of beams



$$\mathcal{L} = \frac{N_1 N_2 f_c N_b}{4\pi \sigma_x \sigma_y} F_c = \frac{N_1 N_2 f_c N_b}{4\pi \sigma_x \sigma_y} \frac{1}{\sqrt{1 + \left(\frac{\sigma_z \theta_c}{2\sigma_x} \right)^2}}$$

$$V_t = \frac{c E_0 \tan(\theta_c/2)}{\omega \sqrt{\beta_{crab} \beta^*} \sin(\psi_{cc \rightarrow ip}^x)}$$



First crabbing concept proposed by R. Palmer (1988)

Deflecting/Crabbing Cavities

- Can be produced by either or by both transverse electric (E_t) and magnetic (B_t) fields
- Lorentz force:
$$\vec{p}_t = \int_{-\infty}^{\infty} \vec{F}_t dt = \frac{q}{v} \int_{-\infty}^{+\infty} \left[\vec{E}_t + j(\vec{v} \times \vec{B}_t) \right] dz$$
- Transverse momentum is related to the gradient of the longitudinal electric field along the beam axis (Panofsky Wenzel theorem)

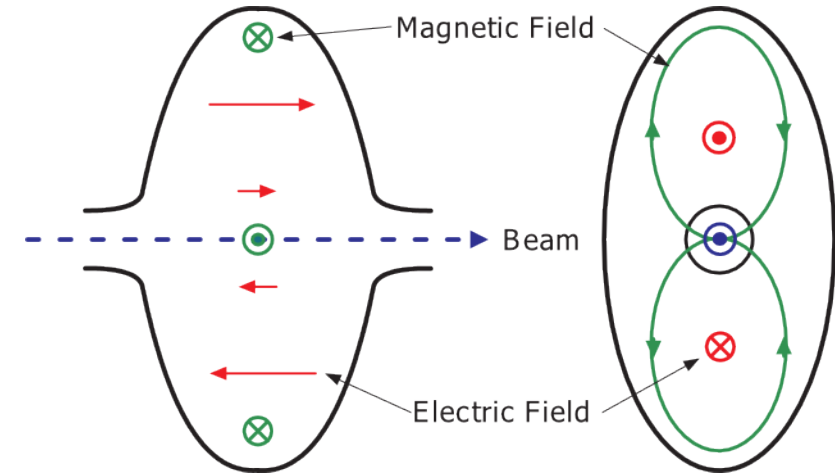
$$\vec{p}_t = -i \frac{q}{\omega} \int_{-\infty}^{+\infty} \vec{\nabla}_t E_z dz$$
$$p_t = -i \frac{q}{\omega} \lim_{r_0 \rightarrow 0} \frac{1}{r_0} \int_{-\infty}^{+\infty} [E_z(r_0, z) - E_z(0, z)] dz$$

- According to the theorem:
 - In a pure TE mode the contribution to the deflection from the magnetic field is completely cancelled by the contribution from the electric field
- Types of designs:
 - TM-type designs \rightarrow Main contribution from B_t
 - TE-like designs \rightarrow Main contribution from E_t
 - TEM-type designs \rightarrow Contribution from both E_t and B_t

TM-Type Deflecting/Crabbing Cavities

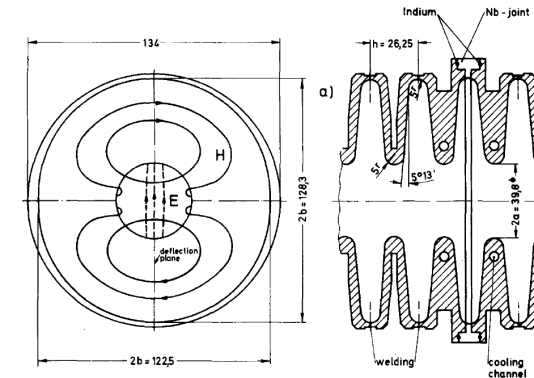
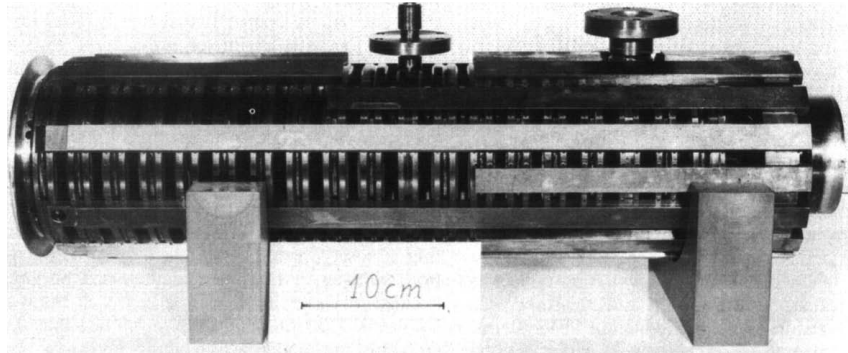
- Operates in TM_{110} -type mode
 - Lowest deflecting mode
- Squashed elliptical geometry: To separate the two polarizations of same frequency
- Contribution to the net deflection is mainly from transverse magnetic field

- Requires damping of the lowest mode (TM_{010}) in the design
- Cavity frequency is inversely proportional to transverse dimensions
- Cavity length $\sim \lambda/2$
- Large with respect to wavelength compared to new designs
 - Disadvantageous for low frequency
 - Advantageous for high frequency
 - Able to accommodate large apertures



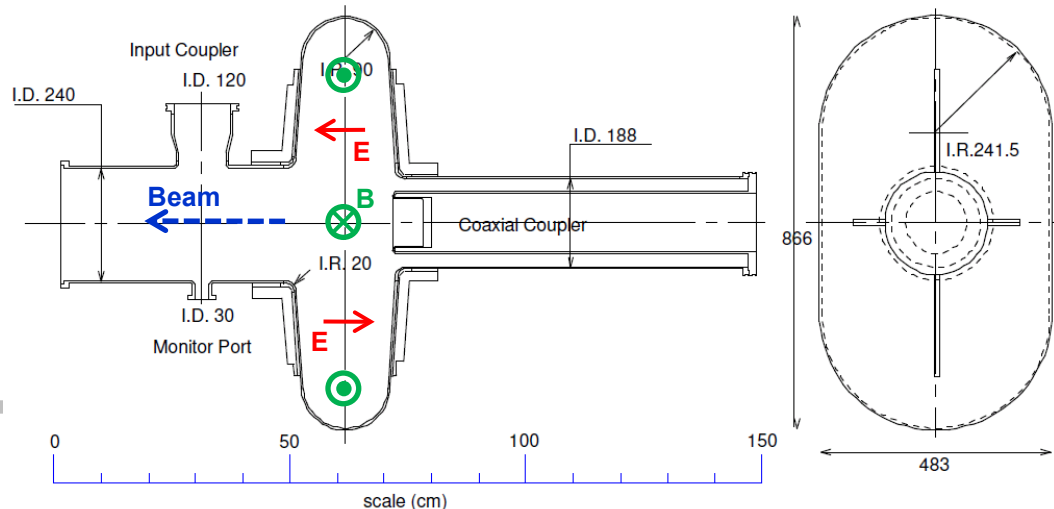
1st Deflecting/Crabbing Cavities

- 1st superconducting deflecting cavity
 - Deflecting cavity: 2.865 GHz Karlsruhe/CERN Separator (104 cells)



Designed 1970, operated 1977-1981
At IHEP since 1998

- 1st superconducting crabbing cavity
 - Crabbing cavity: 508.9 MHz cavity for SuperB Factory at KEK

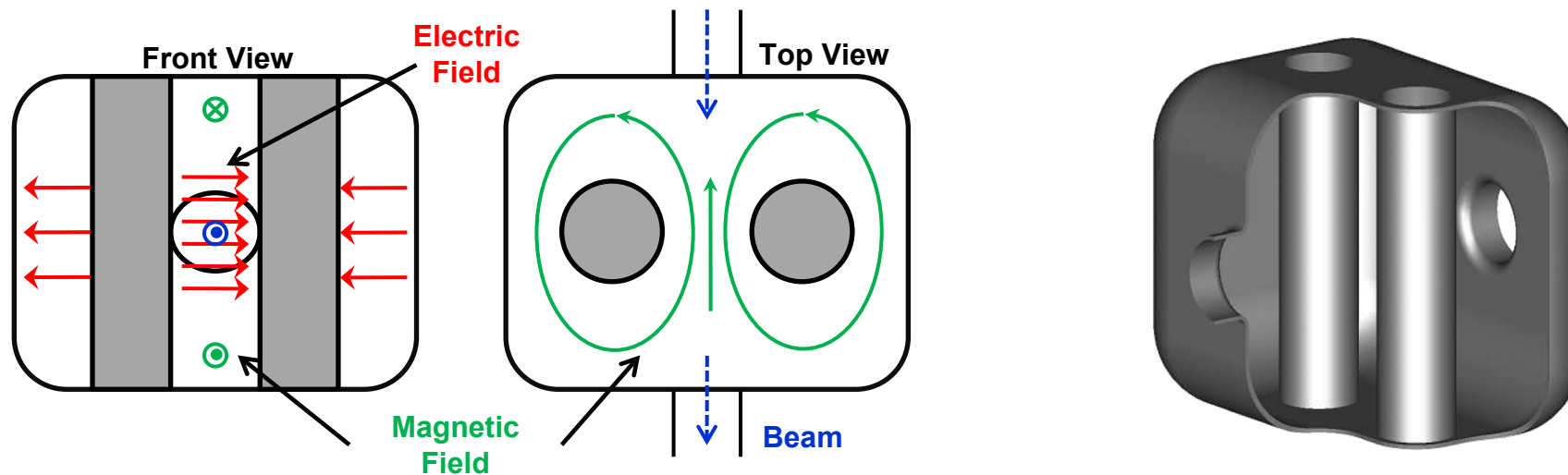


Crab cavities operated from
2007-2010



TEM-Type Deflecting/Crabbing Cavities

- Use both electric and magnetic fields to produce the net deflection

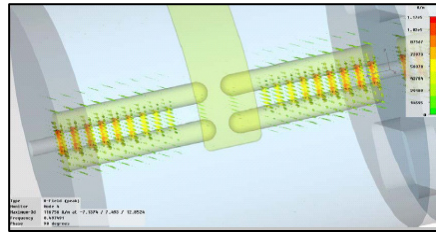
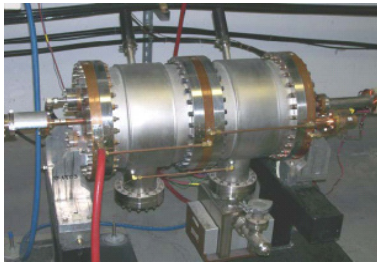


- At low operating frequencies gives:
 - Compact designs
 - Low surface fields and high shunt impedance
 - Some designs have no lower order modes
- New compact deflecting/crabbing designs are originated from TE-like or TEM-type designs

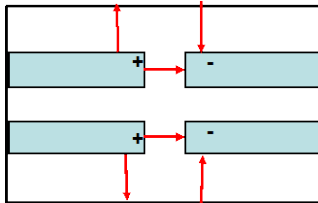
4-Rod Crabbing Cavity

- 4-Rod crabbing cavity – University of Lancaster / Cockcroft Institute
- Adapted from JLab normal conducting cavity
- Proposed for LHC high luminosity upgrade

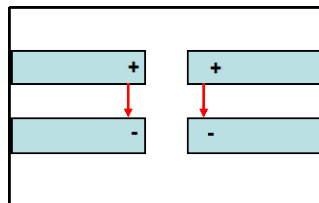
- Operates in a TEM-like mode
 - Uses both electric and magnetic fields
 - Deflecting mode is not the lowest mode



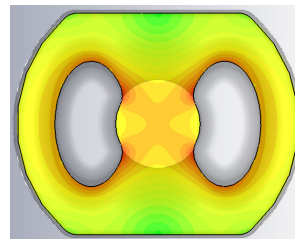
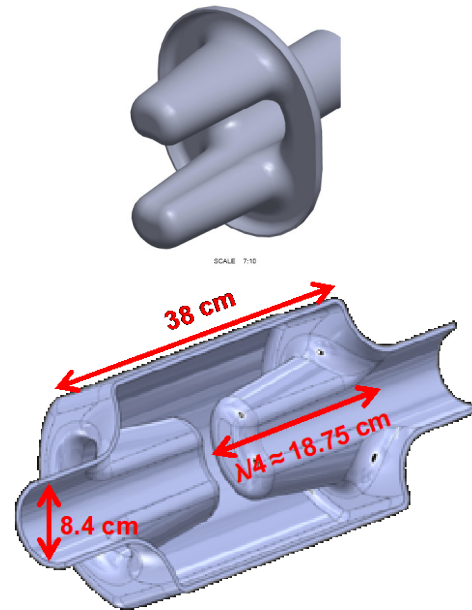
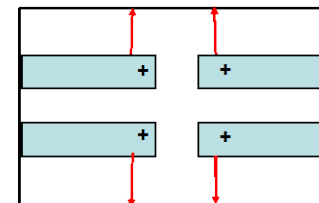
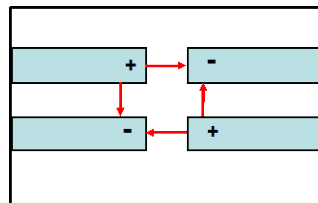
499 MHz normal conducting rf separator at JLab



Accelerating lower order mode



Fundamental deflecting mode

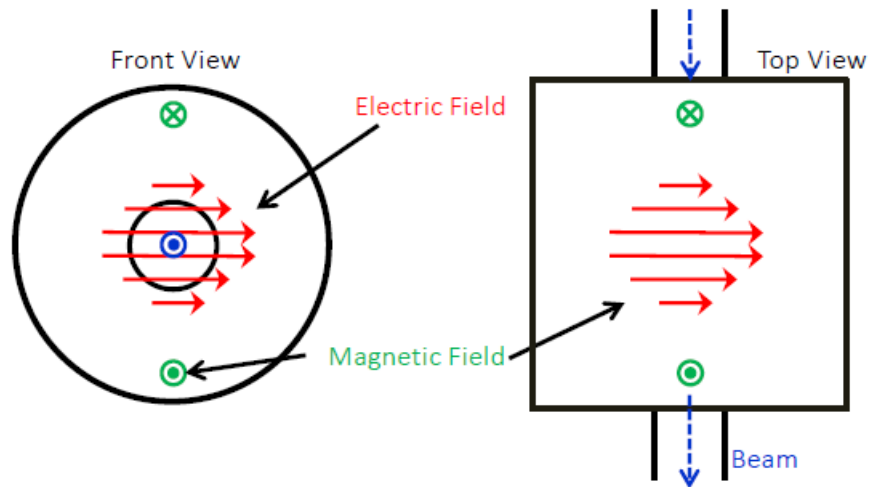


Rod shaping to reduce surface electric and magnetic fields, and the offset nonlinearities

Frequency	400.0	MHz
LOM	375.2	MHz
Nearest HOMs	436.6, 452.1	MHz
E_p^*	4.0	MV/m
B_p^*	7.56	mT
B_p^*/E_p^*	1.89	mT/(MV/m)
$[R/Q]_T$	915.0	Ω
Geometrical Factor (G)	62.8	Ω
$R_T R_S$	5.7×10^4	Ω^2
At $E_T^* = 1$ MV/m		

TE-Like Deflecting/Crabbing Cavities

- Operate in TE_{111} -like mode
 - Cannot be a pure TE_{111} mode where the contribution from electric and magnetic fields cancel each other
 - Main contribution to the transverse voltage is from transverse electric field

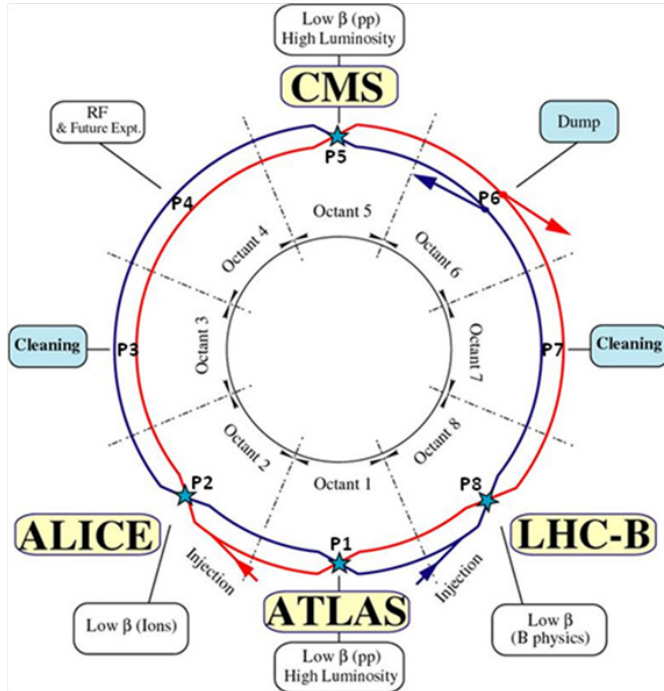


- Pure cylinder would cancel the contribution from E and B fields
- So need deformed shapes

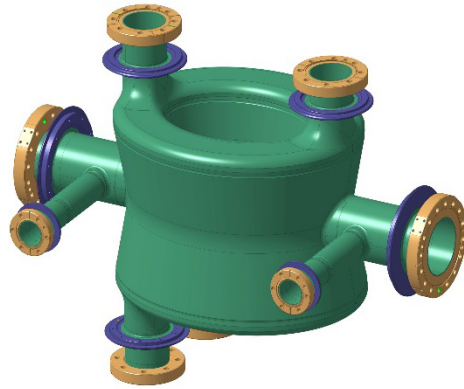
- Has similar rf properties as TEM-type cavities
 - Compact designs
 - Favorable for low frequencies (length $\sim \lambda/2$ and diameter $\sim 1/f$)
 - No lower order mode (TE_{111} is the lowest mode)
 - Have demonstrated transverse voltages at high peak surface fields

TE-Like Cavities for LHC High Luminosity Upgrade

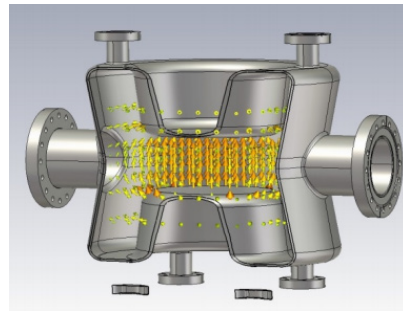
- Crabbing cavities for LHC high luminosity upgrade – Operate at 400 MHz



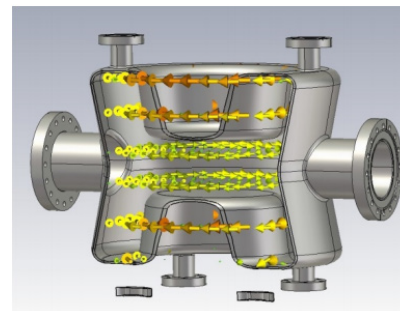
Double Quarter Wave Cavity



For vertical crabbing

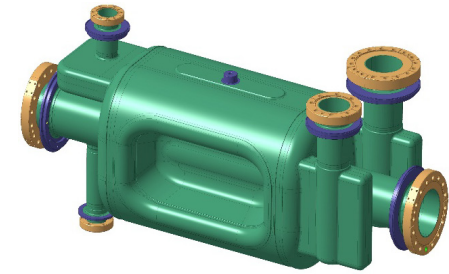


E Field

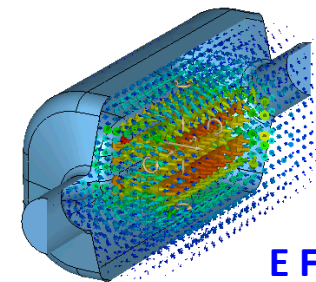


B Field

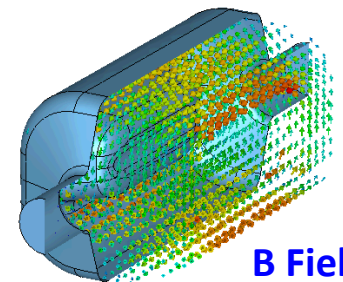
RF-Dipole Cavity



For horizontal crabbing



E Field



B Field

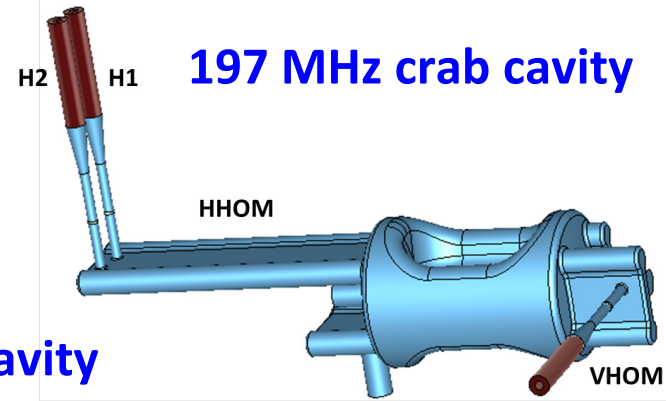
- Two interaction points (P1-ATLAS and P5-CMS)
- 8 cavities per interaction point
- Crossing angle – $285 \mu\text{rad}$

Electron-Ion Collider (EIC) at BNL

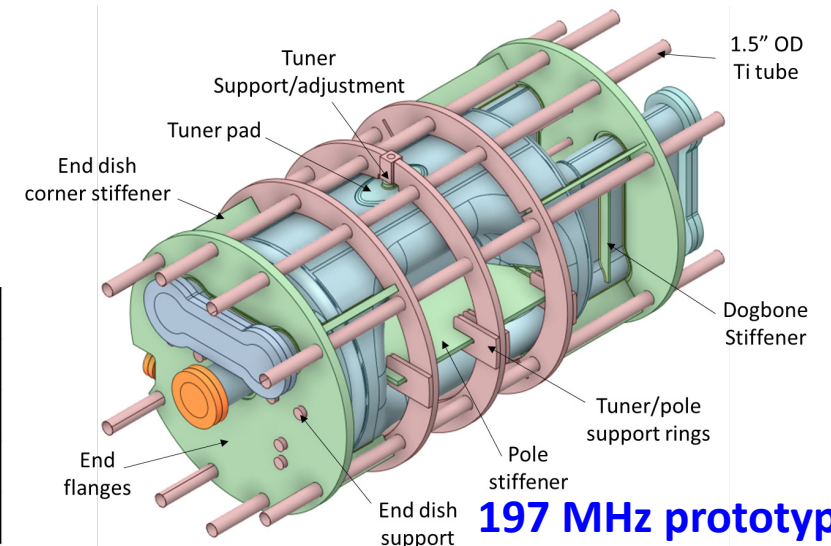
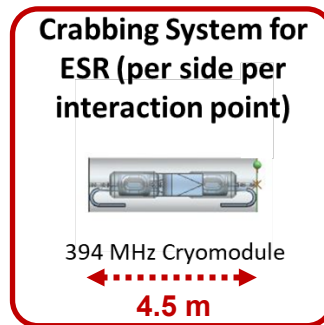
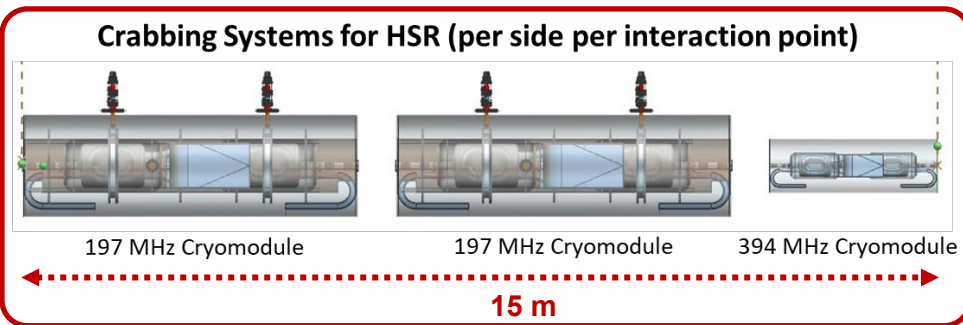
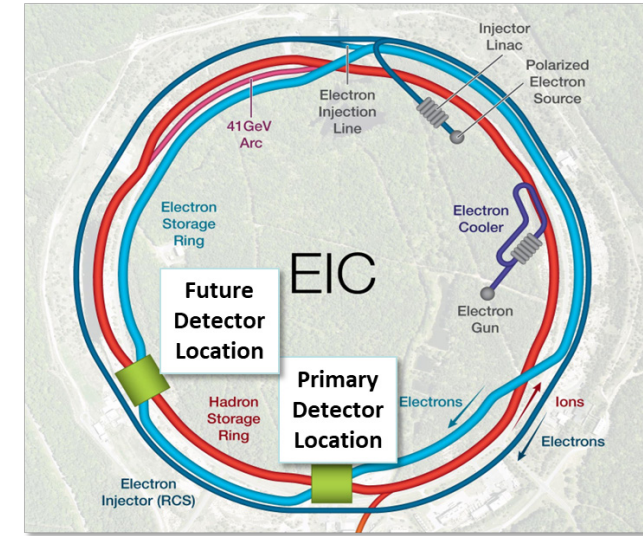
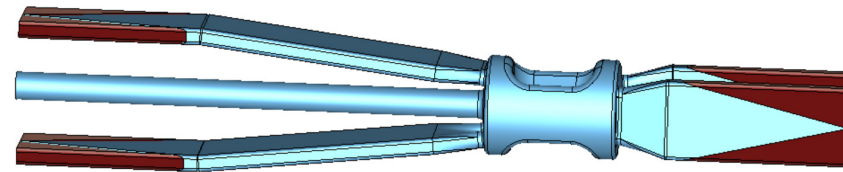
- Collisions of electrons and protons/ions for center mass of $\sim 30\text{-}140$ GeV
- Two rings: Electron storage ring (ESR) and Hadron storage ring (HSR)

	HSR		ESR
Cavity frequency [MHz]	197	394	394
Crossing angle [mrad]	25.0	25.0	25.0
β_{cc} [m]	1300	1300	150
β_{ip} [m]	0.8	0.8	0.45
Piwnski angle [rad]	7.9	7.9	2.4
Crabbing voltage [MV]	33.83	-4.75	2.9

System	V_t [MV]		No. of cavities	
	HSR	ESR	HSR	ESR
197 MHz	33.83	-	4	-
394 MHz	4.75	2.90	2	1



394 MHz crab cavity



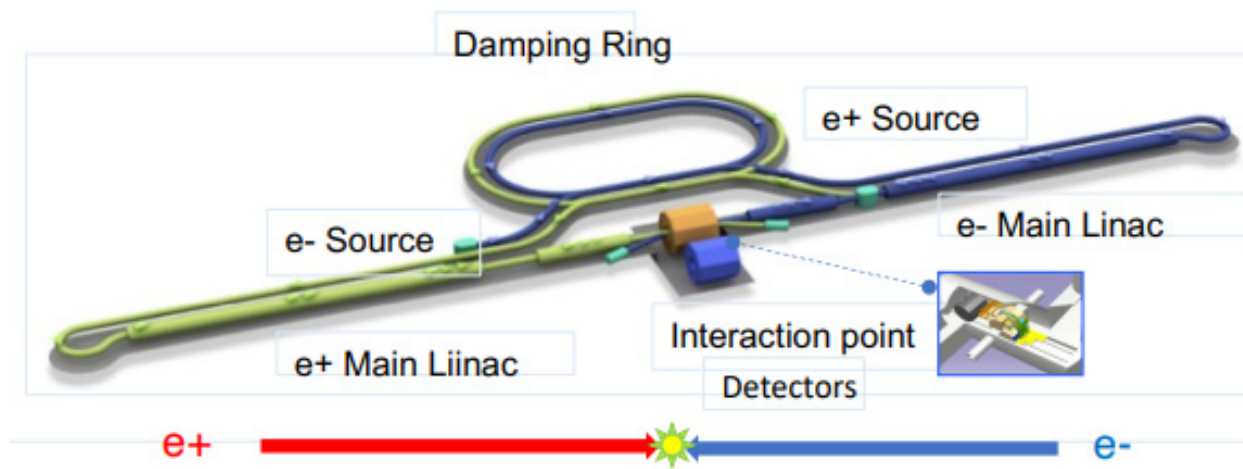
197 MHz prototype crab cavity

	Nb
	NbTi
	Ti
	SS
	Weld

Size and weight: 5' overall long, 30" OD, ~ 800 lb

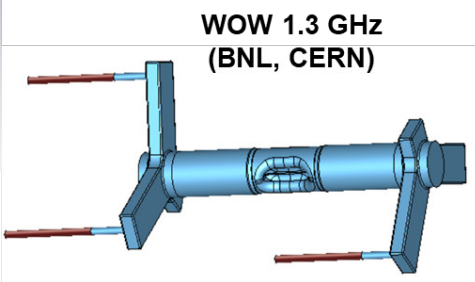
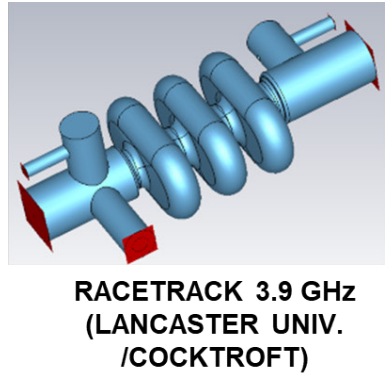
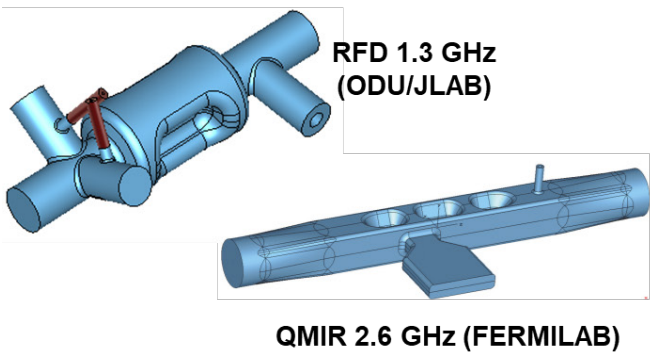
International Linear Collider (ILC)

- e^- and e^+ collision in a linear collider
- Crossing angle 14 mrad
- Luminosity: $1.35/2.7 \times 10^{10} \text{ cm}^2\text{s}^{-1}$



Parameter	Value
Center of mass energy [GeV]	250/1000
Crossing angle [mrad]	14.0
β_{cc} [m]	23.2
β_p [mm]	13.0/22.6
RMS bunch length (σ_z) [mm]	0.3/0.25
RMS horizontal beam size (σ_x^*) [nm]	729.0/481.0

Beam Energy	250 GeV			1 TeV		
	f_0 [GHz]	Tot. V_t [MV]		f_0 [GHz]	Tot. V_t [MV]	
f_0 [GHz]	1.3	2.6	3.9	1.3	2.6	3.9
Tot. V_t [MV]	1.845	0.923	0.615	7.4	3.7	2.5

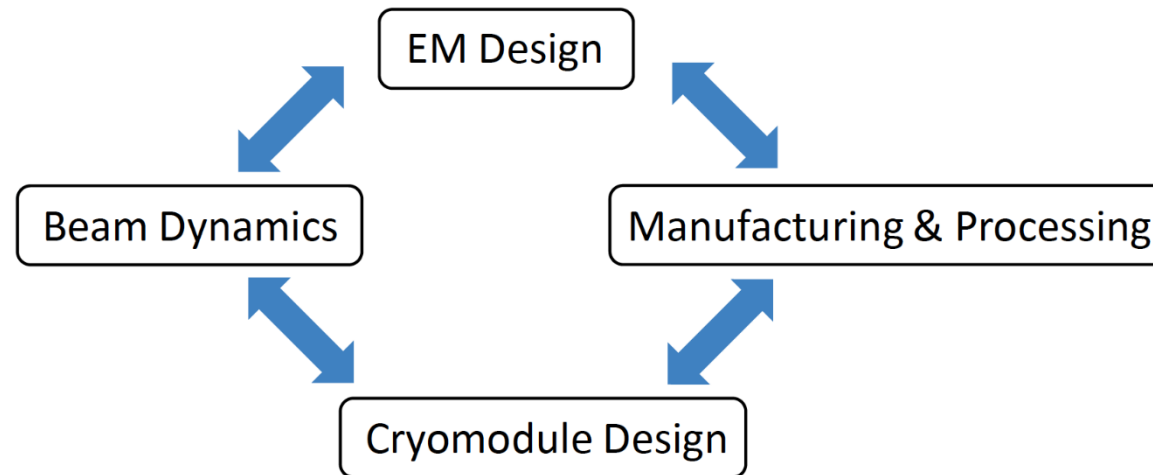


DESIGN CONSIDERATIONS OF NON-ELLIPTICAL ACCELERATING CAVITIES

- (1) DESIGN ASPECTS
- (2) MULTIPACTING
- (3) MECHANICAL DESIGN
- (4) FABRICATION
- (5) CHEMICAL ETCHING
- (6) CRYOMODULE DESIGN

Cavity Design Considerations

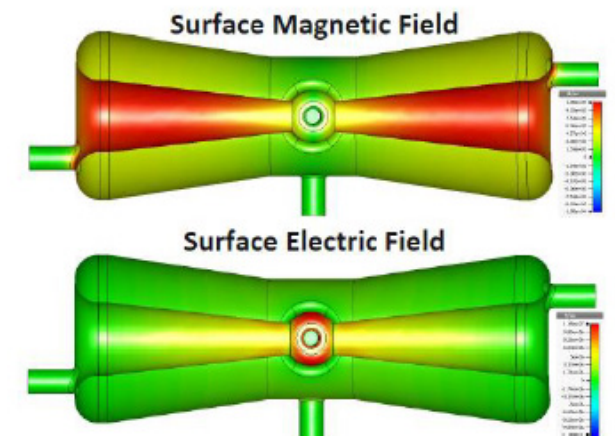
- Beam dynamics:
 - Cavity frequency
 - Beam aperture
 - Voltage acceptance



- Ideally cavities should have:
 - Large accelerating gradient (E_a) / Energy gain (W)
 - Large shunt impedance ($R_{sh}=G*(R/Q)$) for low losses to reduce power consumption
 - Shapes that reduce peak fields (E_p, B_p) for given E_a
 - Efficient energy transfer to the beam (β)
 - Reduce multipacting levels
- In addition, related practical issues of these cavities
 - Reduce pressure sensitivity (df/dp)
 - Microphonics
 - Operation: cw or pulsed
 - Cavity tuning
 - Cavity fabrication
 - Chemical processing, cleaning, and assembly

Designing Non-Elliptical Cavities

- Cavity frequency:
 - To minimize unique number of cavity designs
- Cavity β : Number of cavity designs also depend on required velocity range
 - $T(\beta)$ is efficient over a range of velocities from $0.7\beta_0 < \beta < 2\beta_0$ (Especially for QWR)
 - For $\beta > 0.5$ possible to consider multi-spoke cavities where the reduced transit time factor is compensated by the higher voltage
 - Maintain a certain cavity type until $T(\beta)$ lowers the voltage below the voltage of the next cavity series
 - For post accelerators with different ion acceleration – β profile should be chosen that all ions can be accelerated near the maximum gradient
- Peak surface fields: $E_p \leq 35$ MV/m and $B_p \leq 70$ mT
 - Dominates the optimization between practicality and complexity
- Multipacting analysis: Multipacting levels may not be eliminated, but reduced by optimizing the design
 - Advanced simulation tools exist in simulating multipacting resonance levels in cavities that have matched with measurements

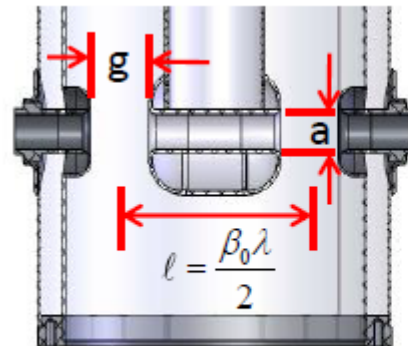


Designing Non-Elliptical Cavities

- Baseline mechanical and fabrication model
 - Consider ports: FPC, HOM couplers, Rinsing and surface treatment
 - Choose material thickness: Based on stress analysis
 - Maintain safe limits in terms of stress
 - Check all pressure differential throughout cavity life cycle
 - Minimize Lorentz force detuning due to radiation pressure → Stiffeners
 - Tuner designs
- Integrating design for fabrication variables include:
 - Cavity performance
 - Complexity in geometry
 - Operational requirements: 4 K or 2 K
 - Stress analysis
 - Material cost vs machining cost

RF Frequency Considerations

- Smaller structures are cheaper to fabricate so favor small λ (high frequency)
- Lower frequencies increase size of stable region in longitudinal phase space and increase cavity active length so increase the effective voltage
- Lower velocities require lower frequencies
- Typically, $\beta\lambda/2$ between gaps (one accelerating cell)
- Gap is typically half a cell or $g \sim \beta\lambda/4$
- Aperture should be less than the gap to improve acceleration efficiency \rightarrow so $a < \beta\lambda/4$



g =gap, a =aperture

Example:

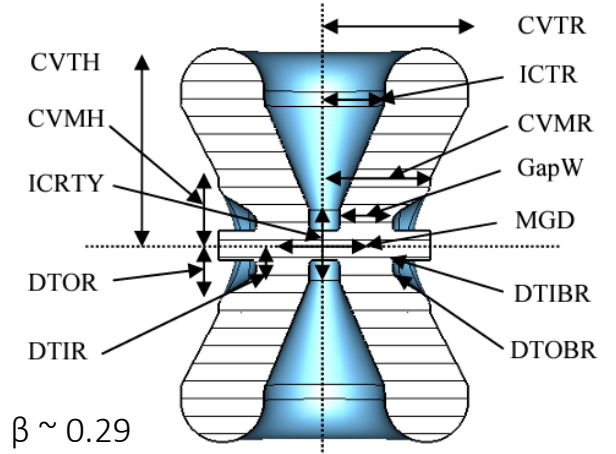
for $a=30\text{mm}$ and $\beta = 4\%$

$\lambda > 3\text{m}$ or $f < 100\text{MHz}$

Design Optimization

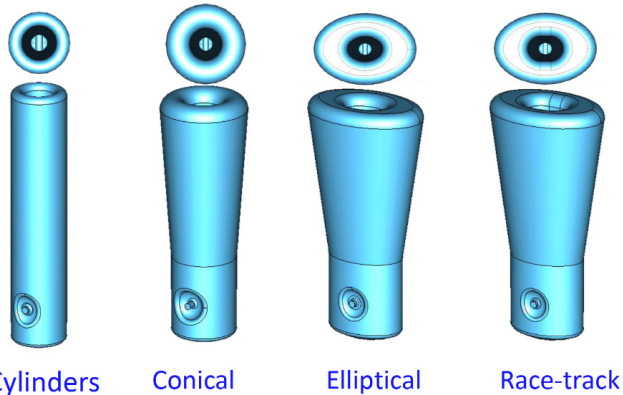
- Good simulation tools exist both on the EM and mechanical aspects (CST, HFSS, COMSOL, ANSYS, SLAC ACE3P Code Suite, ...)
- Many geometrical parameters of optimization compared to elliptical cavities
- Optimize cavity parameters to:
 - Minimize peak surface field ratios – E_p/E_{acc} and B_p/E_{acc} (E_p/E_t and B_p/E_t for deflecting and crabbing cavities)
 - Maximize $R_{sh} = G^*(R/Q)$

HWR



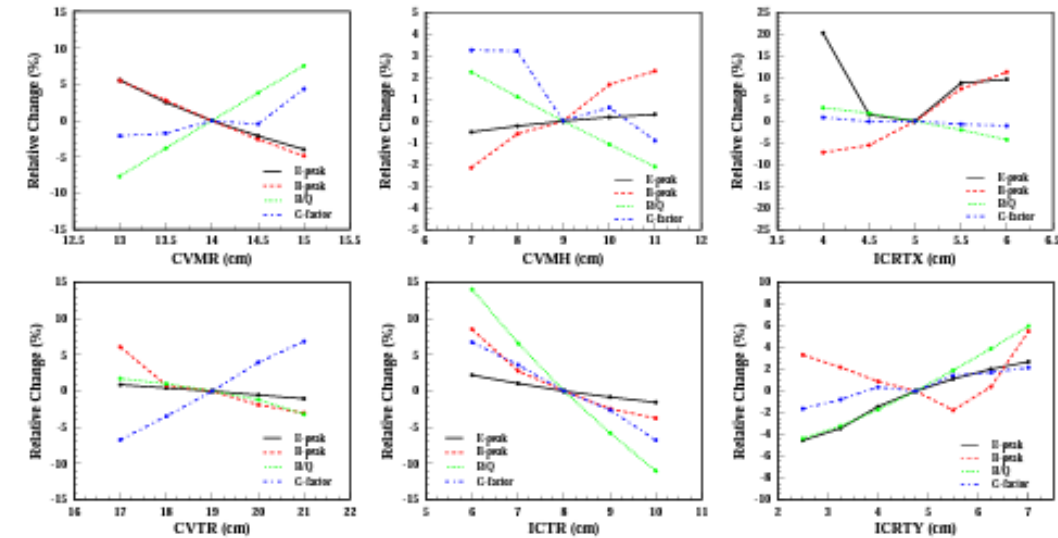
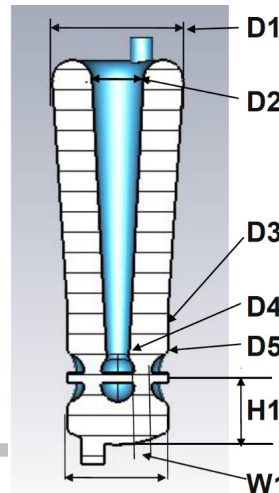
322 MHz, $\beta \sim 0.29$

QWR



Some primary cavity geometrical parameters:

- Cavity Top Diameter (D1)
- Stem Top Diameter (D2)
- Cavity Lower Diameter (D3)
- Stem Bottom Diameter (D4)
- Drift Tube Outer Diameter (D5)
- Drift Tube Gap Width (W1)
- Cavity Bottom Height (H1)

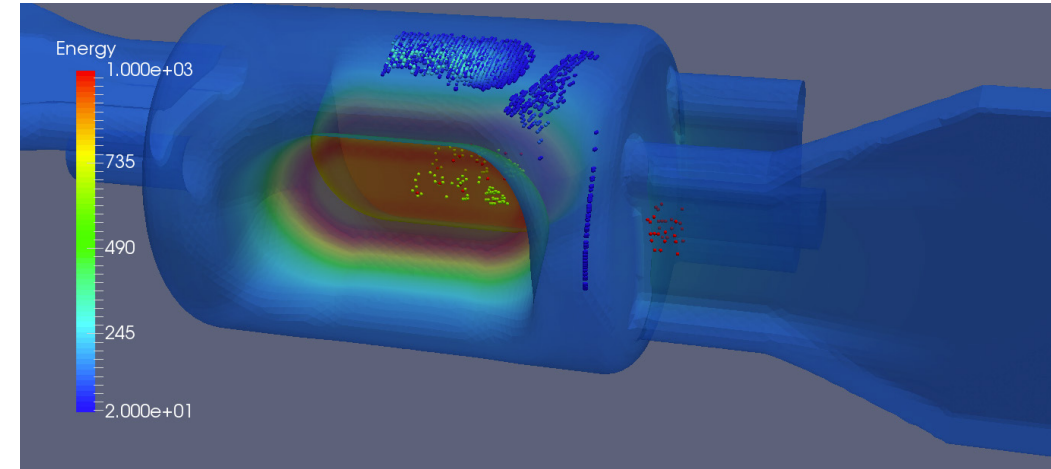
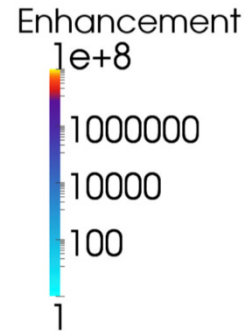
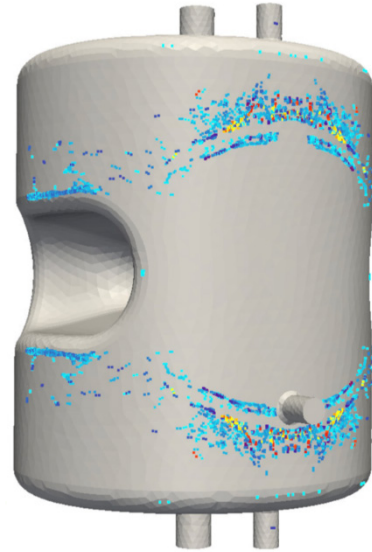
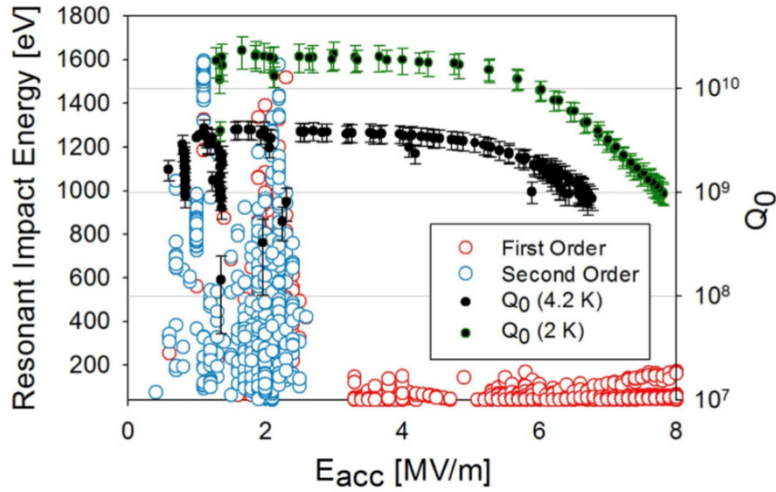


Cylinders Conical Elliptical Race-track

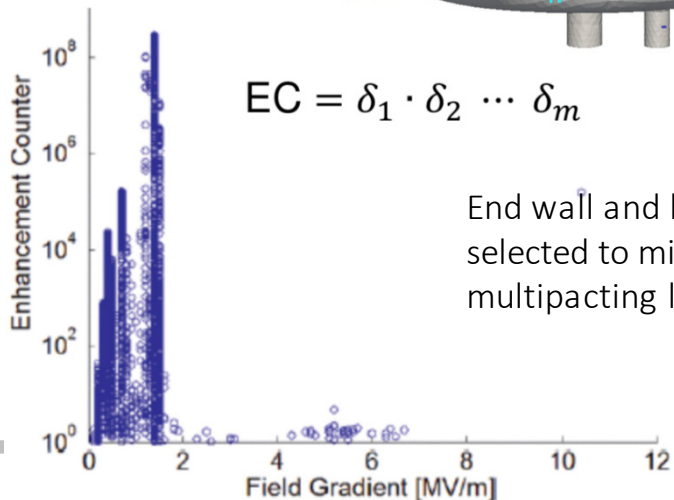
Increasing complexity, lower surface fields

Multipacting in Non-Elliptical Cavities

- Multipacting always occur in non-elliptical cavities due to complex geometries → But not a showstopper
- Now reliable tools exist that can model multipacting resonant levels

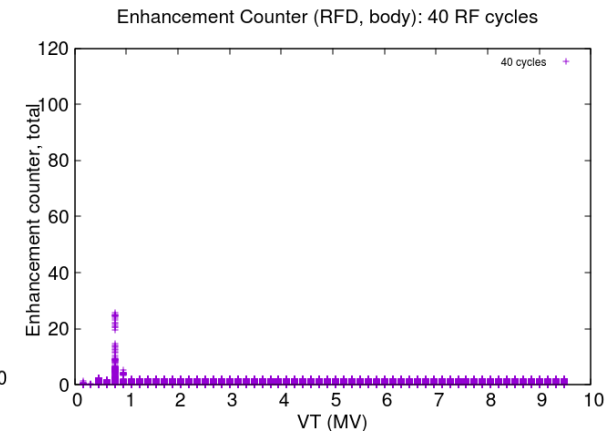
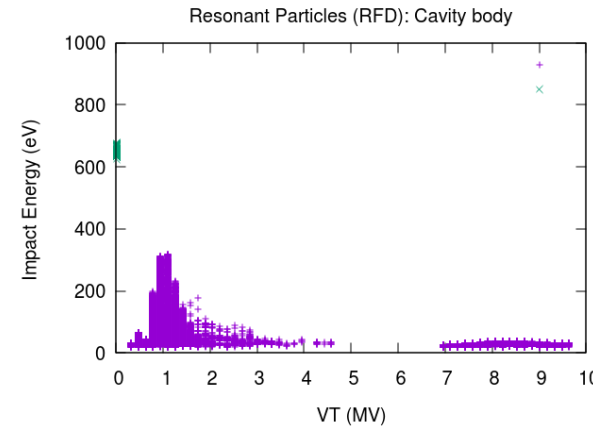


Multipacting experienced up to 2.1 MV/m



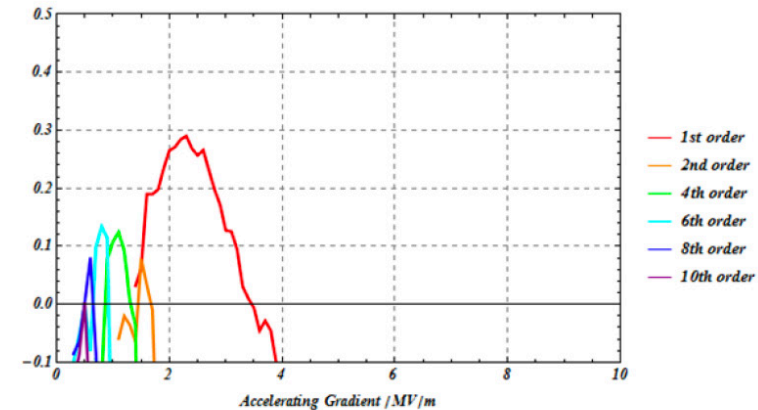
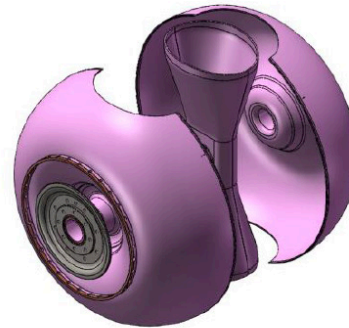
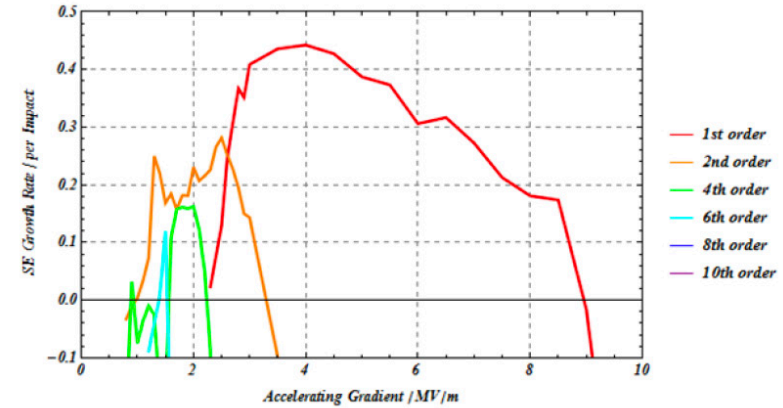
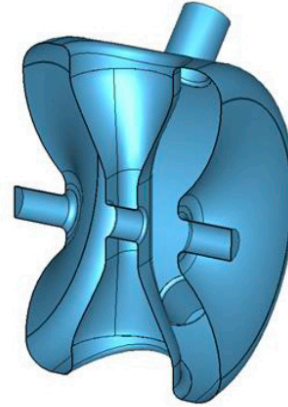
$$EC = \delta_1 \cdot \delta_2 \cdots \delta_m$$

End wall and bend radius selected to minimize multipacting levels



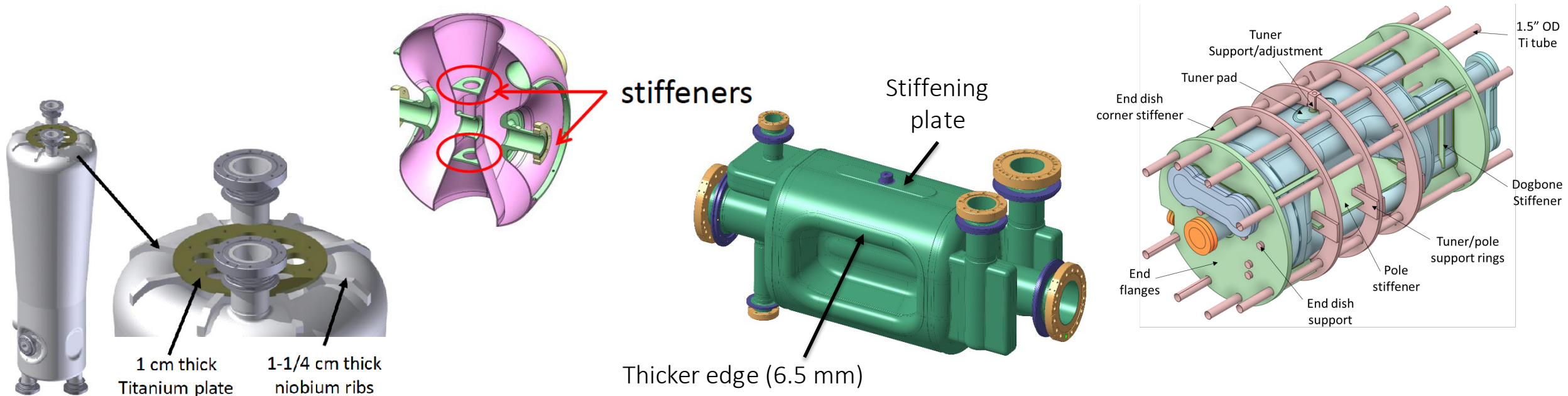
Multipacting in Non-Elliptical Cavities

- Multipacting can be reduced careful cavity design
- Example – single spoke resonator designed by TRIUMF for RISP
- Balloon variant reduces serious multipacting in operating region



Mechanical Designs of Non-Elliptical Cavities

- Mechanical design focuses on reducing internal stresses under the external pressure load for various operational conditions
- Study mechanical stability to microphonics, pressure fluctuations (df/dp), Lorentz force detuning
- Stiffeners can be added strategically to reduce and improve mechanical stability
- Consider thermal performance when considering thicker material to reduce stresses
- Consider tuning range and tuning force required and cavity stresses for maximum tuning range

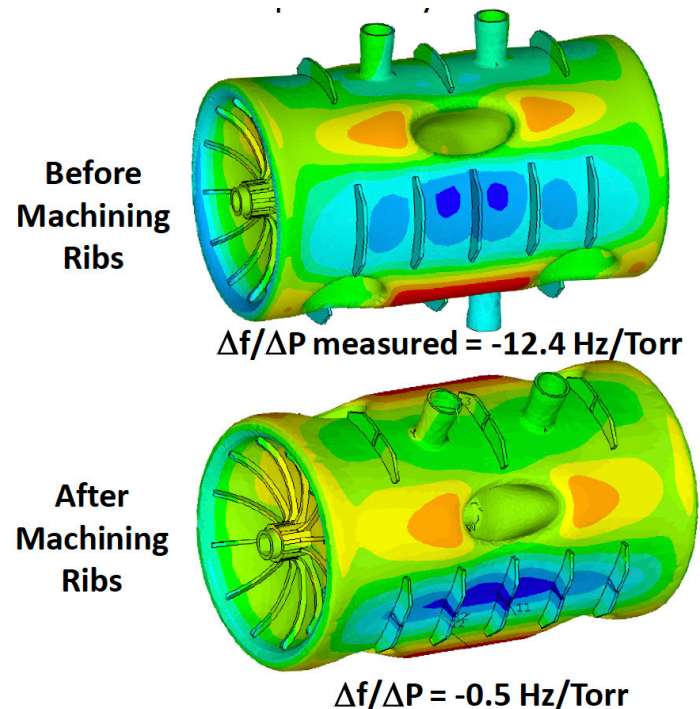


Mechanical Designs of Non-Elliptical Cavities

- Response of the cavity to external pressure changes can be dominant contribution at 4 K operation
- Lorentz force detuning: inward pressure in E field and H field regions produce frequency shifts in opposite direction

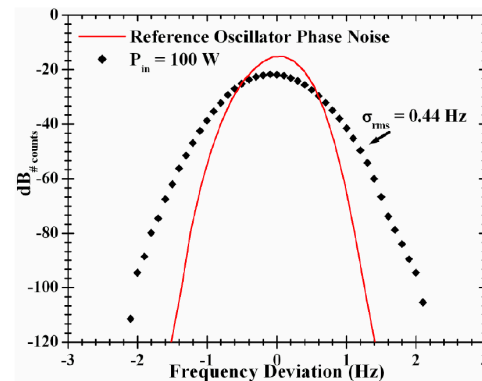
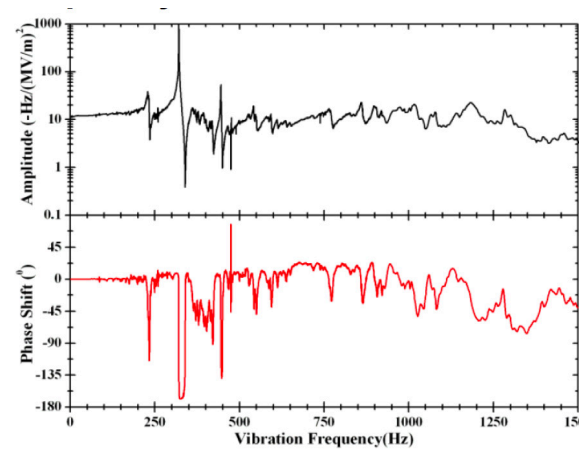
$$\Delta f = k \int (\epsilon_0 E^2 - \mu_0 H^2) dV$$

- Support ribs in E field and H field to balance deflections so that frequency shift can be cancelled



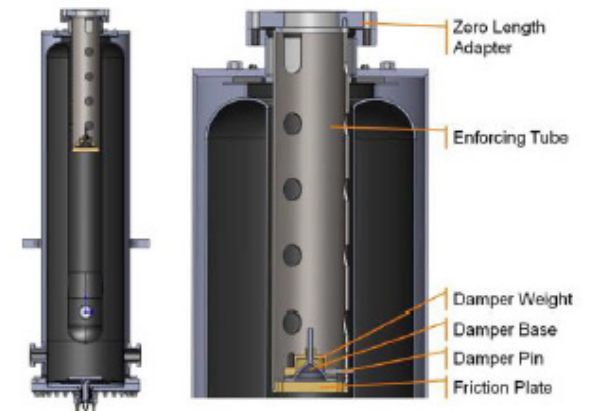
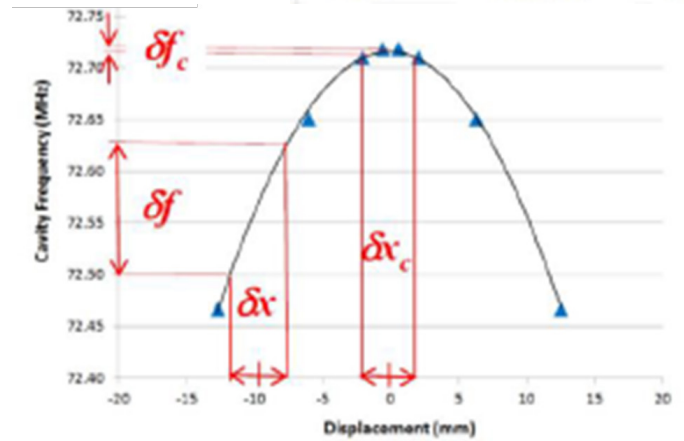
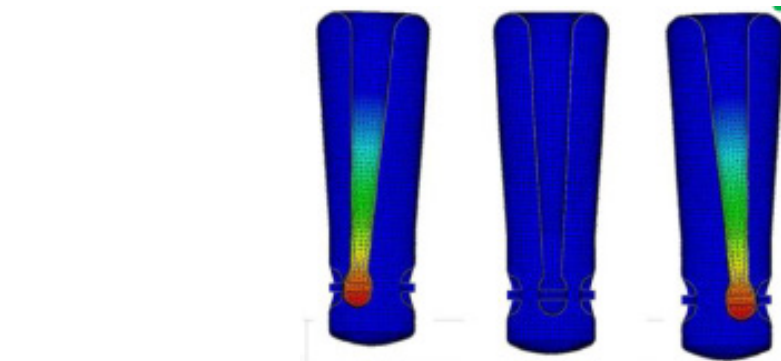
Mechanical Designs of Non-Elliptical Cavities

- Driven by mechanical vibration in the environment
- QWRs inner conductor can have low mechanical frequencies (50-100Hz)
- Need to reduce the RMS detuning to $\ll 10\%$ of the design bandwidth to avoid nuisance unlocks
- Mitigations
 - Stiffening during the design/manufacture to raise the mechanical frequency of the lowest modes
 - Centering the inner conductor by plastic deformation so $df/dx=0$
 - Adding passive dampers in QWR
 - Reducing environmental noise in and around the cryomodule -operation at 2K is quieter than 4K



$$df/dp = -0.4 \text{ Hz/mbar}$$

345 MHz $\beta=0.5$ triple spoke cavity

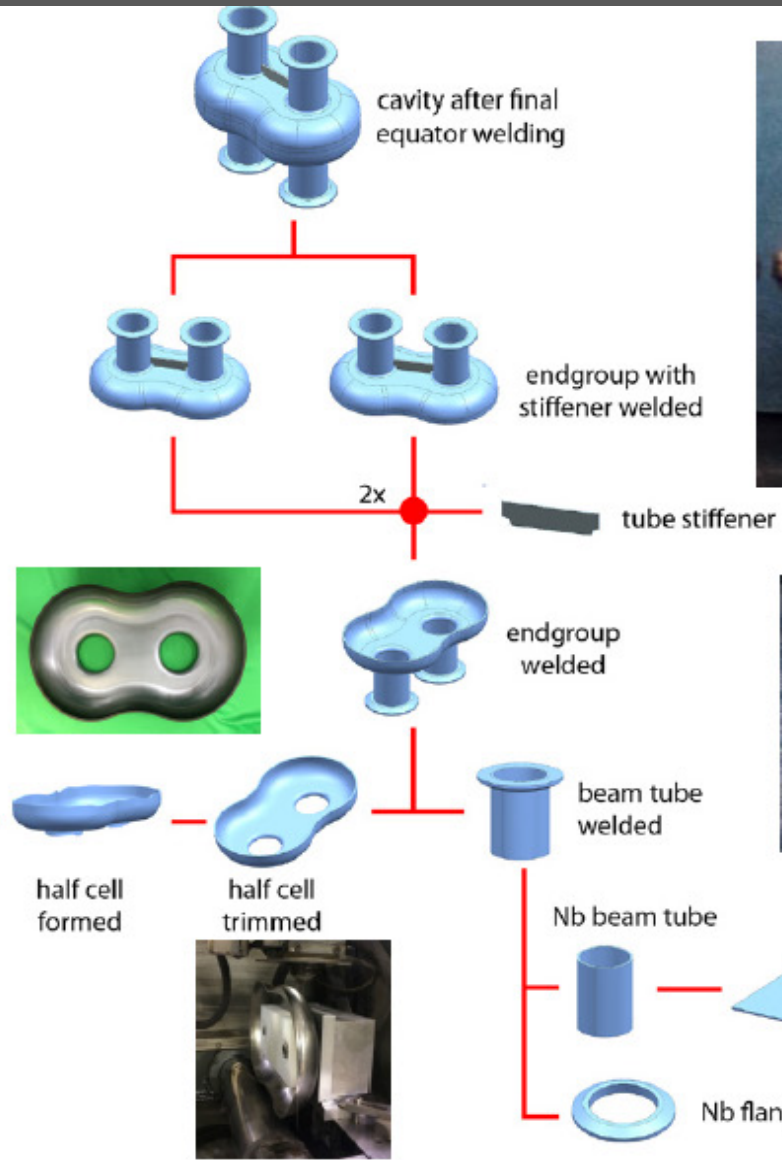
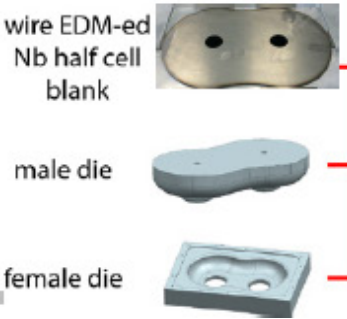


Cavity Fabrication

- Fabrication of TEM mode cavities must follow the same standards as for elliptical cavities – the differences can come from the novel shapes and unique forming steps
 - Consists of many more subcomponents compared to elliptical cavities
 - Requires many more machining/trimming fixtures at different stages of the fabrication process
 - Frequency tuning requires different fixturing compared to elliptical cavities
 - May need additional plastic deformation of welded subcomponents and post fabrication tuning
- New procedures can be qualified in Copper as a `cheap` substitute for Nb or RG Nb
- Fabrication steps can involve
 - Machining of parts from bulk – cost optimization between material cost and fabrication cost – EDM, milling, lathe
 - Forming –deep drawing from sheet –good for production quantities
 - Spinning from sheet –fast, not as precise as forming typically
 - EB welding in high vacuum for RF surfaces to maintain RRR value of Nb
 - TIG in controlled atmosphere for jacket

**SRF 2023 Tutorial –
Cavity Fabrication and
Preparation, Rongli Geng**

Cavity Fabrication



Cavity Fabrication

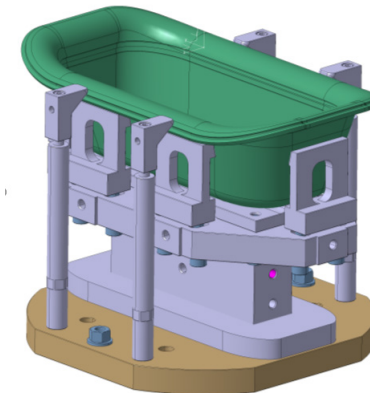
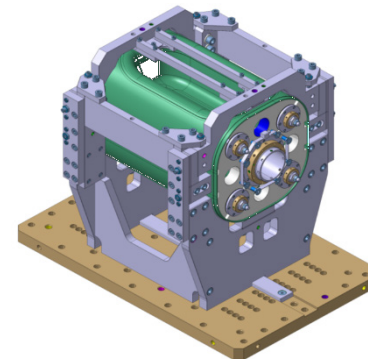
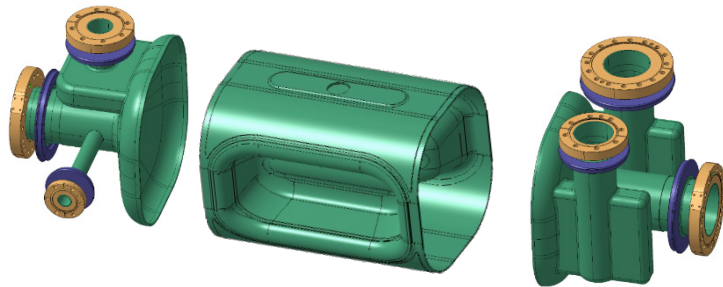
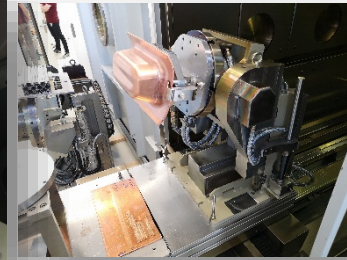
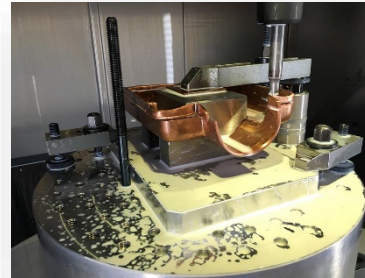
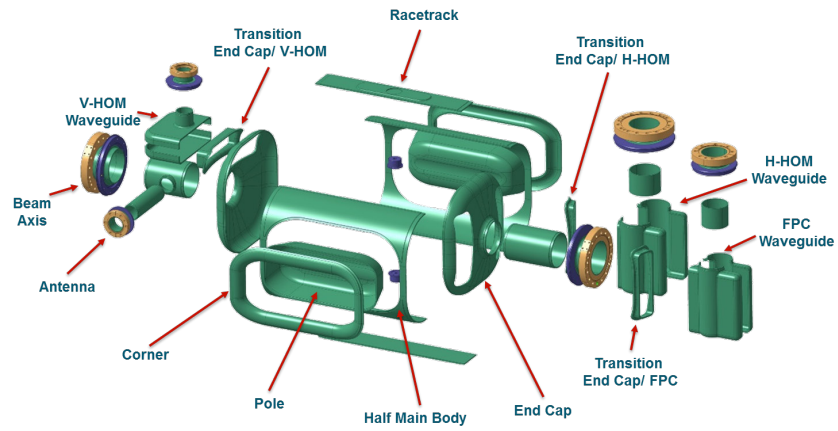
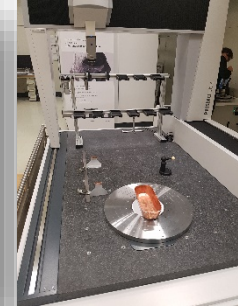
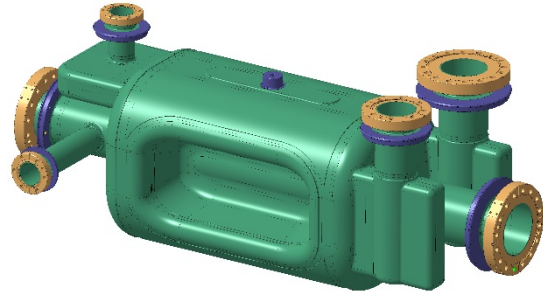
325 MHz Single-Spoke
Fabricated at Niowave Inc.



500 MHz Double-Spoke
Fabricated at Jefferson Lab
(HyeKyoung Park)



Fabrication of Non-Elliptical Cavities



Cavity Surface Treatment, Processing and Assembly

- A standard recipe

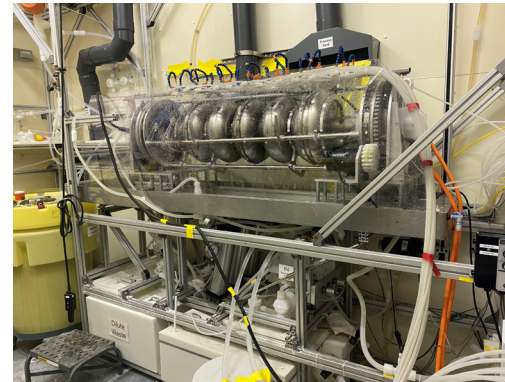
1	Bare Cavity Inspection – Visual, Dimensional, Vacuum, RF
2	US cleaning and rinse
3	BCP 120-150 μm (flip half-way)
4	High-Pressure Rinse
5	Hydrogen Degassing 800 °C, 10 h
6	RF Tuning
7	BCP 20-30 μm
8	HPR (horiz + vert)
9	Clean Room Assembly
10	Low Tem Bake 120 °C, 48 h
11	Vertical Test @ 2K
12	Helium Vessel Welding
13	US cleaning
14	BCP 20-30 μm
15	HPR
16	Clean Room Assembly
17	Low Temp Bake 120 °C, 48 h
18	Horizontal Test @ 2K
19	Ready for String Assembly



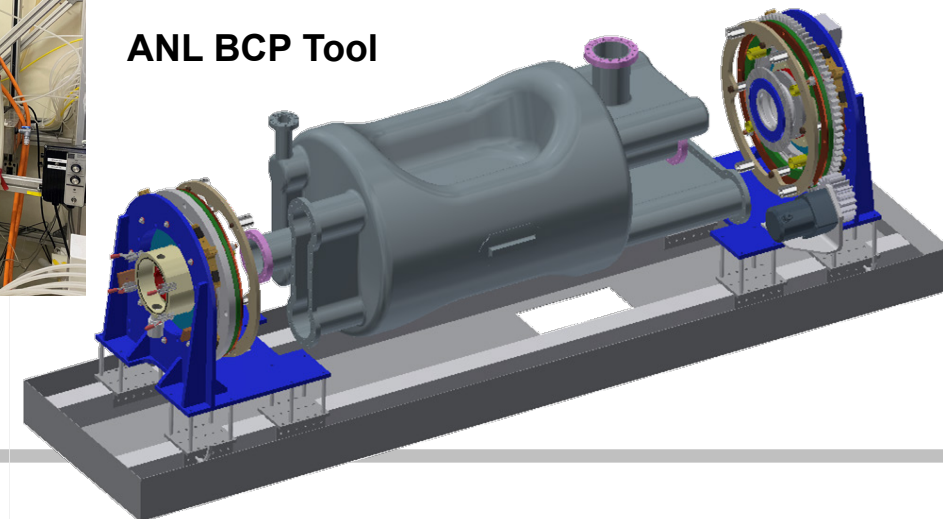
High Temp Furnace - JLAB



Low Temp Bake Box

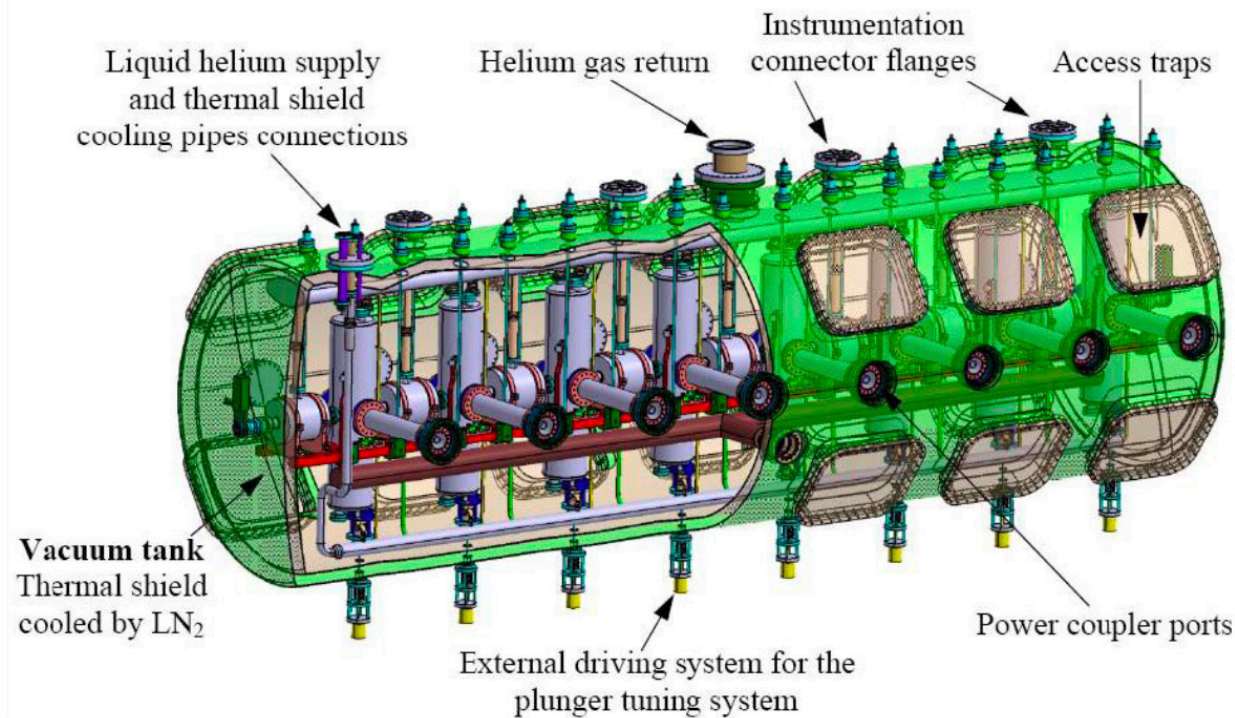


ANL BCP Tool



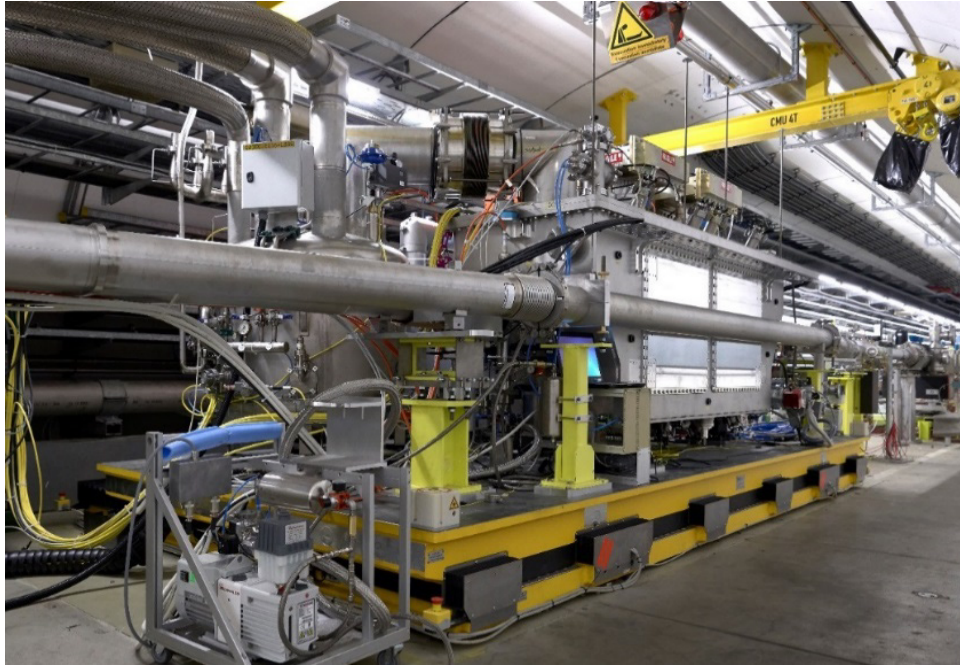
Cryomodule Designs

- Non-elliptical cavities require more complicated cryomodule structures compared to elliptical cavities
- Different solutions investigated for the same cavity types
- Couplers, tuners, and rf lines are often dominant components

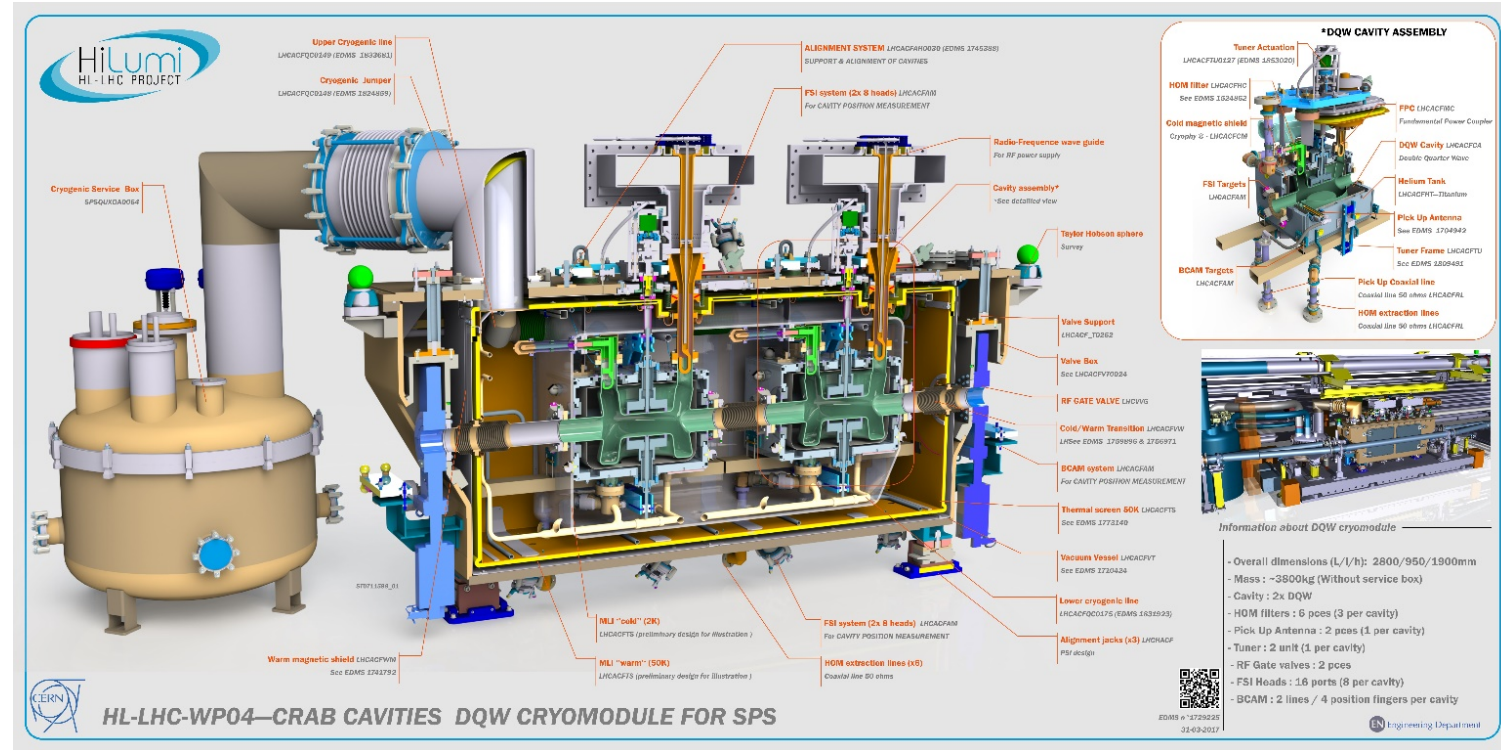


SRF 2023 Tutorial –
Cryomodule Design,
Nicolas Bazin

Cryomodule Designs



400 MHz Double Quarter Wave Cryomodule for SPS, CERN



Final Comments

- Non-elliptical cavities have evolved in the past several decades → Tremendous global interest in superconducting cavities for high energy hadron acceleration, compact high beta cavities, and deflecting/crabbing cavities
- Many accelerators are using non-elliptical cavities such as QWR, HWR, Spoke that have moved the technologies to industrial production aiming high performance and reliability
- Compact crabbing cavity designs have advanced particle colliders in achieving high luminosities
- Parameter, tradeoff, and option space available to the designer is large → No universal design where designs are application specific
- Design process is not reduced to a few simple rules of recipes → Room for pushing the performance of the non-elliptical cavities to higher accelerating/crabbing voltages
- Look out for non-elliptical contributions at SRF 2023

References

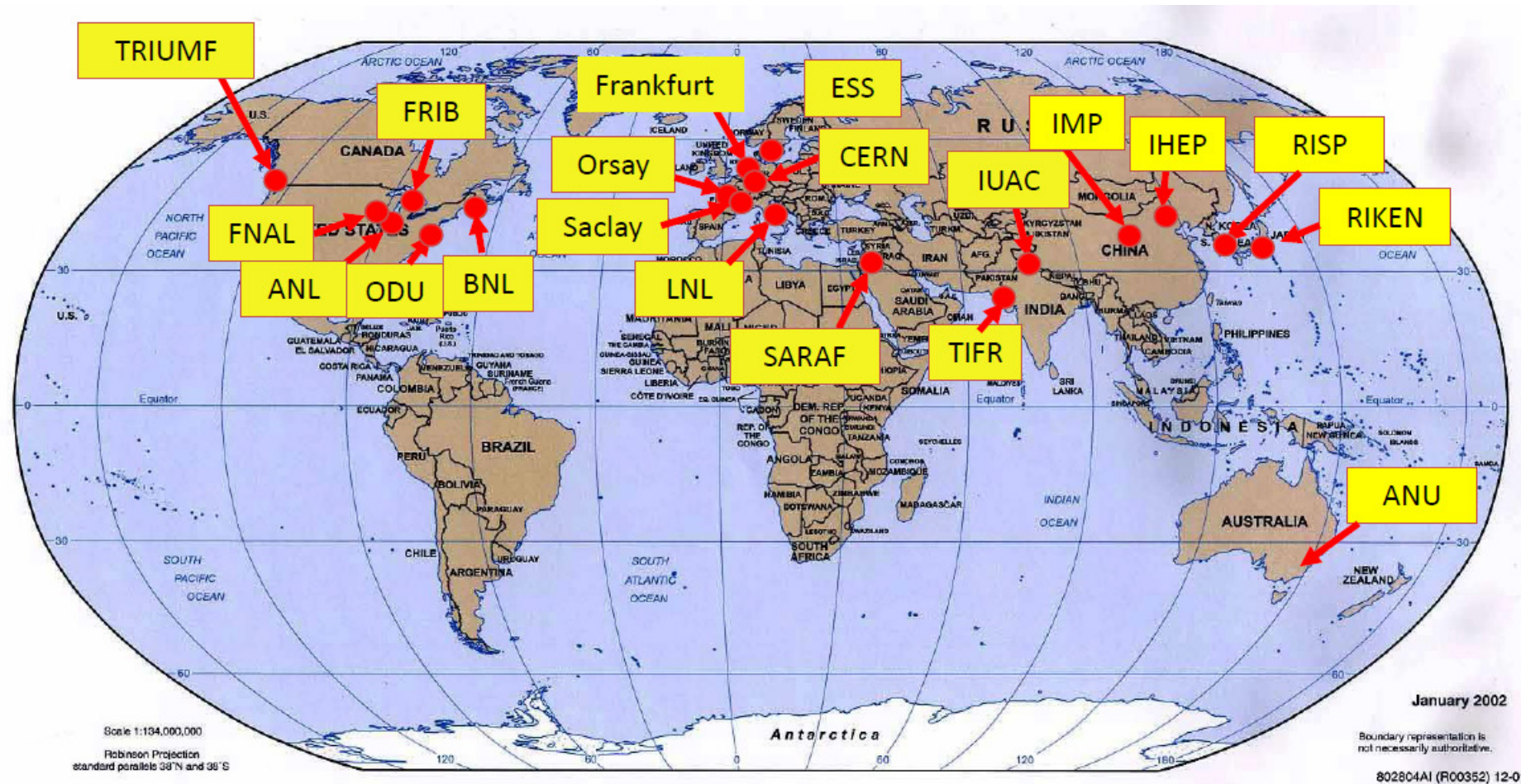
- Past SRF Tutorials:
 - SRF 2021 – Non-Elliptical Cavities – R. Laxdal
 - SRF 2019 – Non-Elliptical Cavities – S. De Silva
 - SRF 2017 – Non-Elliptical Cavities – R. Laxdal
 - SRF 2015 – Non-Elliptical Resonators – A. Facco
 - SRF 2013 – TEM Class Cavity Design – M. Kelly
 - SRF 2007 – Low and Medium β Superconducting Cavities and Accelerators - J. Delayen
- Non-Elliptical Accelerating Cavities
 - Design of Low-Velocity Superconducting Accelerating Structures Using Quarter-Wavelength Resonant Lines – J. R. Delayen, NIM, A259, 341-357 (1987)
 - Longitudinal Transit Time Factors of Short Independently-Phased Accelerating Structures for Low Velocity Ions – J. R. Delayen, NIM, A258, 15-25 (1987)
 - Superconducting Spoke Cavities – M. Kelly, 39th ICFA Advanced Beam Dynamics Workshop on High Intensity High Brightness Hadron Beams (2006), <http://accelconf.web.cern.ch/AccelConf/abdwhb06/PAPERS/THAY07.PDF>
 - Medium- β Superconducting Accelerating Structures – J. R. Delayen, SRF 2001, <https://accelconf.web.cern.ch/accelconf/srf01/papers/fa007.pdf>
 - Applications of Spoke Cavities – J. R. Delayen, LINAC 2010, <https://accelconf.web.cern.ch/accelconf/LINAC2010/papers/tu302.pdf>
 - Superconducting RF Development For FRIB At MSU - K. Saito, et. al., LINAC 2014, <http://inspirehep.net/record/1363521/files/thioa02.pdf>

References

- Non-Elliptical Accelerating Cavities
 - Superconducting Spoke Cavities for High-Velocity Applications – C. S. Hopper and J. R. Delayen, Phys. Rev. ST Accel. Beams, 16, 102001, (2013)
 - High-Velocity Spoke Cavities – C. S. Hopper, SRF 2015, https://digitalcommons.odu.edu/cgi/viewcontent.cgi?article=1178&context=physics_fac_pubs
 - Multipacting Analysis of High-Velocity Superconducting Spoke Resonators – C. S. Hopper and J. R. Delayen, LINAC 2012
 - STUDY OF BALLOON SPOKE CAVITIES, Z.Y. Yao, et. al., SRF 2013, <http://accelconf.web.cern.ch/AccelConf/SRF2013/papers/thp033.pdf>
- Non-Elliptical Deflecting/Crabbing Cavities
 - Energy Scaling Crab Crossing and the Pair Problem – R. B. Palmer, SLAC-PUB-4707 (1988)
 - Design Evolution and Properties of Superconducting Parallel-Bar RF- Dipole Deflecting and Crabbing Cavities – S. U. De Silva and J. R. Delayen, Phys. Rev. ST Accel. Beams 16 (2013) 012004
 - Design, Prototyping and Testing of a Compact Superconducting Double Quarter Wave Crab Cavity - B. P. Xiao et al., Phys. Rev. ST Accel. Beams 18, 041004 (2015)
 - Designing the Four Rod Crab Cavity for the High-Luminosity LHC Upgrade - B. Hall, EuCARD Monograph Vol. 25 (2014)
 - SPX Crab Cavity Development and Testing Result – H. Wang, TTC Meeting (2012), <https://www.jlab.org/indico/event/24/session/9/contribution/36/material/slides/1.pdf>

ADDITIONAL SLIDES

Superconducting Non-Elliptical Cavity Community



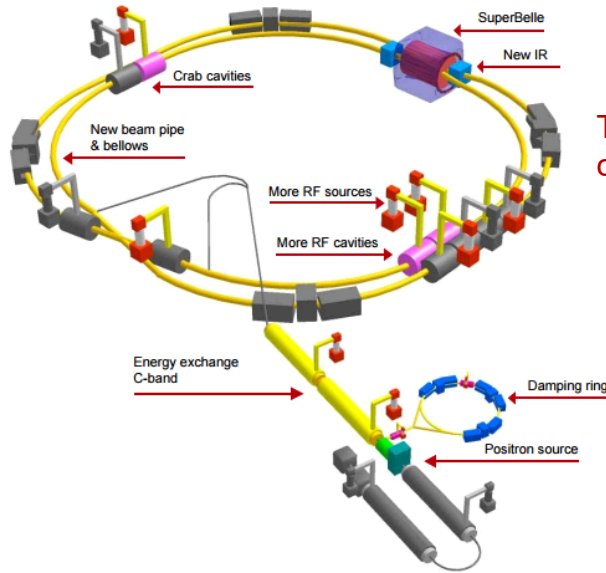
Non-Elliptical Facilities and Projects

Project	Lab	Driver	Post-accelerator	Particle	Structure
ATLAS	ANL		√	HI	Split-ring, QWR
ALPI	LNL		√	P, d / HI	QWR (sputter, bulk)
ISAC-II	TRIUMF		√	HI	QWR
IUAC	IUAC		√	HI	QWR
ReA3/6	NSCL		√	HI	QWR
HIE-Isolde	CERN		√	HI	QWR (sputter)
SARAF	SOREQ	√		P, d	HWR
SPIRAL-II	GANIL	√		P, d, HI	QWR
IFMIF	Saclay	√		P,d	HWR
FRIB	NSCL	√		HI	QWR, HWR
ESS	ESS	√		P	DSR
RAON	RISP	√		HI	QWR, HWR, SSR
ADS	IMP,IHEP	√		p	HWR, SSR
PIP-2	FNAL	√		P	HWR, SSR
Hi-Lumi	CERN			p	Crab cavities - DQWR, RFD

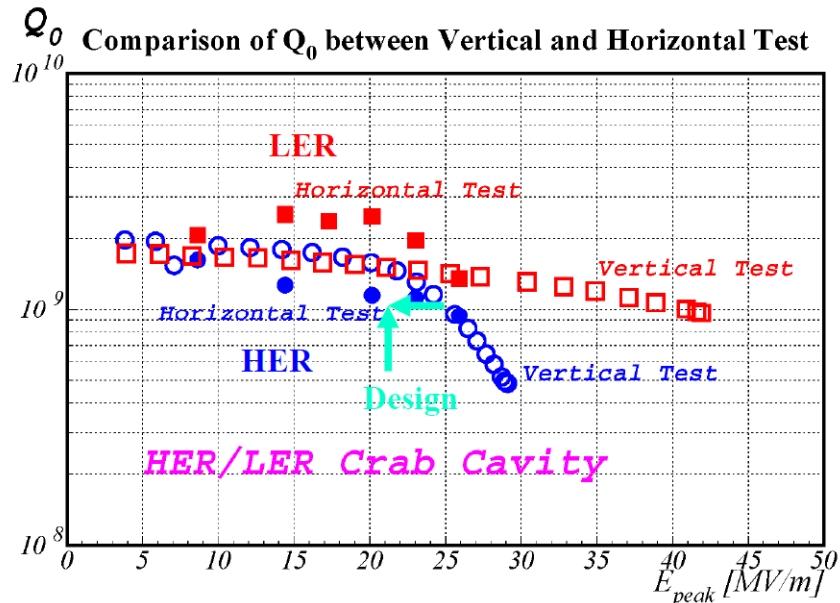
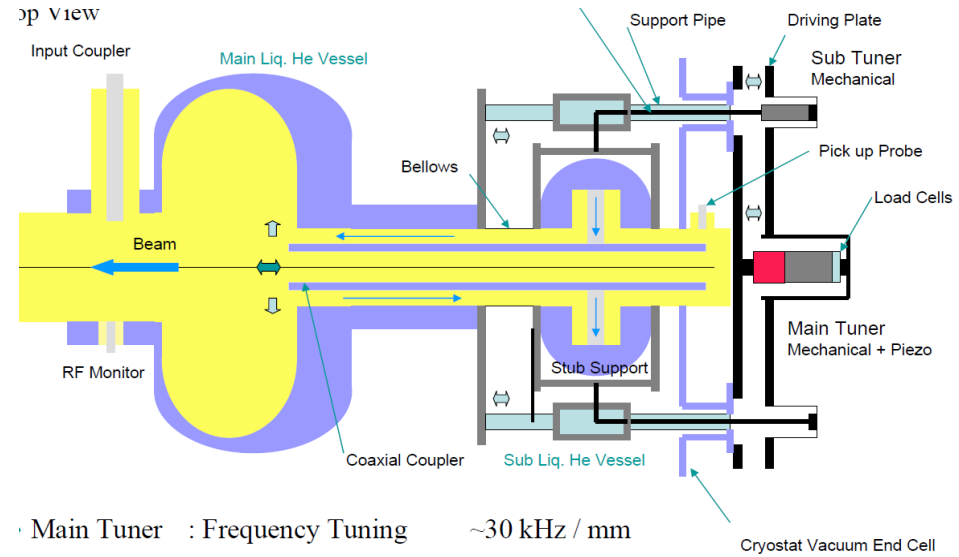
7/10/2017

Bob Level - Non-Elliptical Cavities

KEK Crabbing Cavity



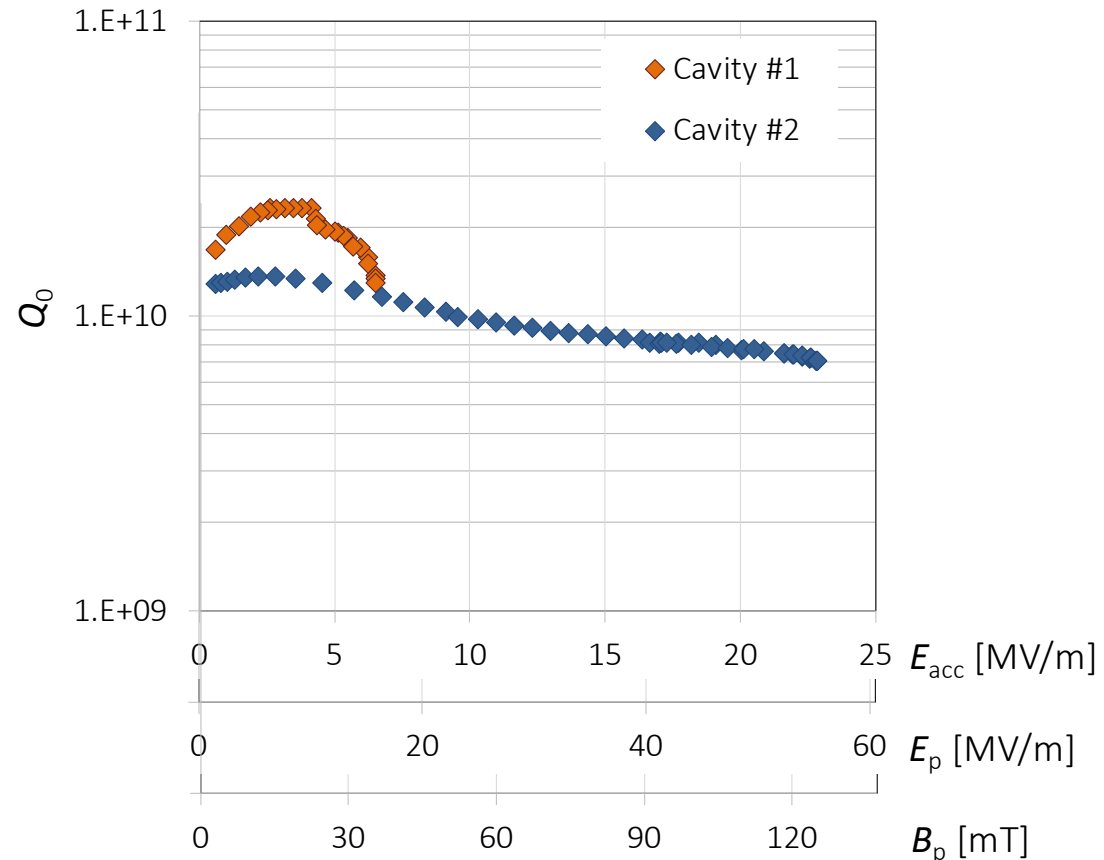
- Two crabbing cavities
- High energy ring (HER)
 - Low energy ring (LER)



Frequency	508.9	MHz
LOM (TM ₁₁₀)	410.0	MHz
Nearest HOMs	630.0, 650.0, 680.0	MHz
E_p^*	4.24	MV/m
B_p^*	12.23	mT
B_p^*/E_p^*	2.88	mT/(MV/m)
$[R/Q]_T$	48.9	Ω
Geometrical Factor (G)	227.0	Ω
$R_T R_S$	1.11×10^4	Ω^2
At $E_T^* = 1$ MV/m		

RF Test Results of Twin Axis Cavity

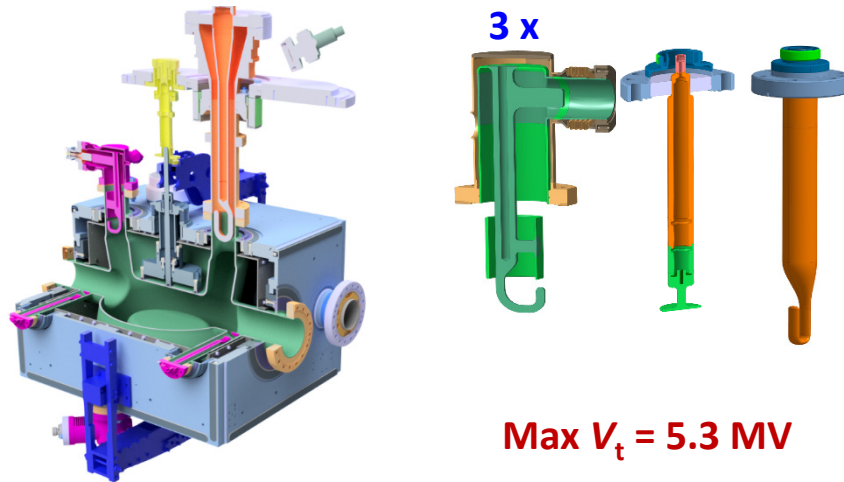
- Cavity 1:
 - Low field Q_0 of 2.3×10^{10}
 - Cavity quenched at 6.5 MV/m
 - Weld defect seems to be the limiting factor according to OST measurements
- Cavity 2:
 - Low field Q_0 of 1.3×10^{10}
 - Achieved an accelerating gradient of 23 MV/m ($E_p = 56$ MV/m & $B_p = 126$ mT)
 - No multipacting levels were observed during the test
 - Given minimal surface treatment (BCP only, no EP) cavity reached high Q_0



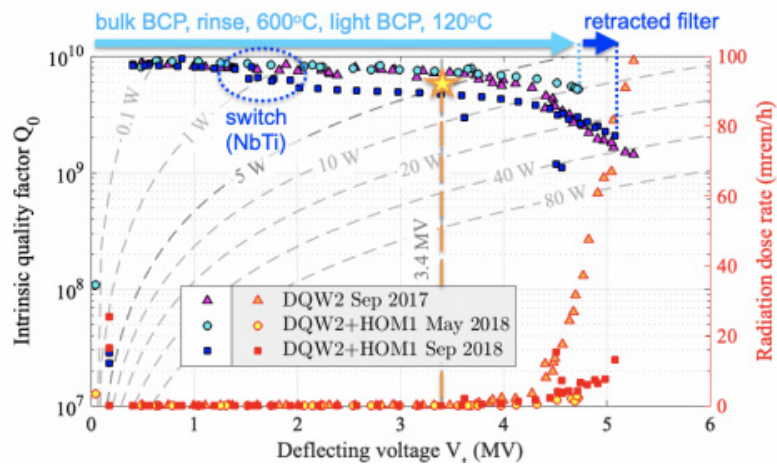
TE-Like Cavities for LHC High Luminosity Upgrade

- Crabbing cavities for LHC high luminosity upgrade – Operate at 400 MHz

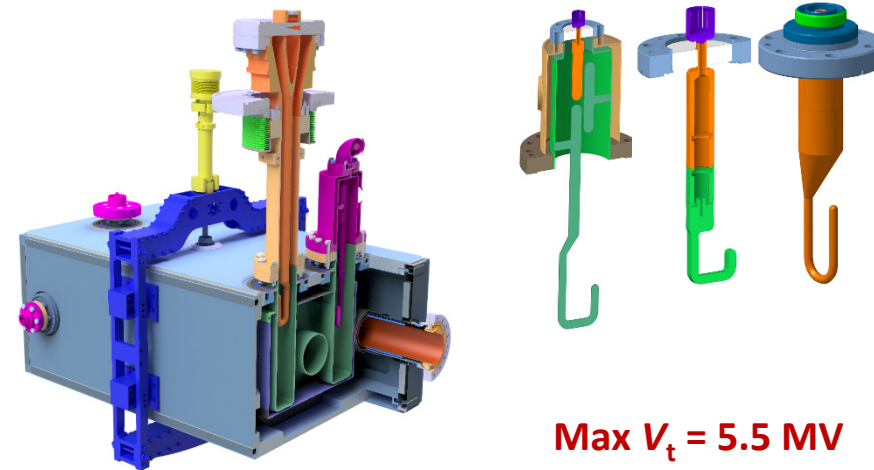
Double Quarter Wave Cavity



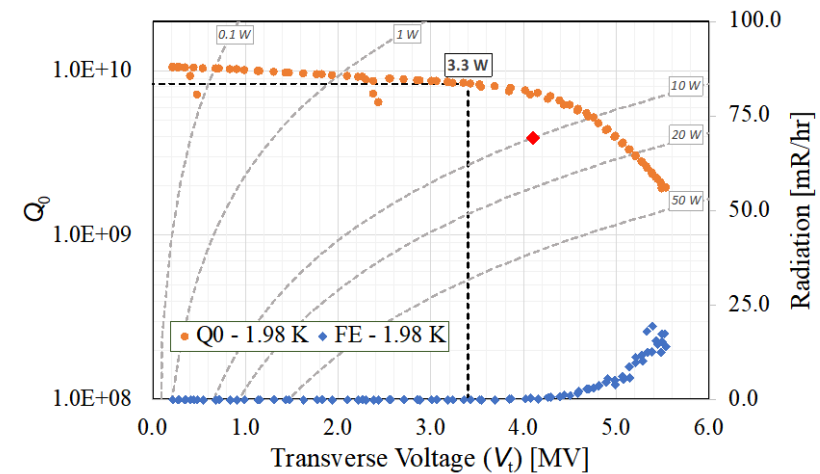
Max $V_t = 5.3$ MV



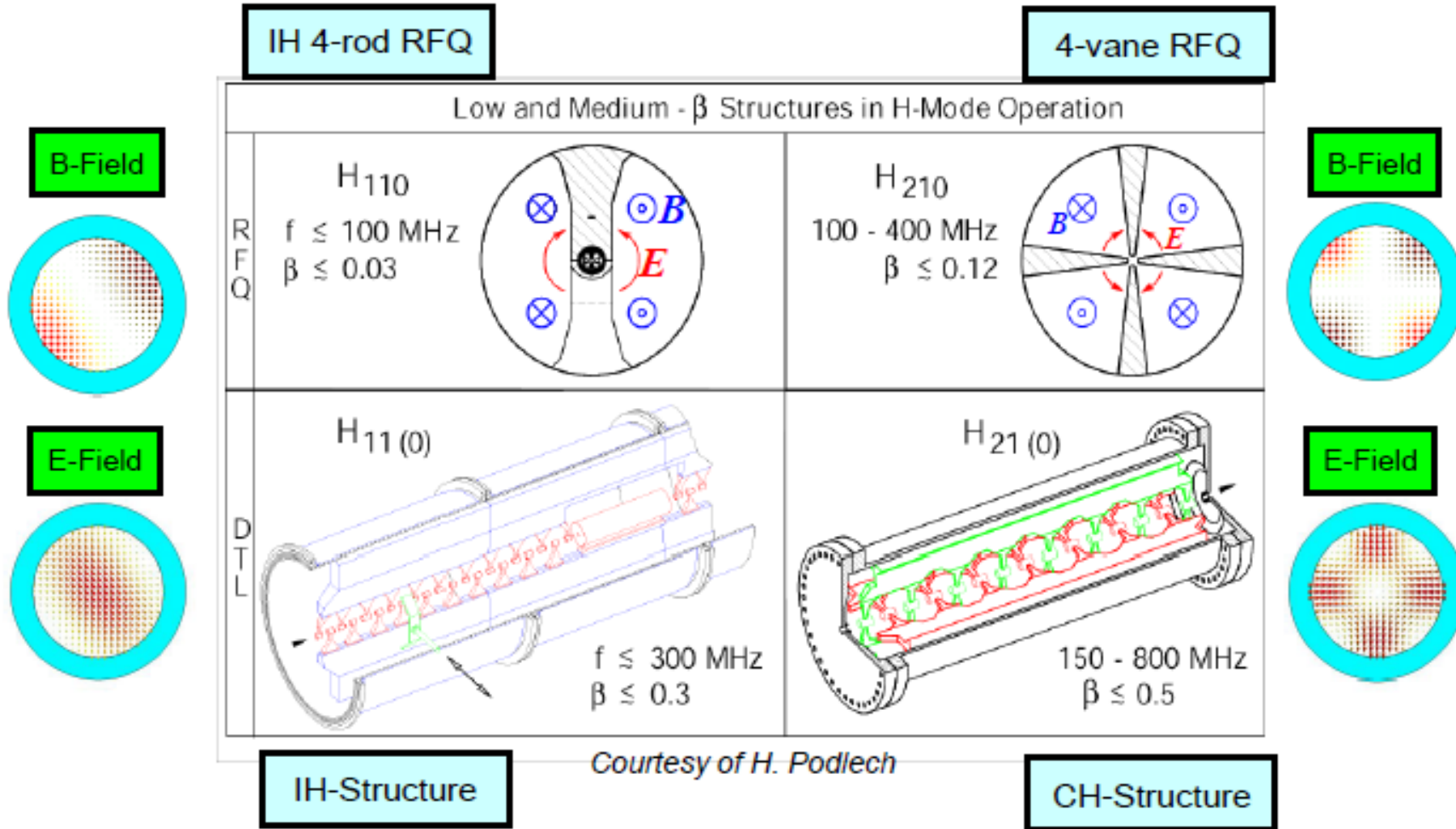
RF-Dipole Cavity



Max $V_t = 5.5$ MV



TE-Type Non-Elliptical Cavities



Courtesy of H. Podlech

Normalized $T(\beta)$

- It is usually convenient to define the normalized transit time factor and include the gap effect in the accelerating gradient

- Normalized transit time factor:
$$T^*(\beta) = \frac{T(\beta)}{T(\beta_0)}$$

- Average accelerating gradient:
$$E_a^* = T(\beta_0)E_a$$

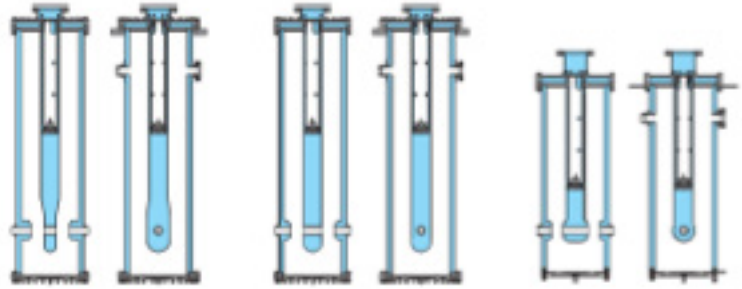
where $\beta_0 \equiv \beta / T(\beta_0) = \max\{T(\beta)\}$ and $T^*(\beta_0) = 1$

- Energy gain definition doesn't change
$$\Delta W_p = qE_a^*LT^*(\beta)\cos\varphi$$

Real QWR Cavities

- Typical range → frequency: 50 MHz – 160 MHz; β_0 : 0.04 – 0.2

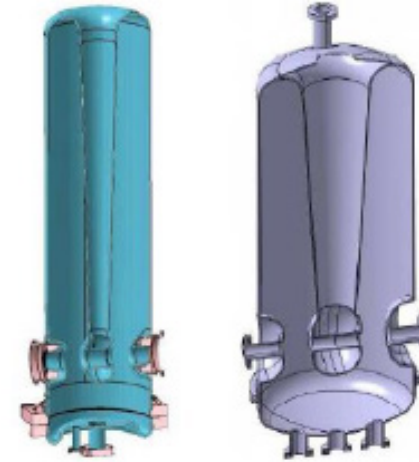
TRIUMF ISAC-II Resonators



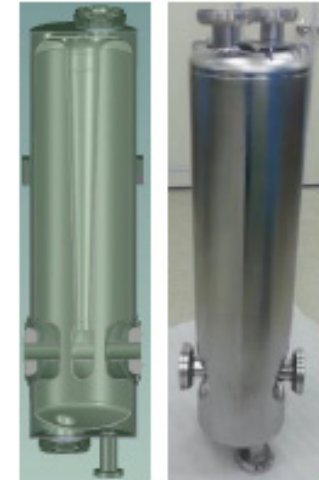
SCB low β (5.7%)
106.08 MHz

SCB medium β (7.1%)
106.08 MHz

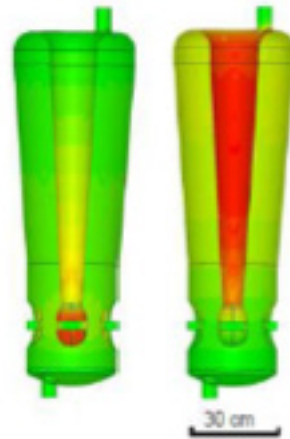
SCB high β (11%)
141.44 MHz



Spiral-2 $\beta=0.007, 0.12$
88.05 MHz



RAON $\beta=0.047$
81.25 MHz



ANL $\beta=0.077, 0.085$ 72.5 MHz



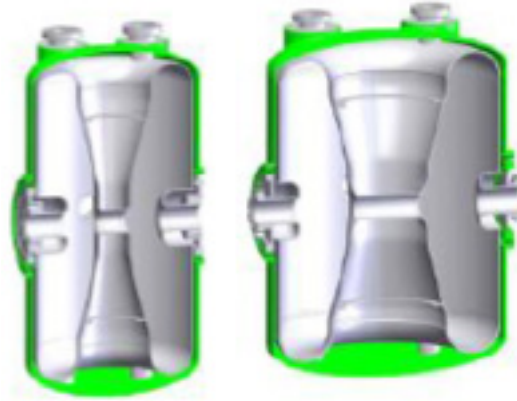
FRIB $\beta=0.041, 0.085$ 80.5 MHz

Real HWR Cavities

- Typical range \rightarrow frequency: 140 MHz – 325 MHz; β_0 : 0.1 – 0.5



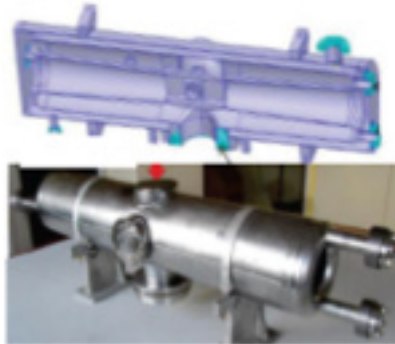
ANL $\beta=0.12$ 325 MHz



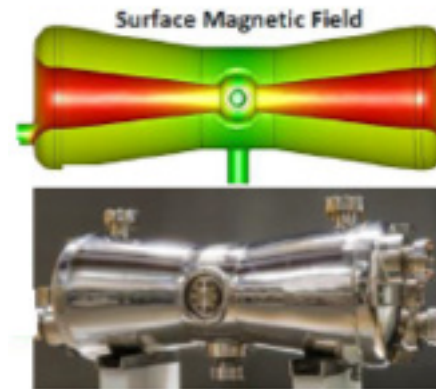
FRIB $\beta=0.29, 0.53$ 322 MHz



IMP $\beta=0.1$ 162.5 MHz



IFMIF $\beta=0.11$ 175 MHz



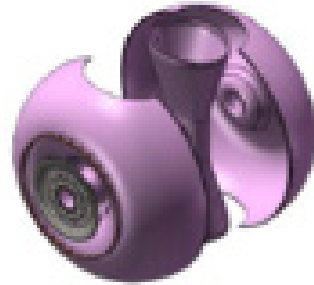
ANL $\beta=0.112$ 162.5 MHz

Real Single Spoke Cavities

- Typical range → frequency: 325 MHz – 700 MHz; β_0 : 0.15 – 0.7



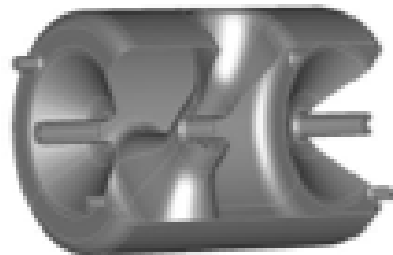
1st SC spoke 1991
ANL $\beta=0.3$ 850 MHz



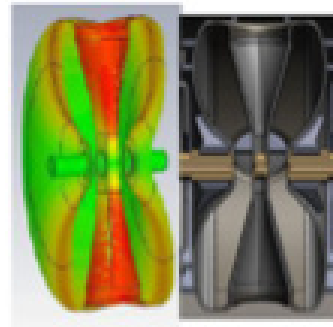
TRIUMF/RISP $\beta=0.3$ 325 MHz



FNAL $\beta=0.215$ 325 MHz



ODU $\beta=0.82$ 325 MHz



IHEP $\beta=0.12$ 325 MHz



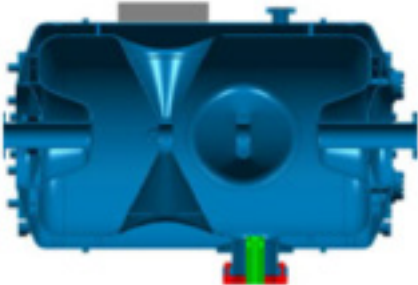
$\beta=0.15$



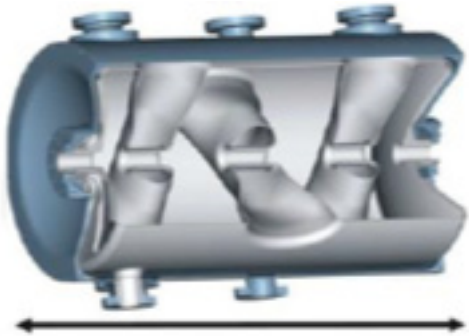
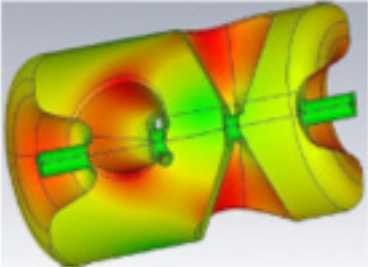
$\beta=0.35$

IPN-Orsay $\beta=0.15, 0.35$ 352 MHz

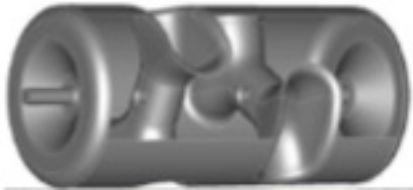
Real Multi Spoke Cavities



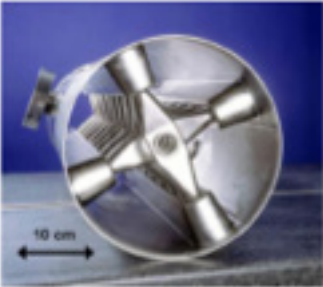
ANL $\beta=0.12$ 325 MHz



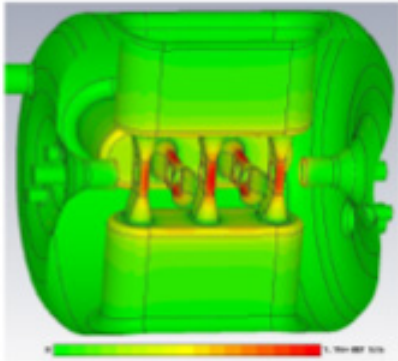
ANL $\beta=0.63$ 345 MHz



ODU $\beta=1.0$ 500 MHz

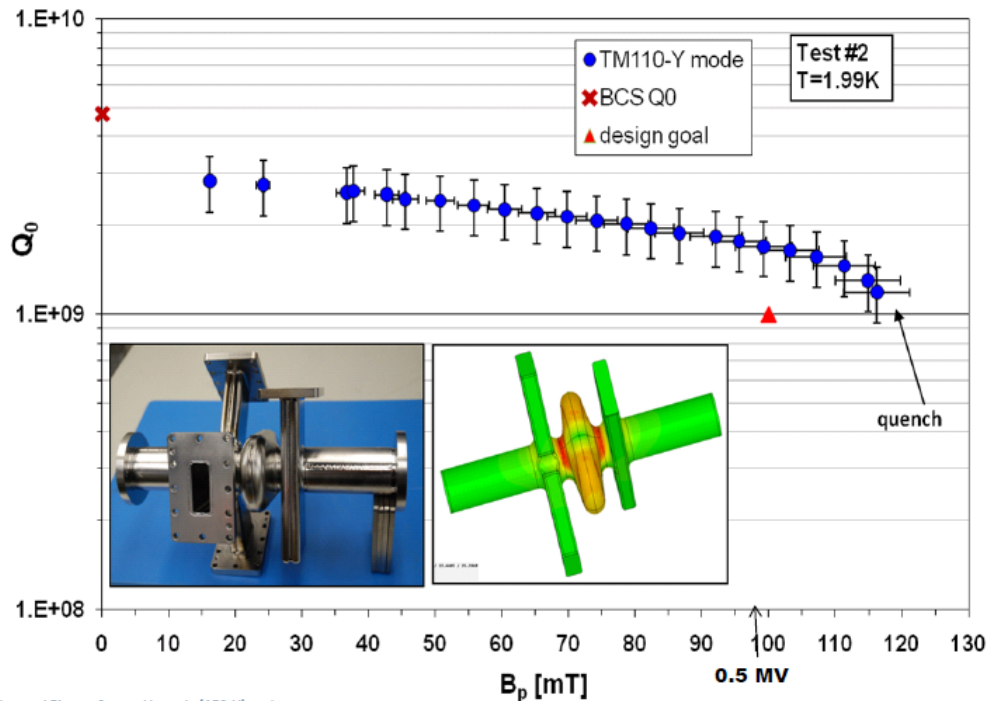
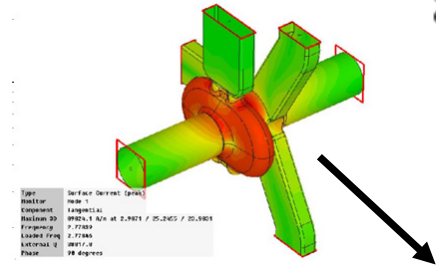
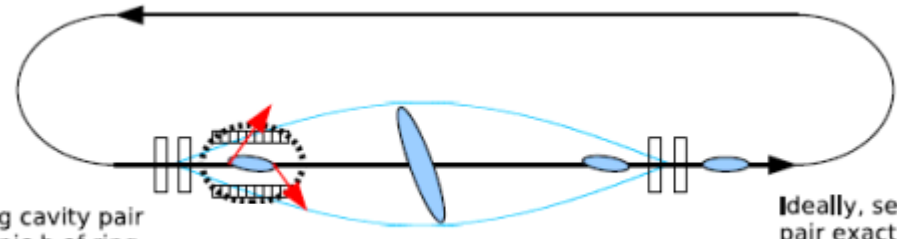
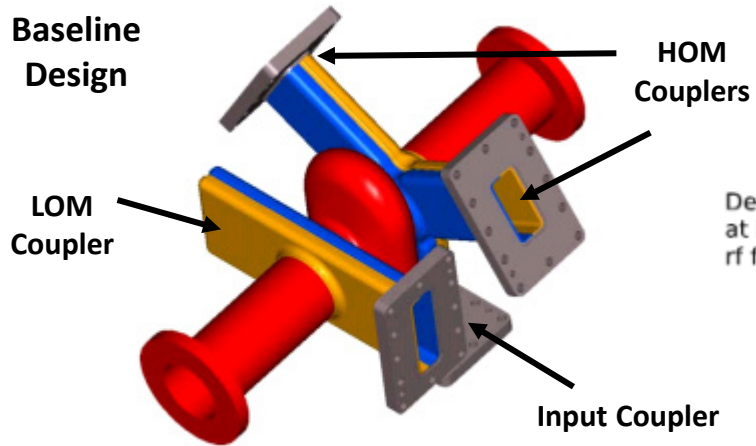


IAP $\beta \sim 0.1$ 360 MHz
19 gap CH resonator



IMP $\beta=0.067$ 162.5 MHz
CH resonator

Crabbing Cavity for Short Pulse X-ray Project

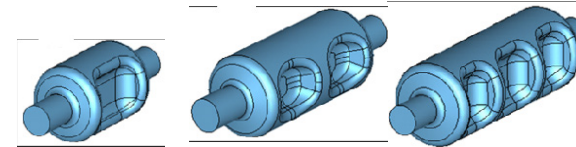


Parameter	Baseline Design	Alternate Design	Units
Frequency	2.815		GHz
Beam iris	25		mm
V_t	0.5		MV
B_t	98	100	mT
E_t	41	42	MV
$[R/Q]_t$	35.8	37.1	Ω
G	227.5	227.8	Ω
Material thickness	3 Nb Sheet	4 Nb Block	mm

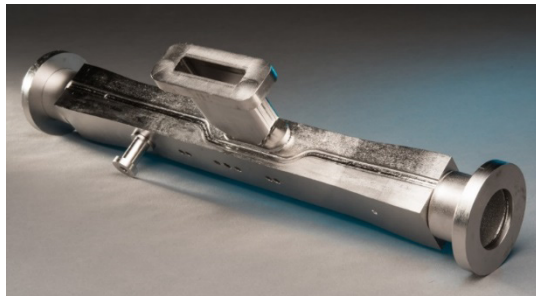
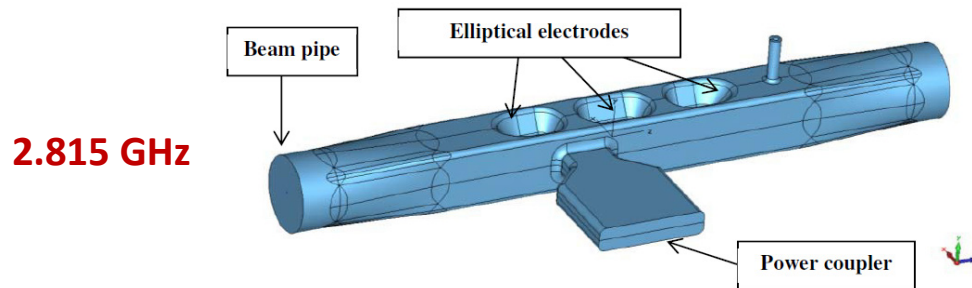
Multi-Cell TE₁₁-Like Cavities

- Multi-cell cavities provide higher gradient with reduced space on the beam line
- No of HOMs multiplies with no. of cells
- Has lower order modes

For JLEIC
e – 12 GeV p – 200 GeV



QMIR (Quasi-waveguide Multi-cell Resonator)



	Single-Cell RFD	Two-Cell RFD	Three-Cell RFD	Unit
Frequency	952.6			MHz
Aperture	70			mm
SOM	None	846	756.8, 862.2	MHz
LOM Mode Type	–	Dipole	Dipole	
1 st HOM	1411.5	1379.5	1335.4	MHz
E_p/E_t	5.4	5.66	5.6	
B_p/E_t	13.6	11.64	11.4	mT/(MV/m)
$[R/Q]_t$	50	147.5	218.8	Ω
G	165.7	169	178.9	Ω
$R_t R_s$	8.3×10^3	2.5×10^4	3.9×10^4	Ω^2
Total V_t (e/p) (per beam per side)	4.2 / 21.5			MV
V_t (per cavity)	0.86	1.9	3.1	MV
No. of cavities (e/p)	5 / 25	3 / 12	2 / 7	
E_p	30	34	39	MV/m
B_p	70	70	70	mT

TRIUMF RF-Deflector

- Due to low performance specifications, fabrication methods include some alternative techniques:

- Machining from bulk reactor grade Nb
 - RRR of 45 compared to usual ~ 300
- Tungsten Inert Gas (TIG) welding
 - Developed as an alternative to electron beam welding

Cavity performance parameters:

- Superconducting Niobium cavity at 4.2 K
- Resonant frequency: 650 MHz
- Deflecting voltage: 0.3 (0.6) MV
- Shunt impedance: 625 Ω
- Geometry factor: 99 Ω
- Peak electric field: 9.5 (19) MV/m
- Peak magnetic field: 12 (24) mT
- RF power dissipation: 0.35 (1.4) W

